

Opg. 24) Matrices in de Dirac-verg.: β en $\vec{\gamma}$, met eigenschappen $\beta^2 = I$, $\beta = \beta^\dagger$, $\{\gamma^i, \beta\} = 0$, $\{\gamma^i, \gamma^k\} = 2\delta^{jk} I$, $\gamma^j = (\gamma^j)^\dagger$ ($j, k = 1, 2, 3$) (D)

(i) $\beta = \begin{pmatrix} I & \phi \\ \phi & -I \end{pmatrix}$ en $\vec{\gamma} = \begin{pmatrix} \phi & \vec{\sigma} \\ \vec{\sigma} & \phi \end{pmatrix}$ voldoen hieraan, want $\beta^2 = I$ en $\beta = \beta^\dagger$ gelden triviaal, $\vec{\gamma} = \vec{\gamma}^\dagger$ volgt uit $\vec{\sigma} = \vec{\sigma}^\dagger$,

$$\{\vec{\gamma}, \beta\} = \begin{pmatrix} \phi & -[\vec{\sigma}, I] \\ [\vec{\sigma}, I] & \phi \end{pmatrix} = 0 \text{ en } \{\gamma^i, \gamma^k\} = \begin{pmatrix} \{\sigma^i, \sigma^k\} & \phi \\ \phi & \{\sigma^i, \sigma^k\} \end{pmatrix} = 2\delta^{jk} I.$$

(ii) Voer de Dirac γ -matrices γ^μ in, met $\gamma^0 \equiv \beta$ en $\vec{\gamma} \equiv \beta \vec{\gamma}$. Dan geldt:

$$\{\gamma^\mu, \gamma^\nu\} = \begin{cases} \mu = \nu = 0 & : 2\beta^2 \stackrel{(D)}{=} 2I \\ \mu = 0, \nu = j & : \{\beta, \beta \gamma^j\} = \beta^2 \gamma^j + \beta \gamma^j \beta = \beta \{\beta, \gamma^j\} \stackrel{(D)}{=} 0 = 2g^{\mu\nu} I \\ \mu = j, \nu = k & : \{\beta \gamma^j, \beta \gamma^k\} = \beta \gamma^j \beta \gamma^k + \beta \gamma^k \beta \gamma^j \stackrel{(D)}{=} -\beta^2 \{\gamma^j, \gamma^k\} \stackrel{(D)}{=} -2\delta^{jk} I \end{cases}$$

$$(\gamma^\mu)^\dagger = \begin{cases} \mu = 0 & : \beta^\dagger \stackrel{(D)}{=} \beta = \gamma^0 \\ \mu = j & : (\beta \gamma^j)^\dagger \stackrel{(D)}{=} \gamma^j \beta \stackrel{(D)}{=} -\beta \gamma^j = -\gamma^j \end{cases}$$

(iii) Hieruit volgt dat $(\gamma^0)^2 = \frac{1}{2} \{\gamma^0, \gamma^0\} \stackrel{(ii)}{=} I$, $(\gamma^j)^2 = \frac{1}{2} \{\gamma^j, \gamma^j\} \stackrel{(ii)}{=} -I$,

$$\gamma^0 \gamma^\mu \gamma^0 = \begin{cases} \mu = 0 & : (\gamma^0)^2 \gamma^0 \stackrel{(ii)}{=} \gamma^0 \\ \mu = j & : \gamma^0 \gamma^j \gamma^0 \stackrel{(ii)}{=} -(\gamma^0)^2 \gamma^j \stackrel{(ii)}{=} -\gamma^j \end{cases} \stackrel{(ii)}{=} (\gamma^\mu)^\dagger \text{ en}$$

$$\text{Tr}(\gamma^\mu) = \begin{cases} \mu = 0 & : \text{Tr}(\gamma^0) = -\text{Tr}([\gamma^j]^2 \gamma^0) \stackrel{(ii)}{=} \text{Tr}(\gamma^j \gamma^0 \gamma^j) \stackrel{\text{cycl.}}{=} \text{Tr}([\gamma^j]^2 \gamma^0) = 0 \\ \mu = j & : \text{Tr}(\gamma^j) = \text{Tr}([\gamma^0]^2 \gamma^j) \stackrel{(ii)}{=} -\text{Tr}(\gamma^0 \gamma^j \gamma^0) \stackrel{\text{cycl.}}{=} -\text{Tr}([\gamma^0]^2 \gamma^j) = 0 \end{cases}$$

(iv) "slash"-notatie: $\not{a} \equiv \gamma^\mu a_\mu \Rightarrow \not{a}^2 = (\gamma^\mu a_\mu)(\gamma^\nu a_\nu) = \frac{1}{2} a_\mu a_\nu \{\gamma^\mu, \gamma^\nu\} \stackrel{(ii)}{=} a_\mu a_\nu g^{\mu\nu} I = a^2 I$.

(v) Dirac-operator: $(i\hbar \not{\partial} - mc)$

Klein-Gordon operator: $(\square + m^2 c^2 / \hbar^2)$

Verband tussen beide operatoren: $(i\hbar \not{\partial} + mc)(i\hbar \not{\partial} - mc)$

$$= -\hbar^2 \not{\partial}^2 - m^2 c^2 I \stackrel{(iv)}{=} (-\hbar^2 \square - m^2 c^2) I = -\hbar^2 I (\square + m^2 c^2 / \hbar^2).$$

Opg. 25) Dirac-vergelijking en Lorentz-transformaties (rotaties + boosts):
 uit het relativiteitsprincipe volgt dat als $x'^{\mu} \stackrel{\text{inf.}}{\approx} (g^{\mu}_{\nu} + d\omega^{\mu}_{\nu}) x^{\nu}$,
 dan moet gelden dat $\psi'(x) \approx (I + s(d\omega)) \psi(x)$ met
 $[\gamma^{\rho}, s(d\omega)] = d\omega^{\rho}_{\nu} \gamma^{\nu}$ ($\rho=0, \dots, 3$) ① en $s(d\omega) = -s(-d\omega)$ ②.

Om de oplossing voor de 4×4 matrix $s(d\omega)$ te vinden ontbinden we t.o.v. de basis

(303) : $s(d\omega) = c I + c_{\mu} \gamma^{\mu} + c_{\mu\nu} \sigma^{\mu\nu} + \bar{c}_{\mu} \gamma^{\mu} \gamma^5 + \bar{c} \gamma^5$
 # basismatrices: ① ④ ⑥ ④ ①

Vervolgens gaan we de complexe coëfficiënten $c, c_{\mu}, c_{\mu\nu}, \bar{c}_{\mu}, \bar{c}$ voor $\mu, \nu=0, \dots, 3$ bepalen. Opmerking: $d\omega_{\mu\nu}, \sigma_{\mu\nu}$ en $c_{\mu\nu}$ zijn antisymmetrisch onder $\mu \leftrightarrow \nu$.

(i) $\{\gamma^{\mu}, \gamma^{\nu}\} = \gamma^{\mu} \gamma^{\nu} + \gamma^{\nu} \gamma^{\mu} = 2g^{\mu\nu} I$ geldt per definitie voor de γ -matrices
 $\Rightarrow \gamma^{\mu} \gamma^{\nu} = -\gamma^{\nu} \gamma^{\mu}$ als $\mu \neq \nu$, $\gamma^{\mu} \gamma^{\mu} = \frac{1}{2} \{\gamma^{\mu}, \gamma^{\mu}\} = g^{\mu\mu} I$ als $\mu = \nu$,
 $\gamma^{\mu} (i\gamma^0 \gamma^1 \gamma^2 \gamma^3) \equiv \gamma^{\mu} \gamma^5 = (-i)^3 \gamma^5 \gamma^{\mu} = -\gamma^5 \gamma^{\mu}$ (optewel: $\{\gamma^{\mu}, \gamma^5\} = 0, \mu=0, \dots, 3$),
 $\sigma^{\mu\nu} = \frac{i}{2} [\gamma^{\mu}, \gamma^{\nu}] = \frac{i}{2} \gamma^{\mu} \gamma^{\nu} - \frac{i}{2} \gamma^{\nu} \gamma^{\mu} = i\gamma^{\mu} \gamma^{\nu}$ als $\mu \neq \nu$, $\sigma^{\mu\nu} = 0$ als $\mu = \nu$.

(ii) ① met $\rho=0$ en $d\omega^0_{\alpha}$ anti-symm. ① (i) $c * 0 + c_0 * 0 + \sum_{k=1}^3 c_k * 2\gamma^k + \sum_{j,k=1}^3 c_{jk} * 0$
 $+ \sum_{k=1}^3 (c_{0k} * 2i\gamma^k - c_{k0} * 2i\gamma^k) + \bar{c}_0 * 2\gamma^5 + \sum_{k=1}^3 \bar{c}_k * 0 + \bar{c} * 2\gamma^5 \stackrel{①}{=} \sum_{k=1}^3 d\omega_{0k} \gamma^k = \sum_{k=1}^3 d\omega_{0k} \gamma^k$
 antisymm.: $4i c_{0k} \gamma^k$

basis $\Rightarrow c_k = \bar{c}_0 = \bar{c} = 0$ en $c_{0k} = -\frac{i}{4} d\omega_{0k}$ ($k=1, 2, 3$).

(iii) ① met $\rho=j$ en $d\omega^j_{\alpha} = 0$ (i), (ii) $c * 0 + c_0 * 2\gamma^j + \sum_{k \neq j} c_k * 0 + \sum_{k \neq j} c_{kj} * 0$
 $+ \sum_{k \neq j} (c_{jk} * 2i\gamma^k - c_{kj} * 2i\gamma^k) + c_{j0} * 2i\gamma^0 + \sum_{k \neq j} \bar{c}_k * 0 + \bar{c}_j * (-2\gamma^5)$
 antisymm.: $-4i c_{jk} \gamma^k$ antisymm.: $-4i c_{j0} \gamma^0$

① $d\omega^j_{\alpha} \gamma^{\alpha} + \sum_{k \neq j} d\omega^j_{\alpha k} \gamma^k = -d\omega^j_{\alpha 0} \gamma^0 - \sum_{k \neq j} d\omega^j_{\alpha k} \gamma^k$ basis $\Rightarrow c_0 = \bar{c}_j = 0$ en $c_{jk} = -\frac{i}{4} d\omega_{jk}$ ($j, k=1, 2, 3$)

(iv) We hebben nu gevonden dat $c_{\mu\nu} = -\frac{i}{4} d\omega_{\mu\nu}$, maar de coëfficiënt c is nog steeds onbepaald. Echter, uit conditie ② volgt een extra voorwaarde:

$s(d\omega) = c I - \frac{i}{4} d\omega_{\mu\nu} \sigma^{\mu\nu} = -s(-d\omega) = -c I - \frac{i}{4} d\omega_{\mu\nu} \sigma^{\mu\nu} \Rightarrow c = 0$, zodat
 $s(d\omega) = -\frac{i}{4} d\omega_{\mu\nu} \sigma^{\mu\nu}$
 (circled note: $\sigma^{\mu\nu}$ genereert inf. Lor. transp.)