

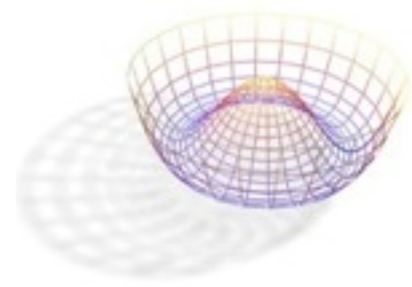
The Search for the Higgs Boson: Discovery in Sight!

Frank Filthaut
Radboud Universiteit Nijmegen / Nikhef

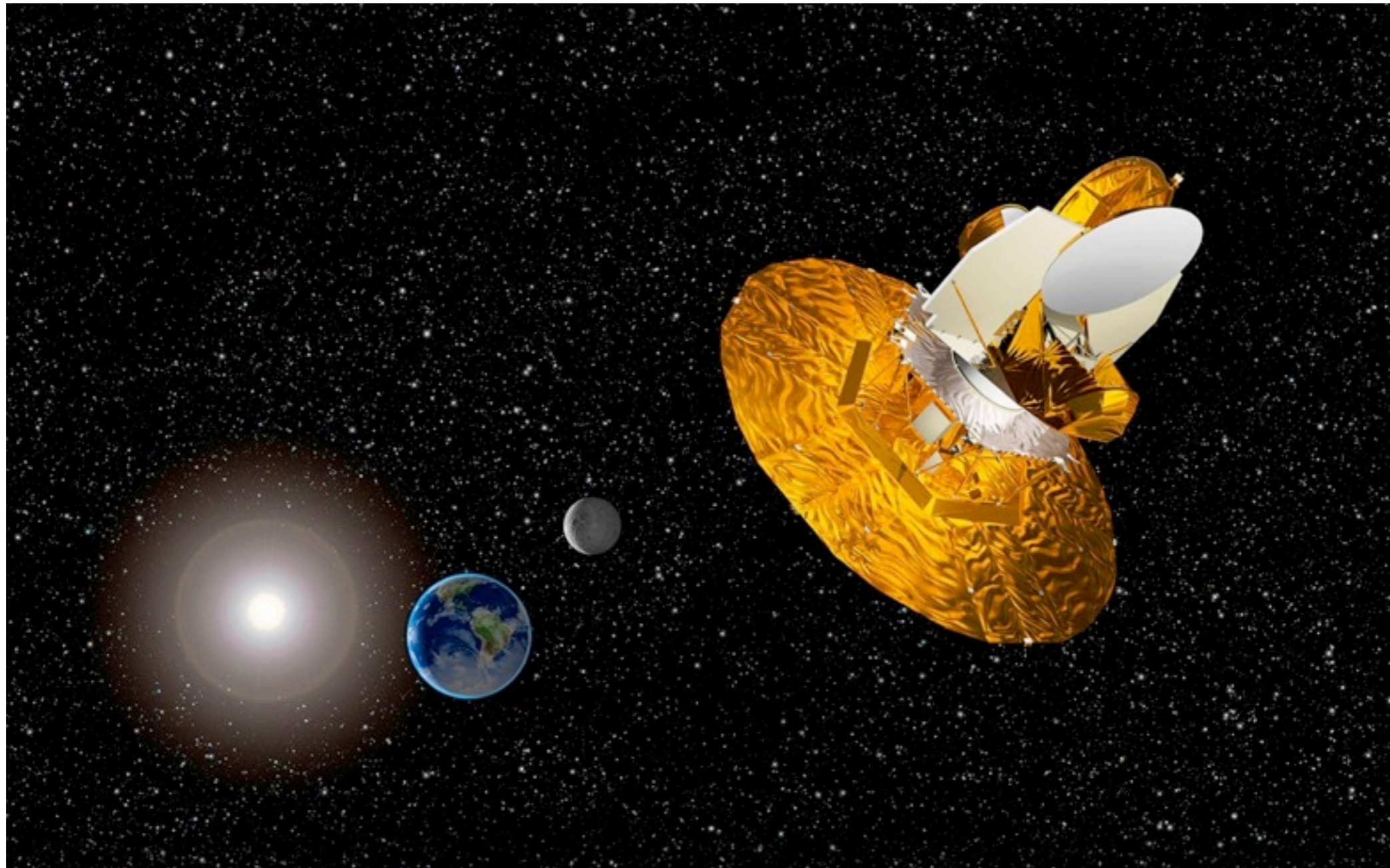
Particle physics: general picture
The Higgs boson
Higgs hunting at the LHC

Particle Physics: the General Picture

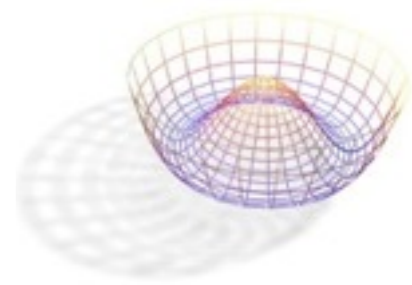
Rotational Symmetry



The Universe at large is isotropic and homogeneous!

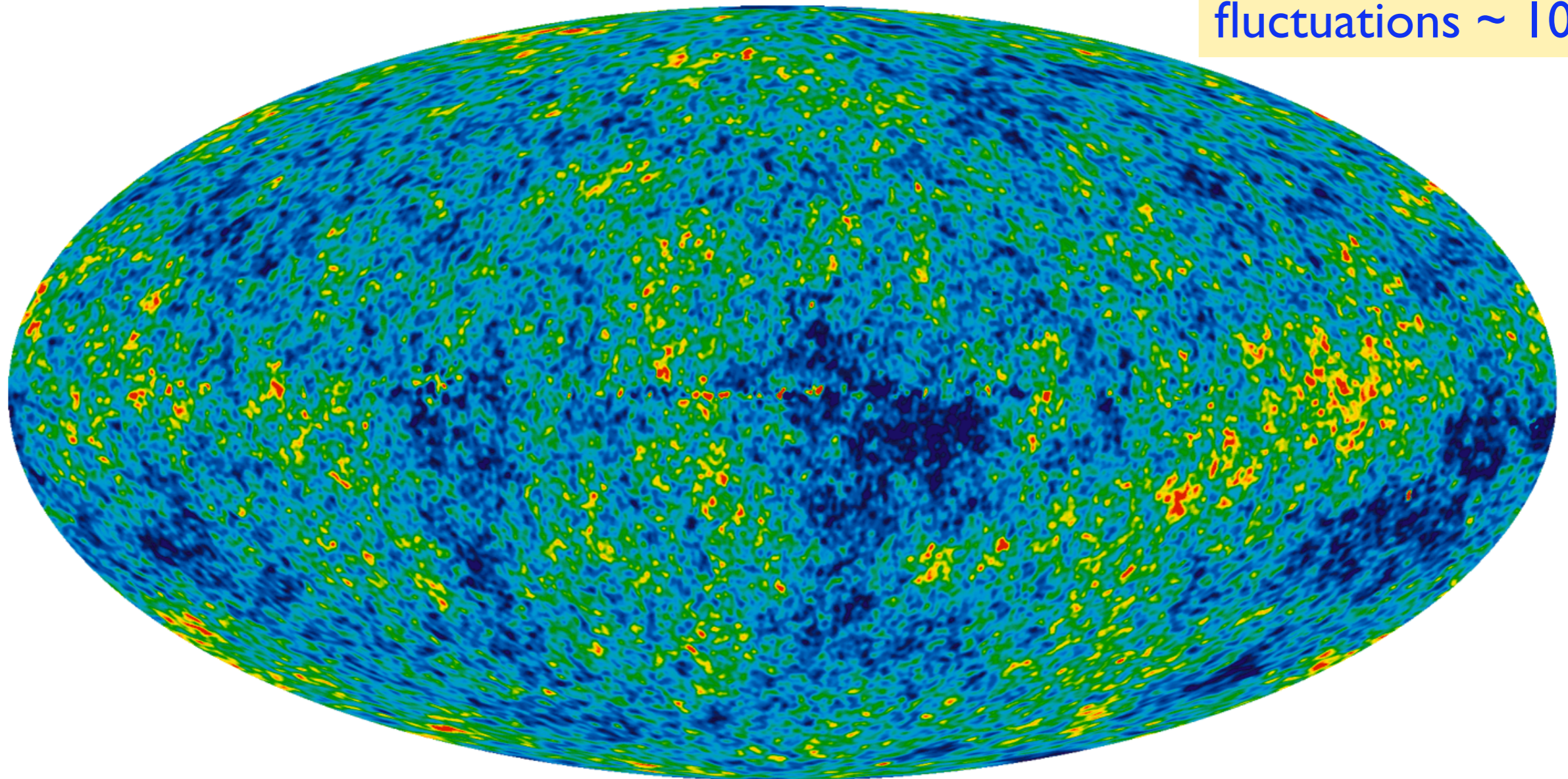


Rotational Symmetry



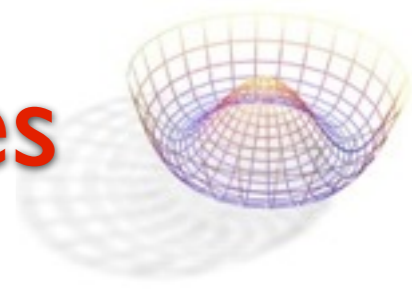
The Universe at large is isotropic and homogeneous!

CMWB: temperature
fluctuations $\sim 10^{-5}$ K



WMAP 7-year results, full sky

Symmetries and Conserved Quantities



Rotational symmetry: laws of physics do not depend on any direction.
Symmetries are important in many areas of physics

- e.g. conserved quantities like angular momentum in the case of rotational symmetry

Particle physics extends these concepts to internal symmetries, preserved even under arbitrary space-time dependent (gauge) transformations

Allows for extraordinarily successful description of electromagnetism: QED

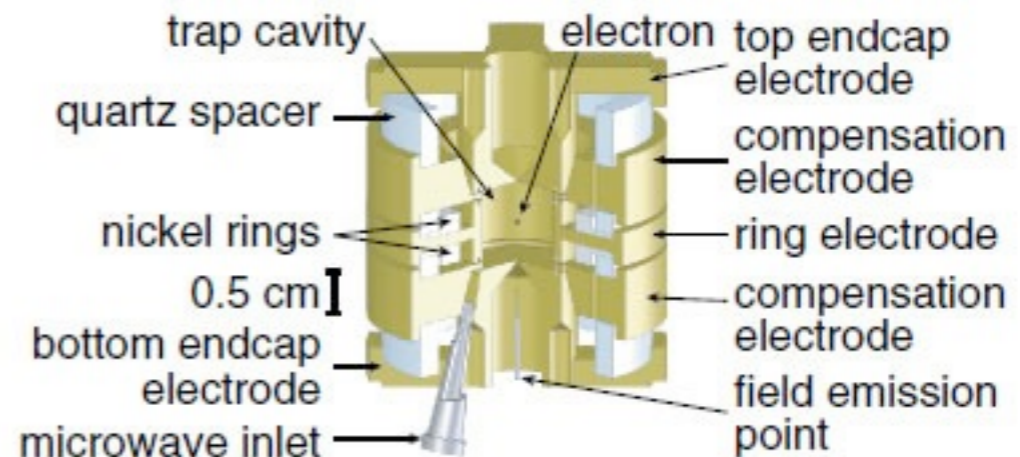
- interaction of magnetic dipole moment with external magnetic field:

$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

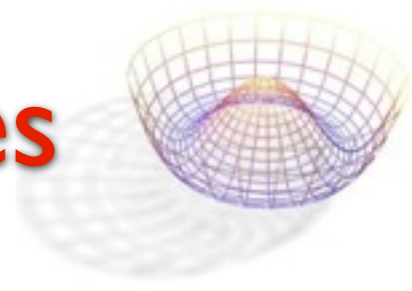
In contrast to “ordinary” QM, g can be computed from first principles! Only input:

$$\alpha \equiv \frac{e^2}{4\pi}$$

Observe trapped single electron for months



Symmetries and Conserved Quantities



Rotational symmetry: laws of physics do not depend on any direction.
Symmetries are important in many areas of physics

- e.g. conserved quantities like angular momentum in the case of rotational symmetry

Particle physics extends these concepts to internal symmetries, preserved even under arbitrary space-time dependent (gauge) transformations

Allows for extraordinarily successful description of electromagnetism: QED

- interaction of magnetic dipole moment with external magnetic field:

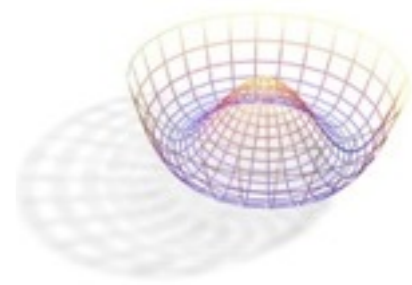
$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

In contrast to “ordinary” QM, g can be computed from first principles! Only input:

$$\alpha \equiv \frac{e^2}{4\pi}$$

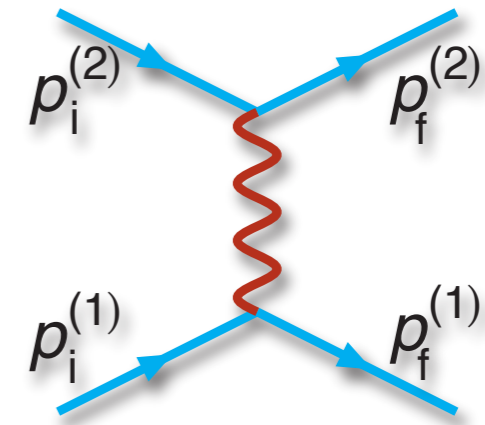
$$g/2 = \begin{cases} 1.001\,159\,652\,180\,73(28) & \text{(experiment)} \\ 1.001\,159\,652\,180\,85(76) & \text{(theory)} \end{cases}$$

Particle Paradigm

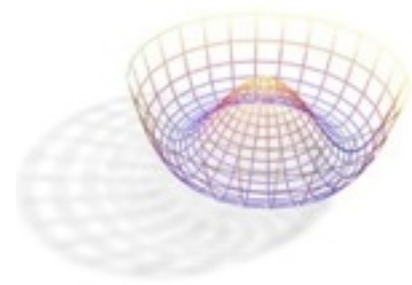


QED implements EM interaction as exchange of (massless) photon

- same for (massless) gluons as force carriers of strong interaction



Particle Paradigm

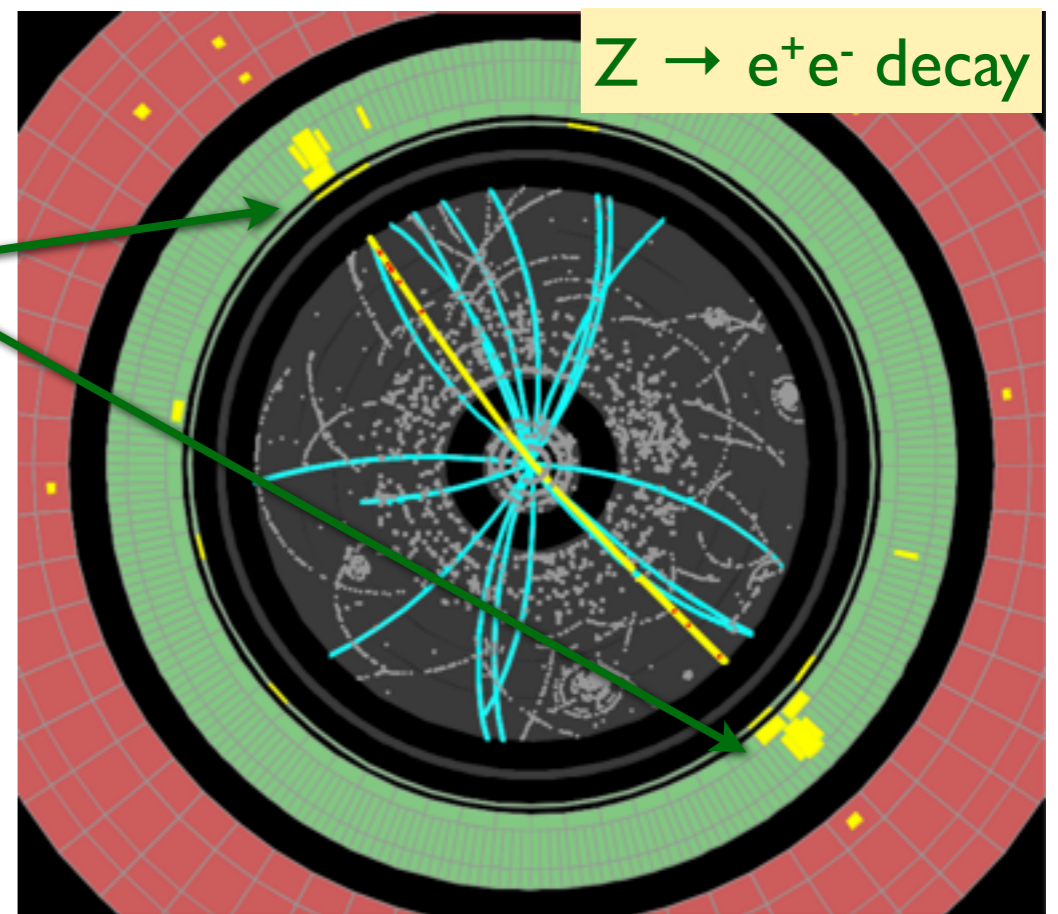
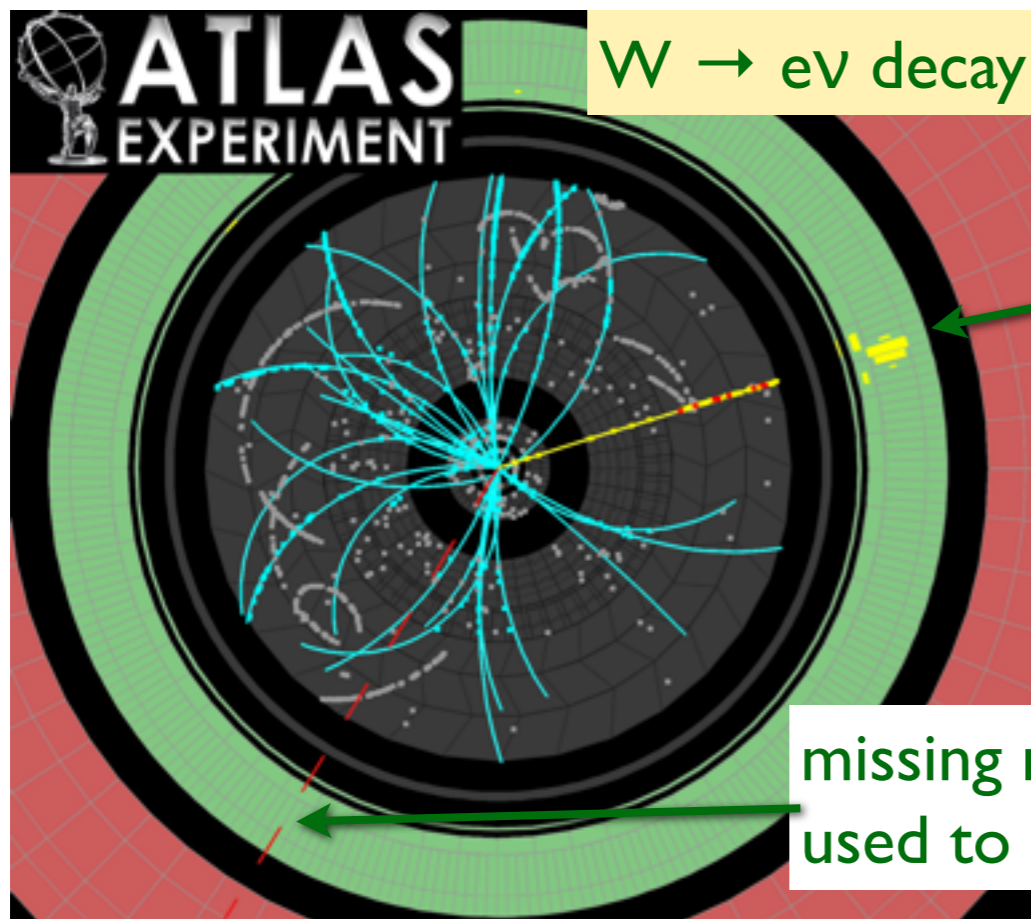
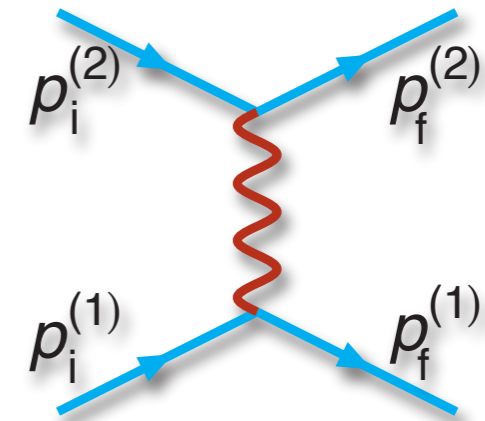


QED implements EM interaction as exchange of (massless) photon

- same for (massless) gluons as force carriers of strong interaction

The same paradigm also applies to the weak interaction, but the W and Z bosons are heavy!

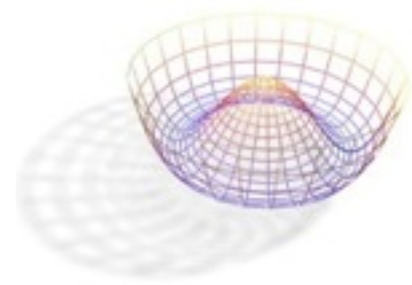
- $M_W = 80.387 \pm 0.017 \text{ GeV}$
- $M_Z = 91.188 \pm 0.002 \text{ GeV}$



e^\pm

missing momentum used to infer ν

Particle Summary

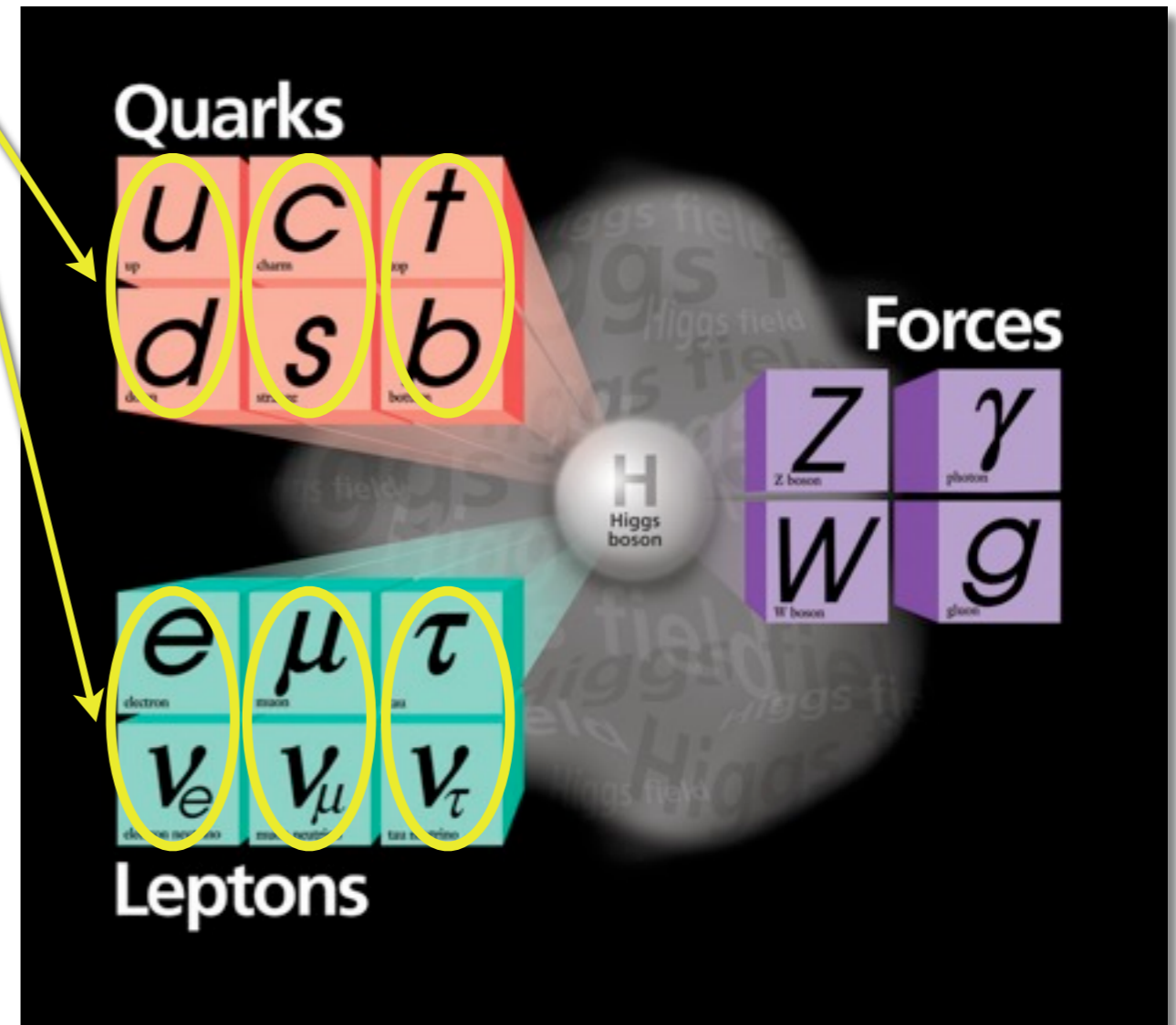


Internal symmetries can turn particles into one another

- most clearly visible for weak interaction
- but this symmetry must be broken!
For the masses of partners in doublets of the weak symmetry are different
- force carriers mediating interactions should be massless: clearly invalidated by heavy W and Z particles

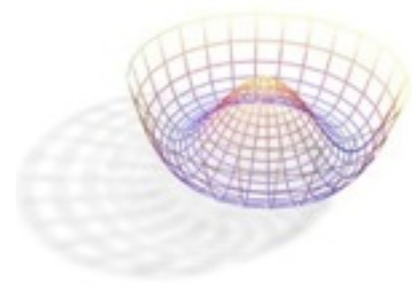
Consequence of introducing the Higgs field:

- interactions obey symmetry
 ▮ theory remains meaningful
- ground state does not:
 spontaneous symmetry breaking



The Standard Model is invalid without the Higgs boson!

Particle Summary



Internal symmetries can turn particles into one another

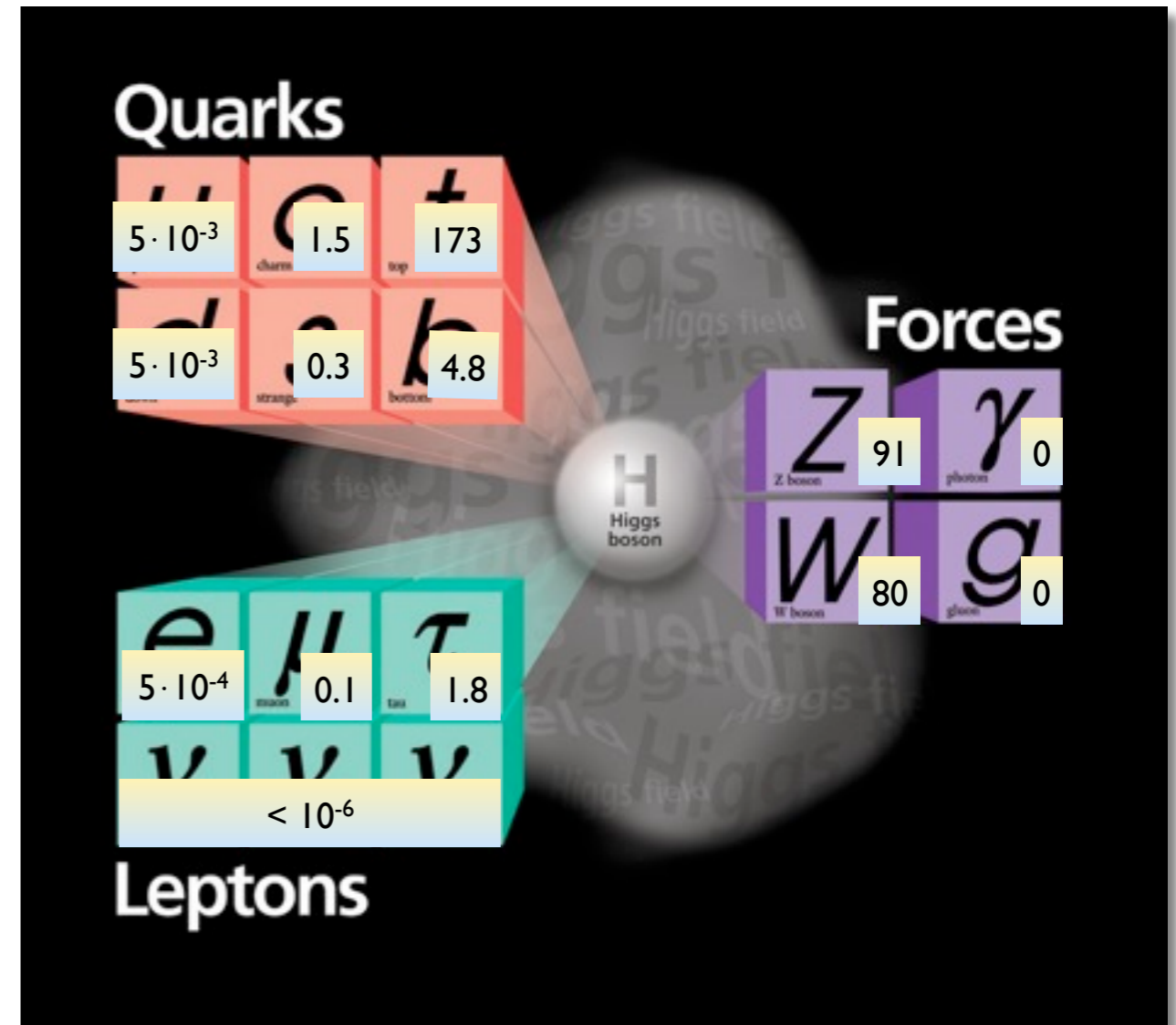
(masses in GeV)

- most clearly visible for weak interaction
- but this symmetry must be broken!
For the masses of partners in doublets of the weak symmetry are different
- force carriers mediating interactions should be massless: clearly invalidated by heavy W and Z particles

Consequence of introducing the Higgs field:

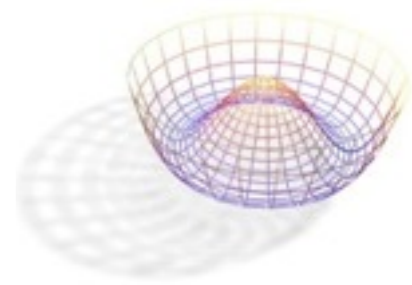
- interactions obey symmetry
 ▣ theory remains meaningful
- ground state does not:
 spontaneous symmetry breaking

The Standard Model is invalid without the Higgs boson!

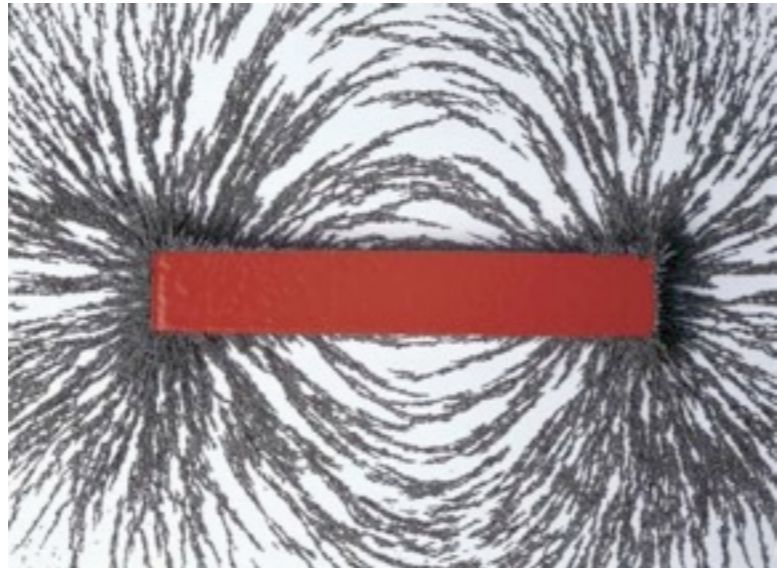


The Higgs Boson

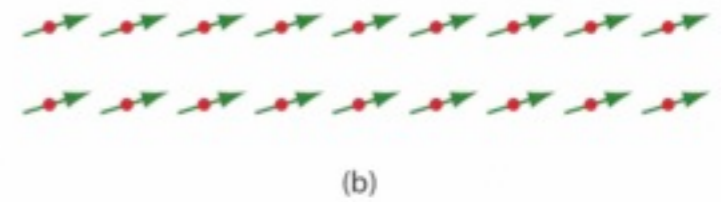
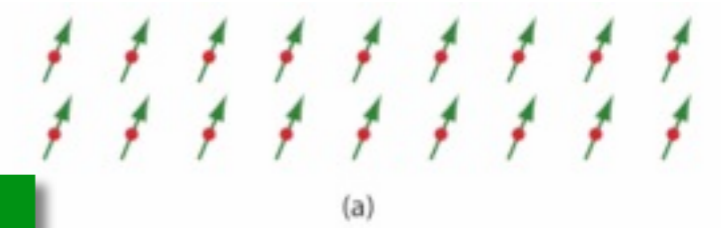
Magnetic Analogues



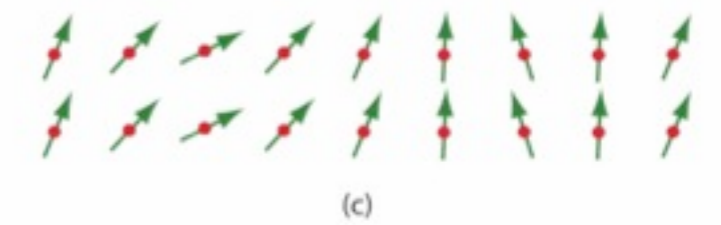
Spontaneous symmetry breaking



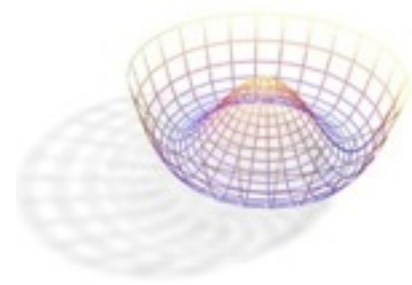
equivalent
ground states



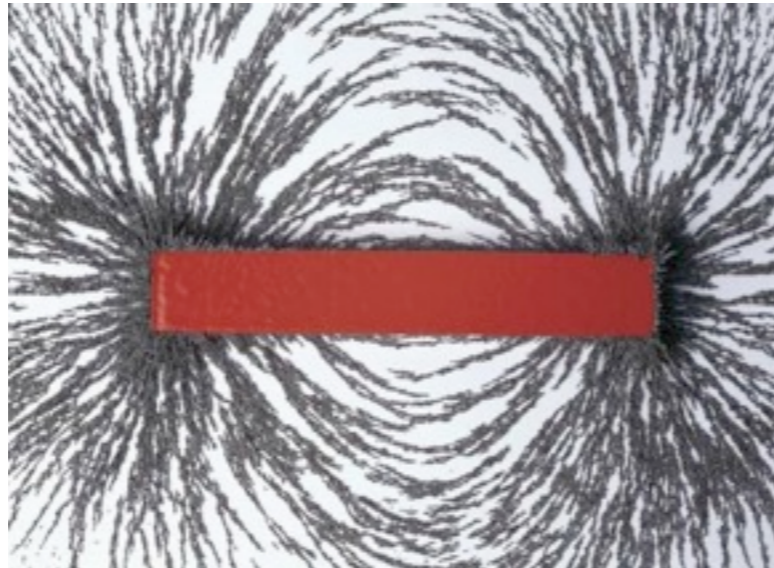
excitation



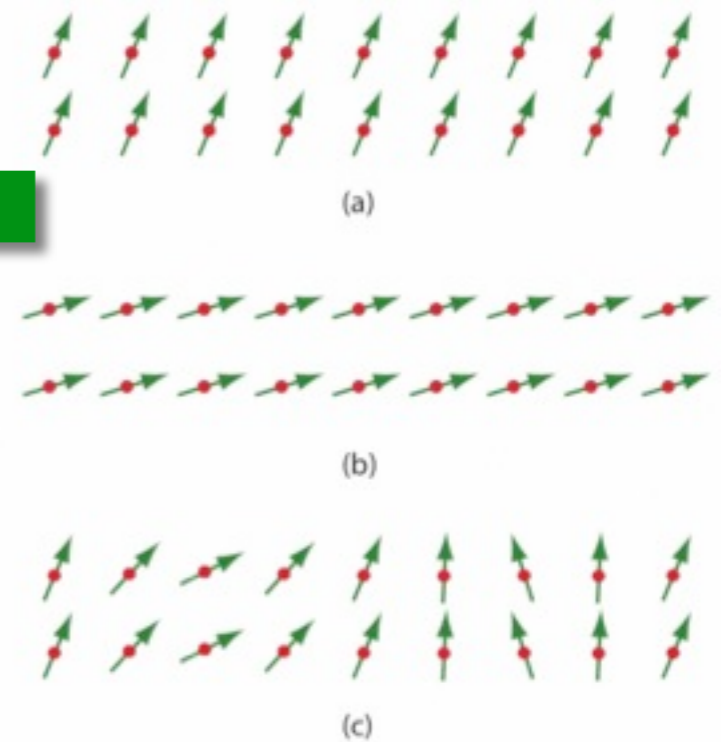
Magnetic Analogues



Spontaneous symmetry breaking

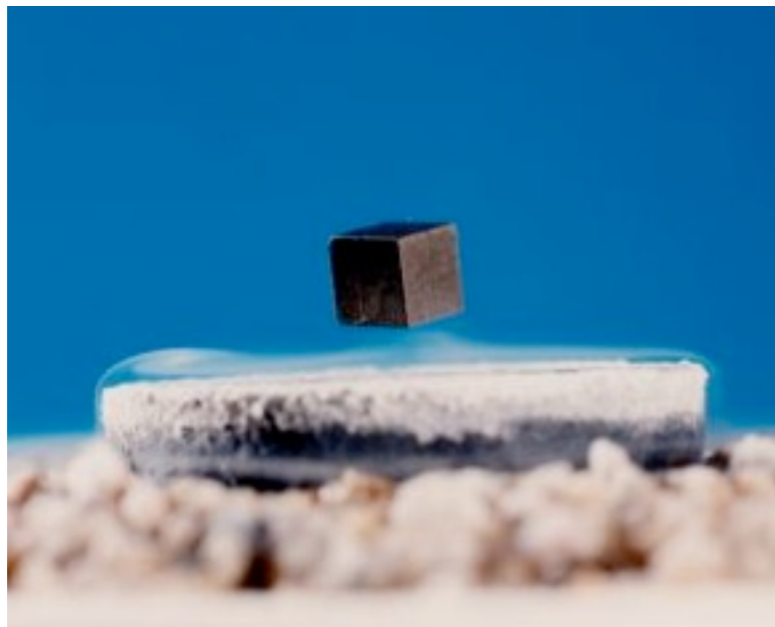


equivalent ground states



excitation

Massive photons

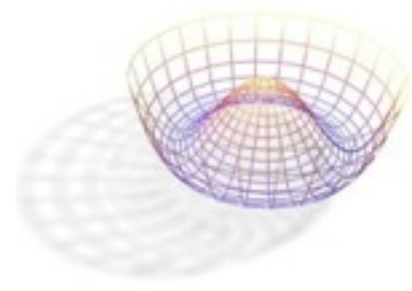


Meißner effect: superconductor repels magnetic field lines

- massive photons
- but needs a medium (e^- pair condensate)!

In the particle physics case, the “medium” is the vacuum!

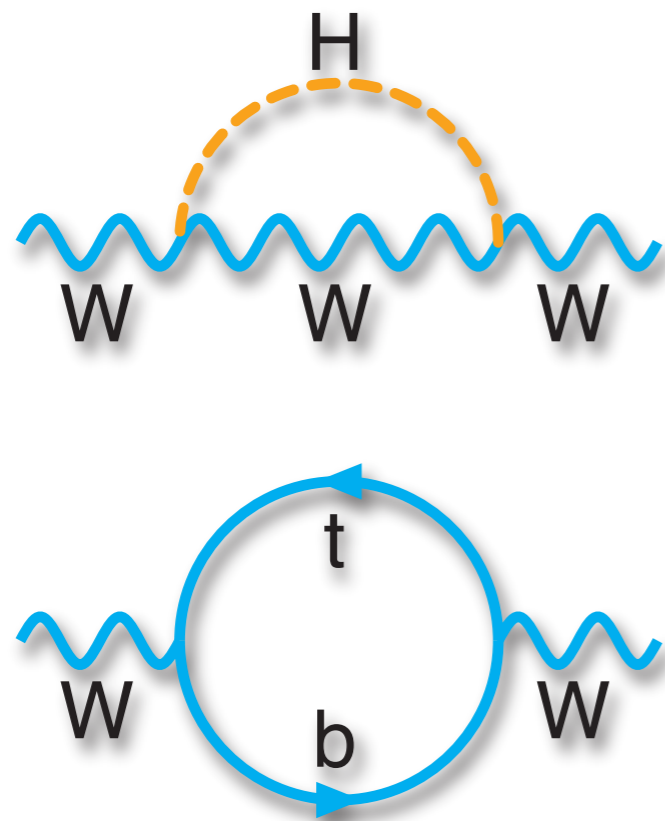
Electroweak Constraints



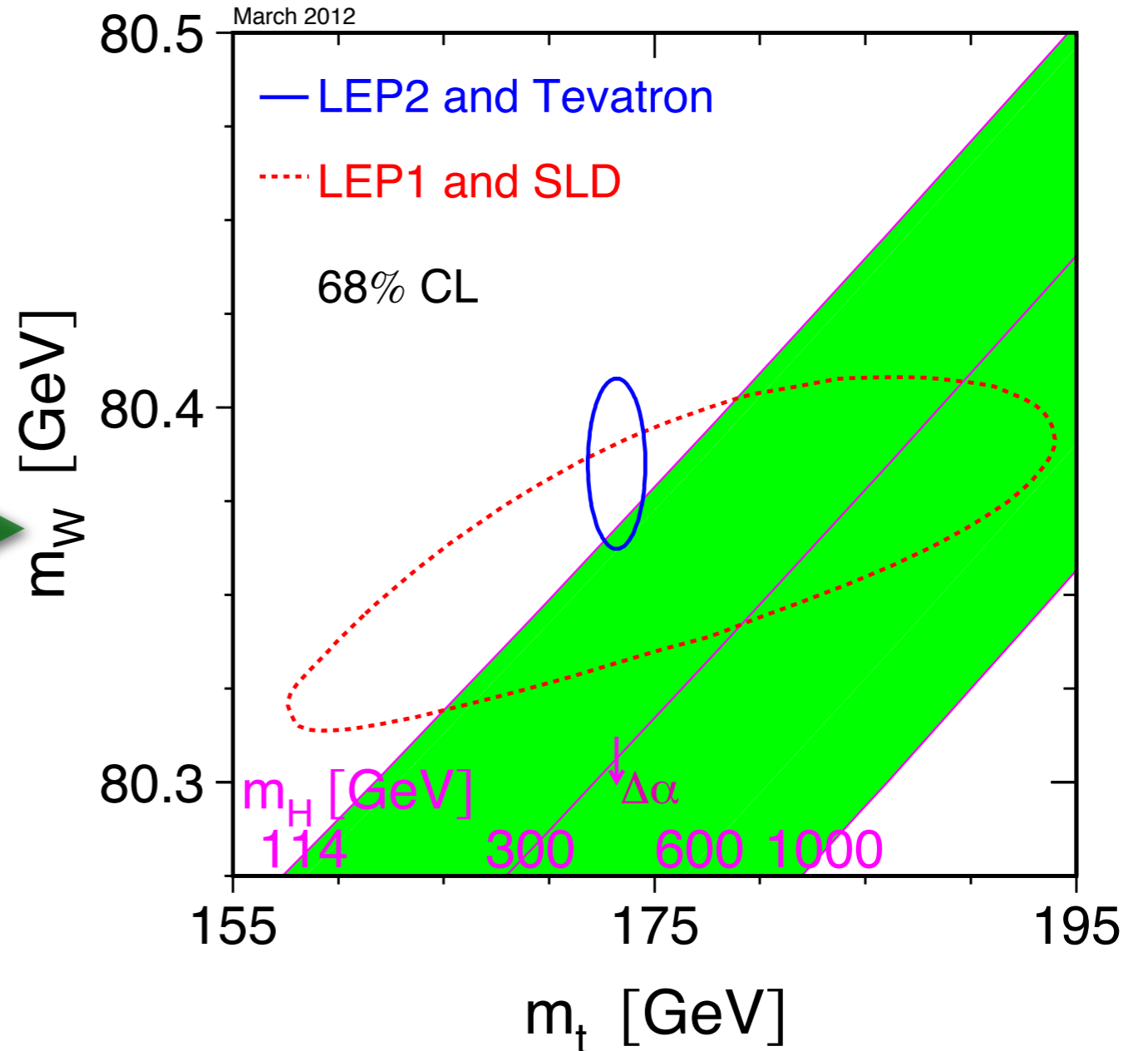
M_H unknown, but for given M_H all Higgs boson properties are fixed

⇒ know “exactly” what to look for

- many constraints! Most important: masses of W boson and top quark

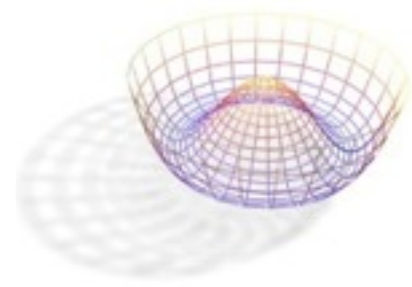


small modifications to M_W from radiative corrections

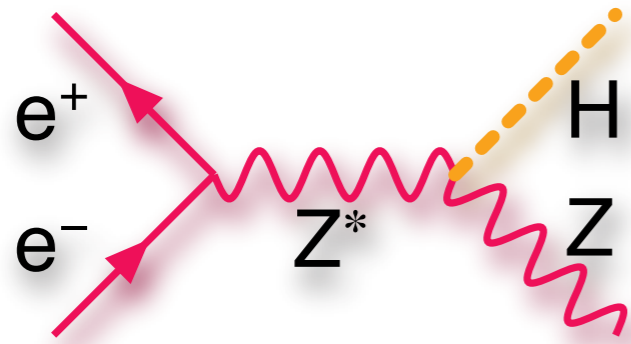


Preference for a “light” Higgs boson!

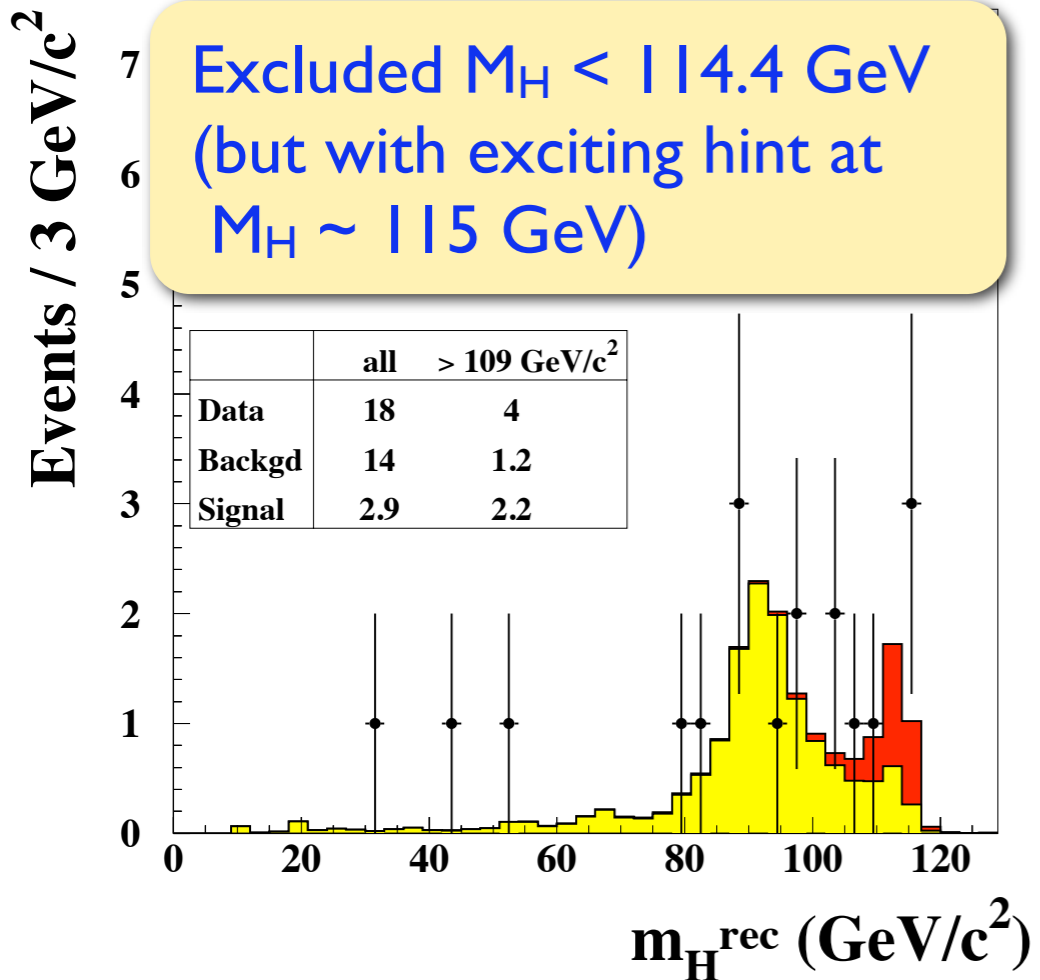
Previous Higgs Boson Searches



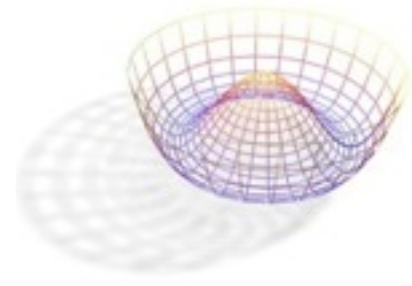
LEP: e^+e^- , $E_{CM} < 210$ GeV



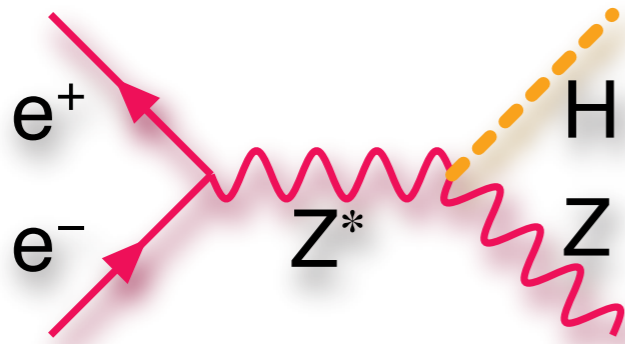
Excluded $M_H < 114.4$ GeV
(but with exciting hint at
 $M_H \sim 115$ GeV)



Previous Higgs Boson Searches

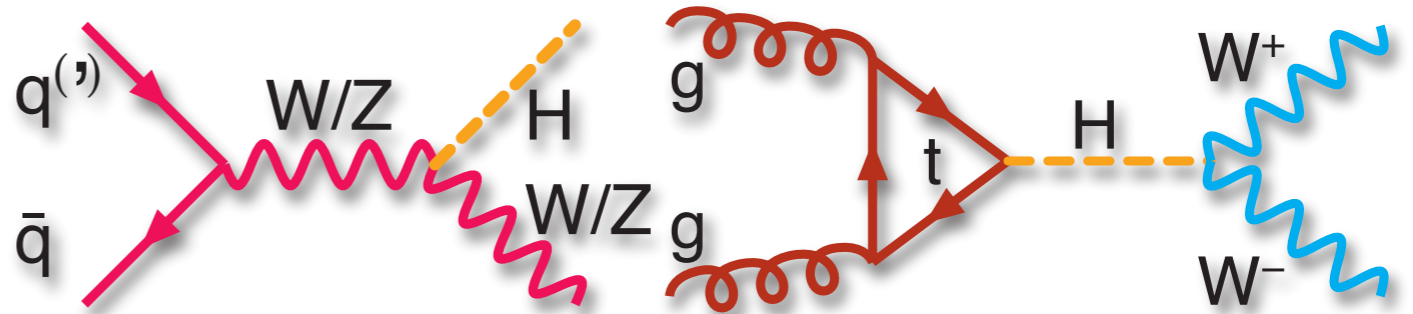


LEP: e^+e^- , $E_{CM} < 210$ GeV

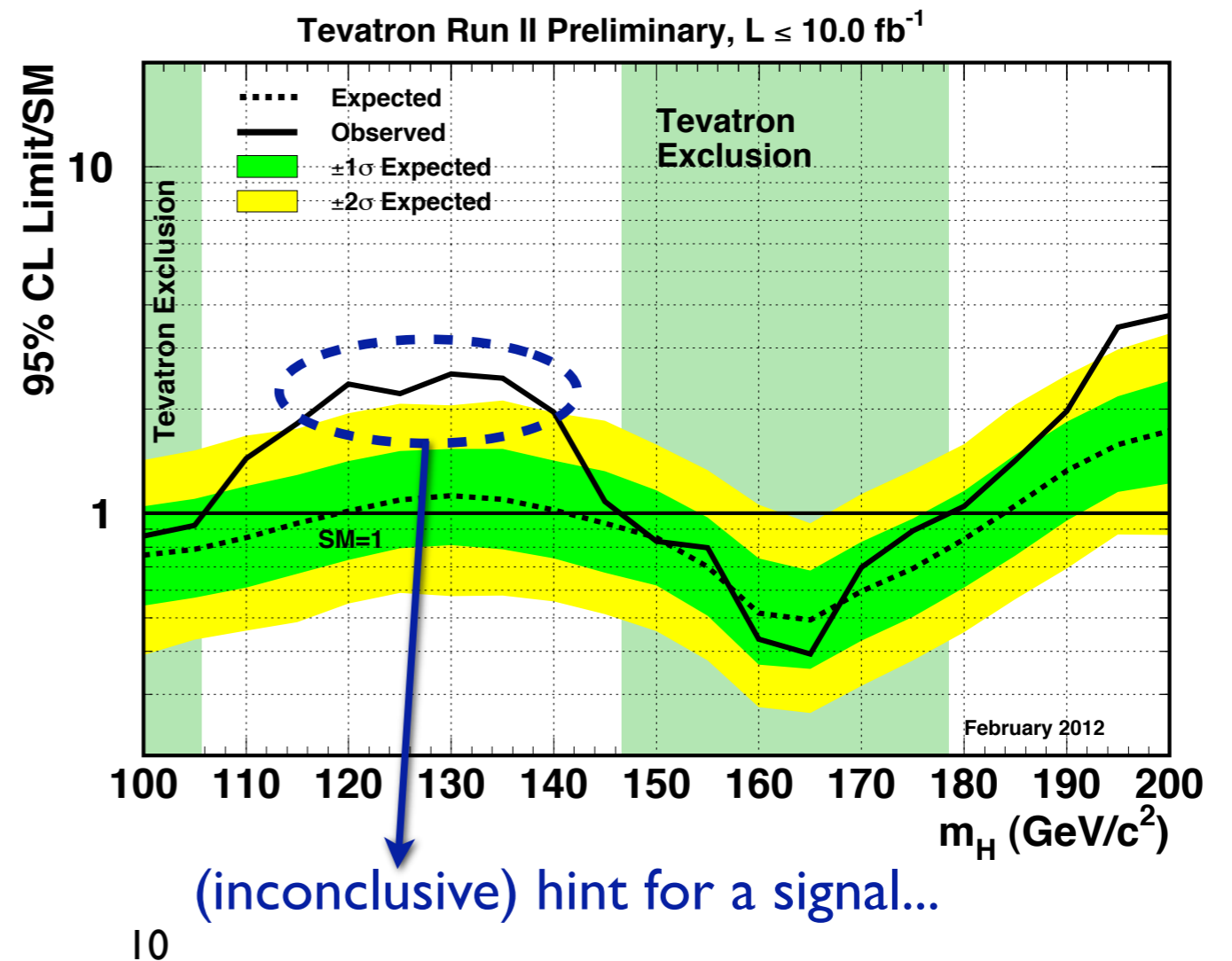
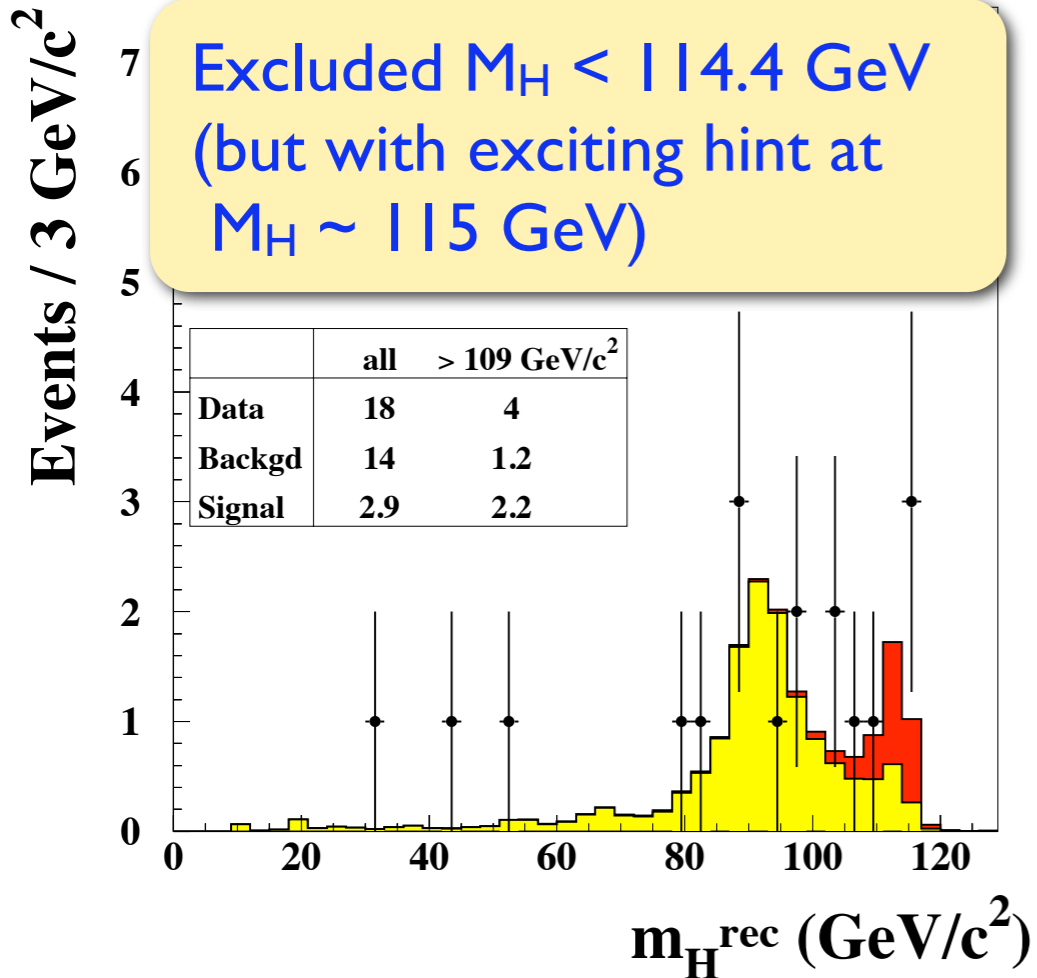


Tevatron: $p\bar{p}$, $E_{CM} = 1.96$ TeV

• but much lower E_{CM} of colliding partons!

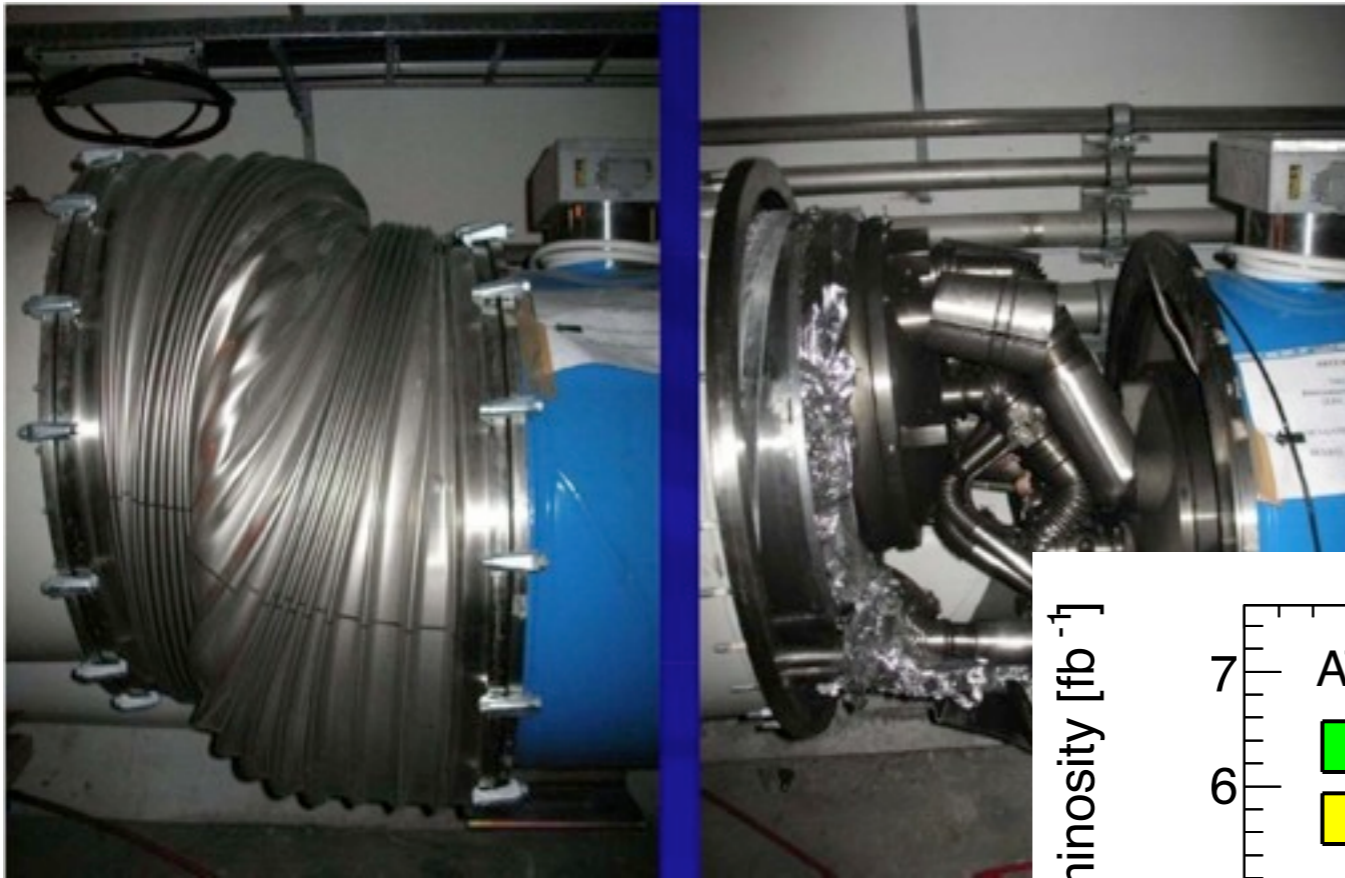
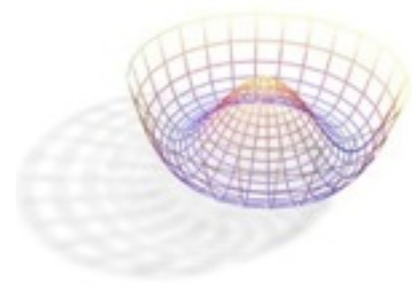


Excluded $M_H < 114.4$ GeV
(but with exciting hint at $M_H \sim 115$ GeV)



Higgs Hunting at the LHC

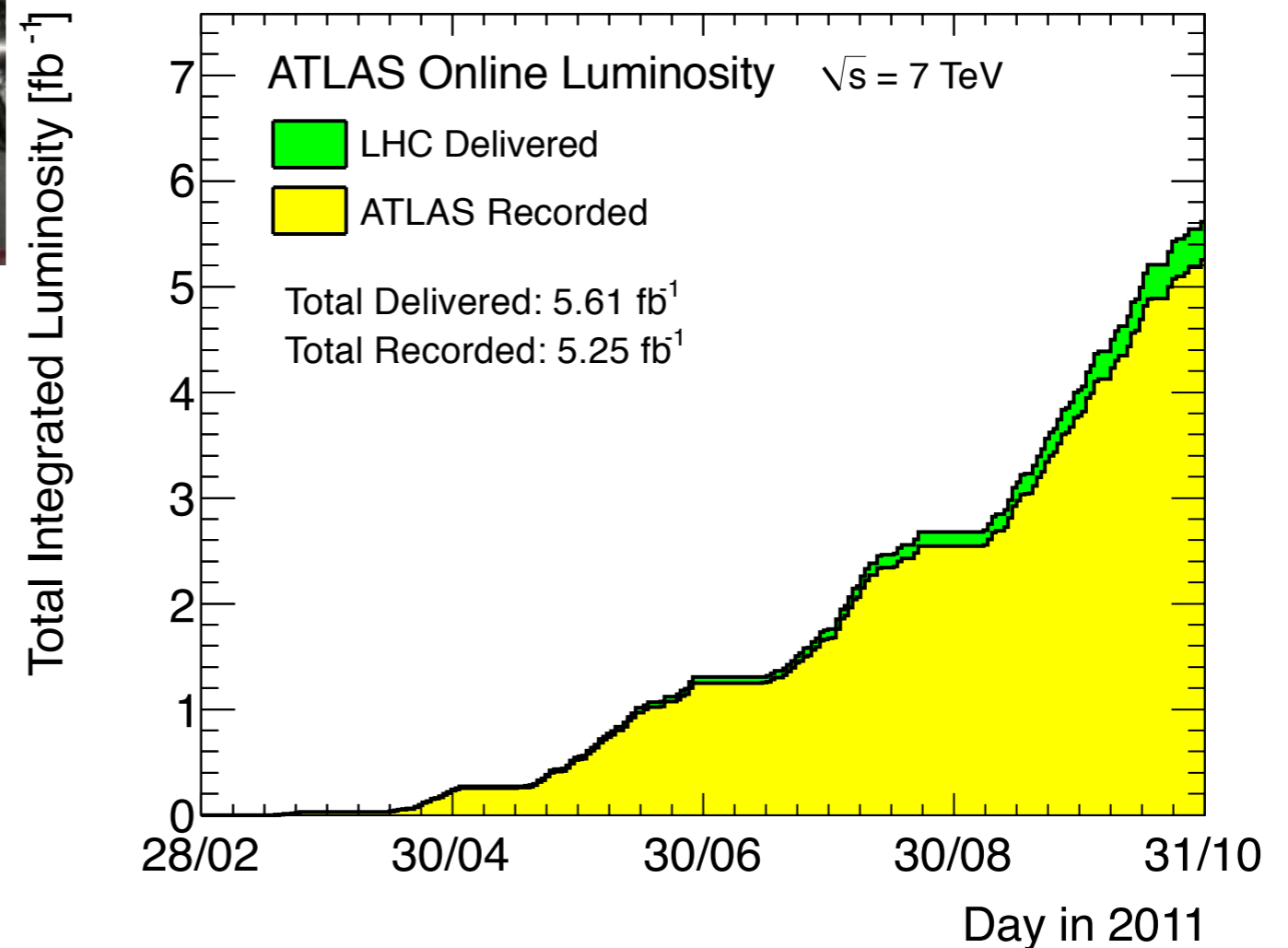
The LHC: a Success Story!



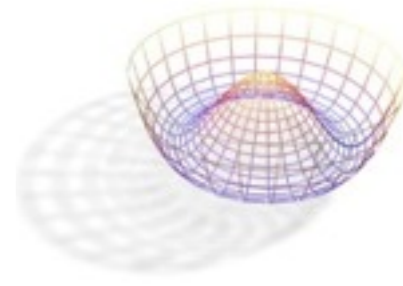
1 fb⁻¹: integrated “intensity”
needed to produce 1 interaction
for a process with a production
cross section of 1 fb = 10⁻⁴³ m²

Expectations for 2011
exceeded by a factor 5

- even if at half the CM energy initially foreseen



Experimental conditions



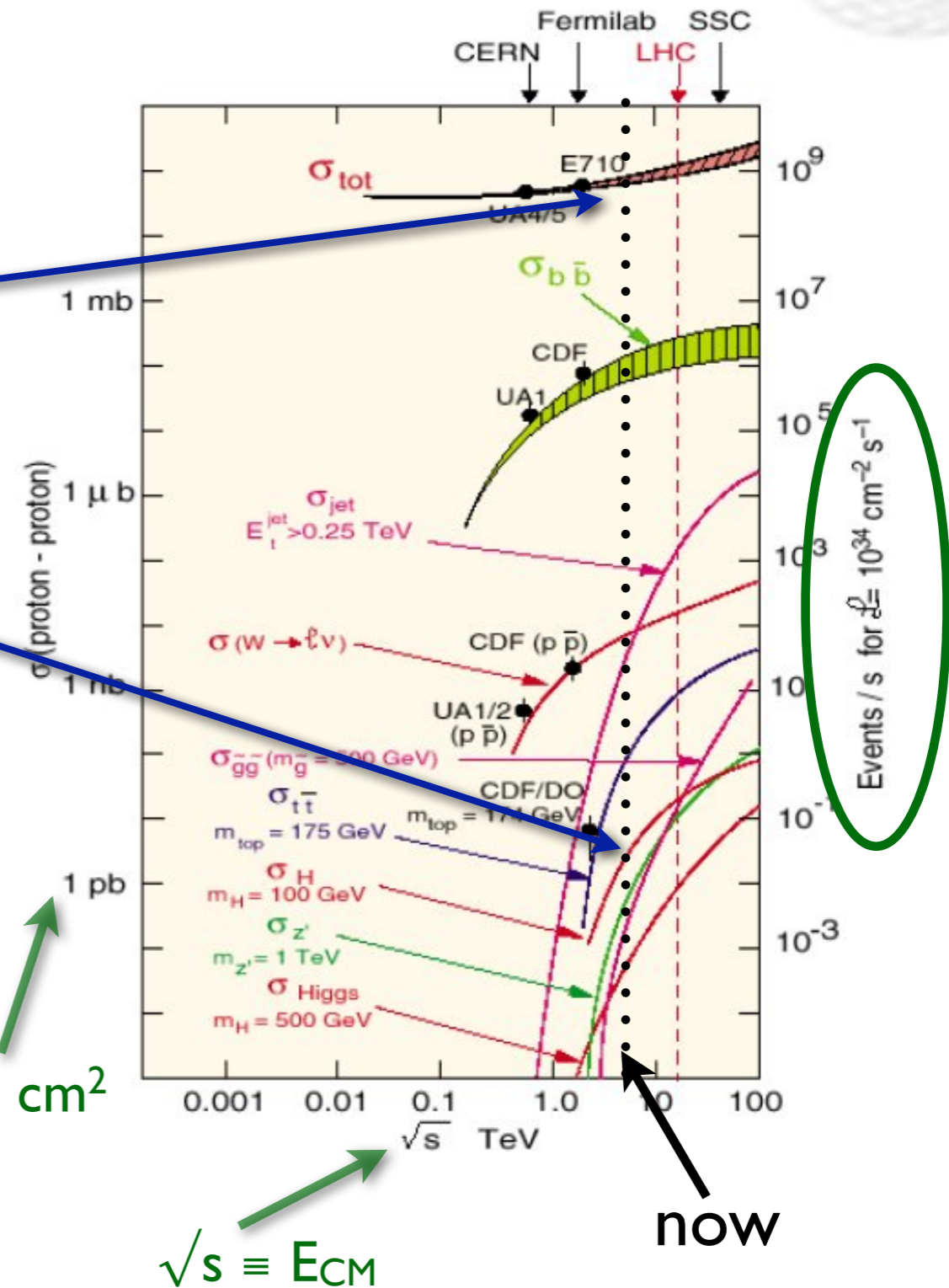
We need this performance!

- interactions at hadron colliders dominated by strong interaction
- when searching for Higgs boson production, need to **suppress backgrounds by $\sim 10^{10}$**

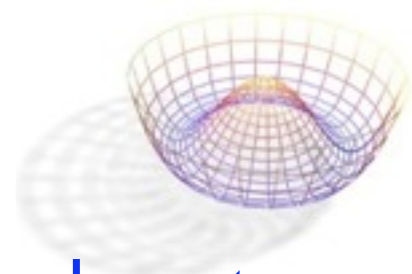
Look for striking signatures setting the Higgs boson apart from more ubiquitous "background" processes

$1 \text{ pb} = 10^{-36} \text{ cm}^2$

$\sqrt{s} \equiv E_{\text{CM}}$

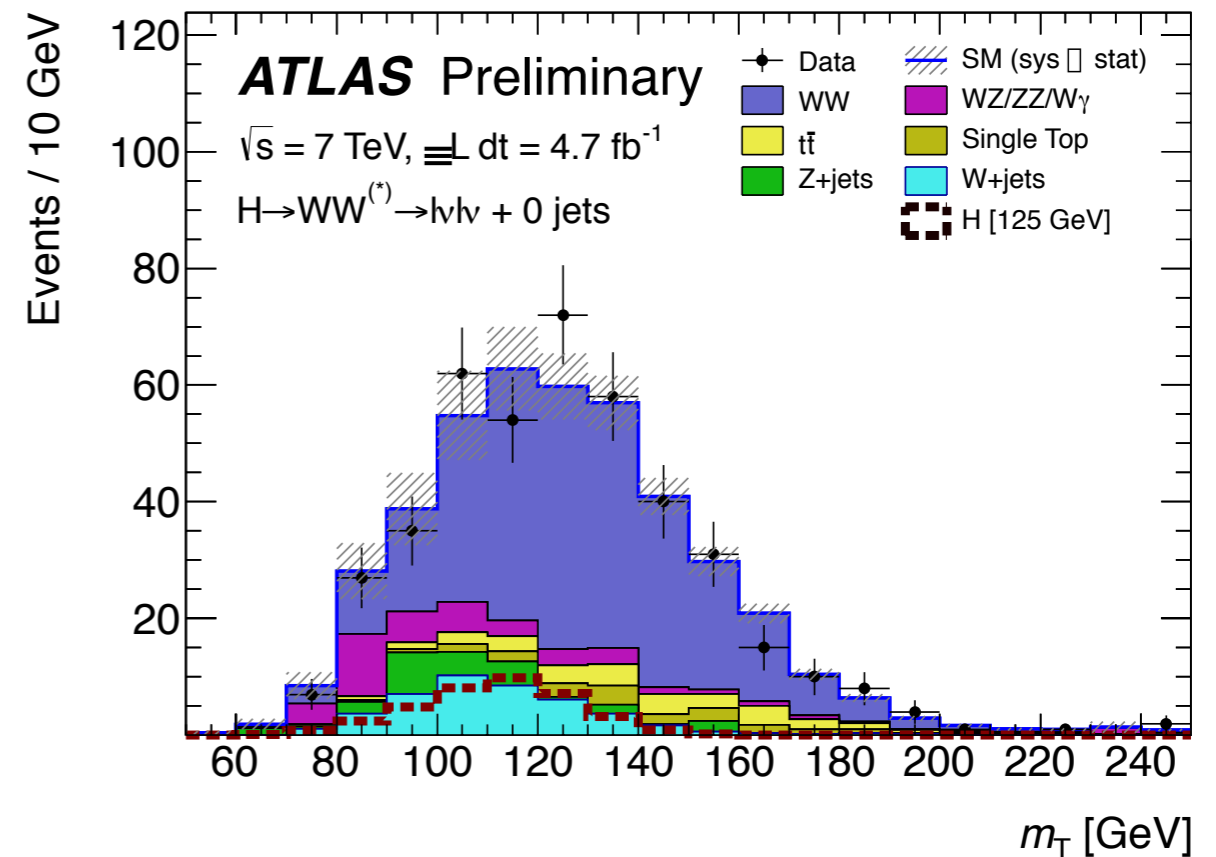
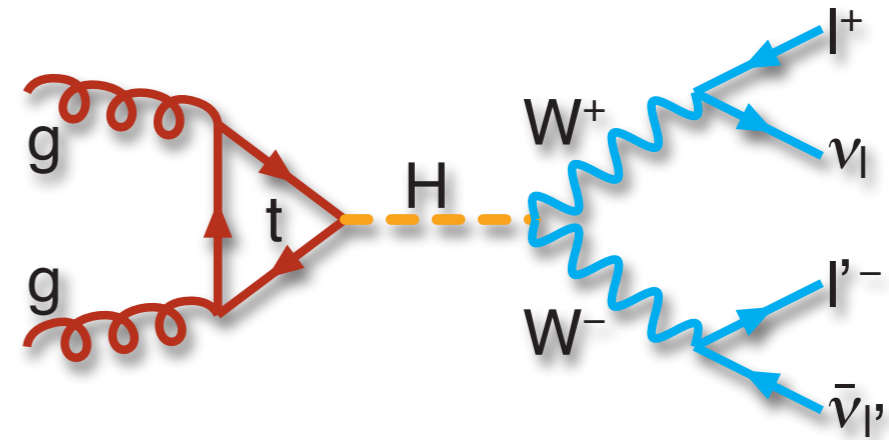


Higgs Boson Searches

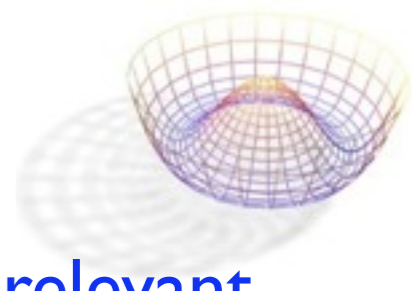


Many possible production and decay modes! Here, focus on channels relevant in the most “interesting” mass range:

- $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$: relatively large event rate but cannot reconstruct mass of event candidates due to escaping neutrinos
- rely on shapes of kinematic variables
- also substantial backgrounds

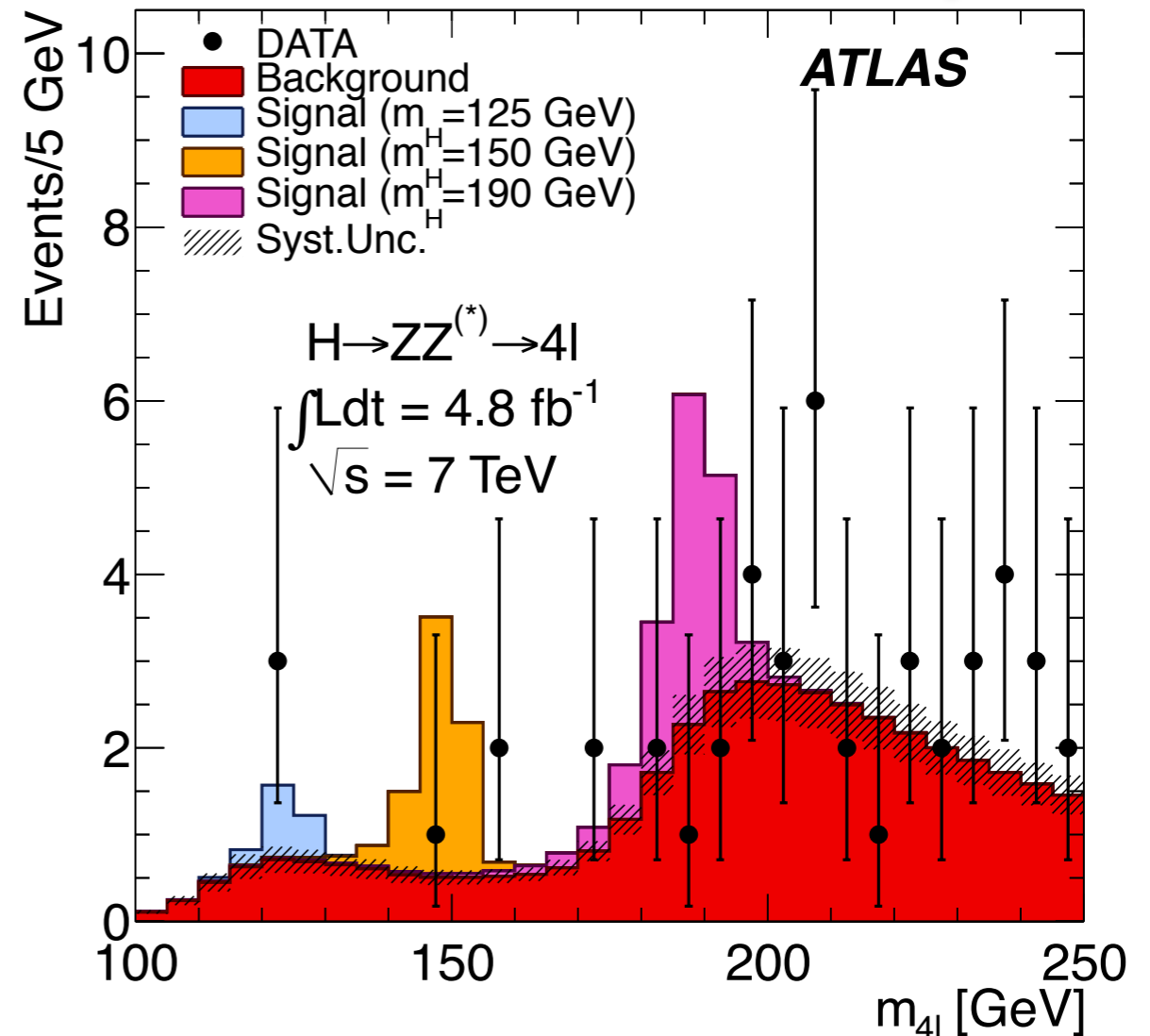
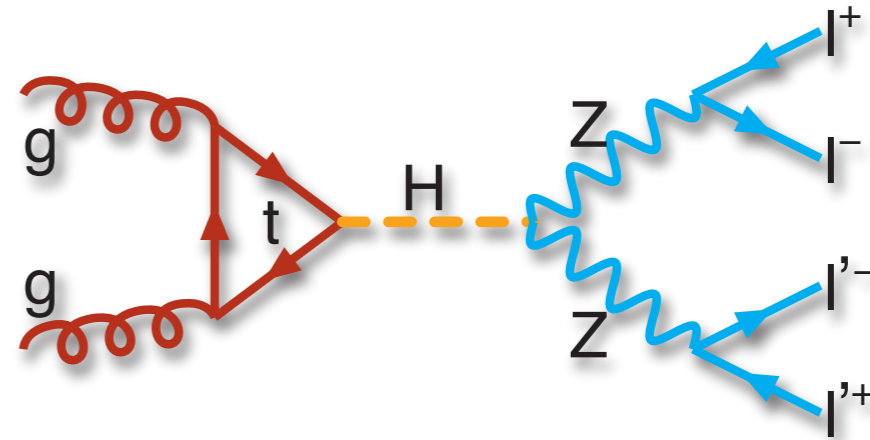


Higgs Boson Searches

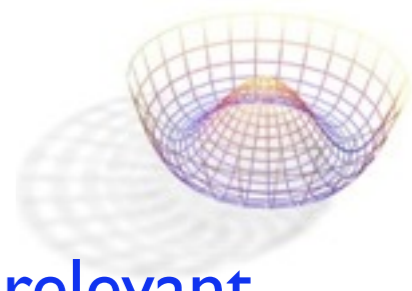


Many possible production and decay modes! Here, focus on channels relevant in the most “interesting” mass range:

- $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$: relatively large event rate but cannot reconstruct mass of event candidates due to escaping neutrinos
 - rely on shapes of kinematic variables
 - also substantial backgrounds
- $H \rightarrow ZZ^{(*)} \rightarrow l^+l^-l'^+l'^-$: precise mass reconstruction, rare but very pure

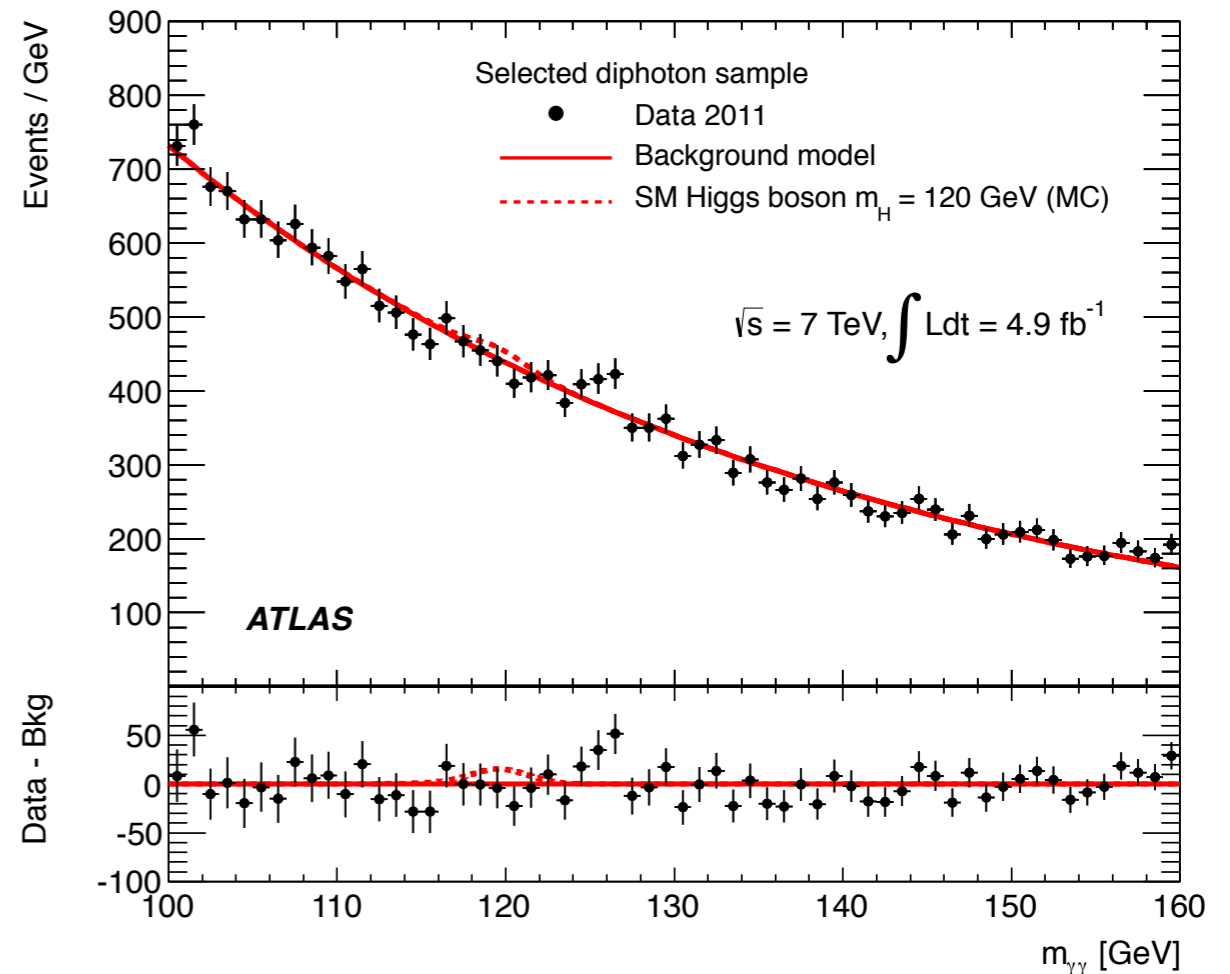
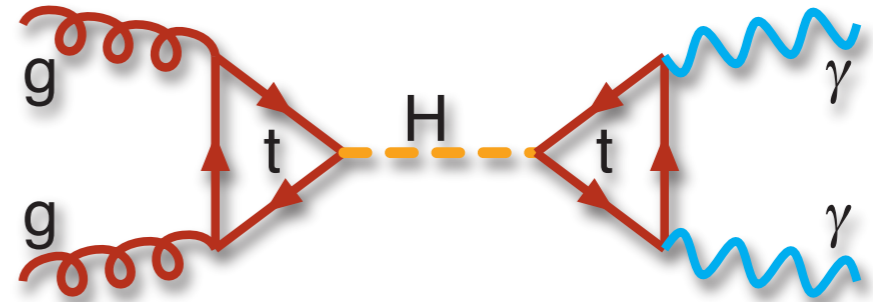


Higgs Boson Searches



Many possible production and decay modes! Here, focus on channels relevant in the most “interesting” mass range:

- $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$: relatively large event rate but cannot reconstruct mass of event candidates due to escaping neutrinos
 - rely on shapes of kinematic variables
 - also substantial backgrounds
- $H \rightarrow ZZ^{(*)} \rightarrow l^+l^- l^+l^-$: precise mass reconstruction, rare but very pure
- $H \rightarrow \gamma\gamma$: precise mass reconstruction, modest rate but large background

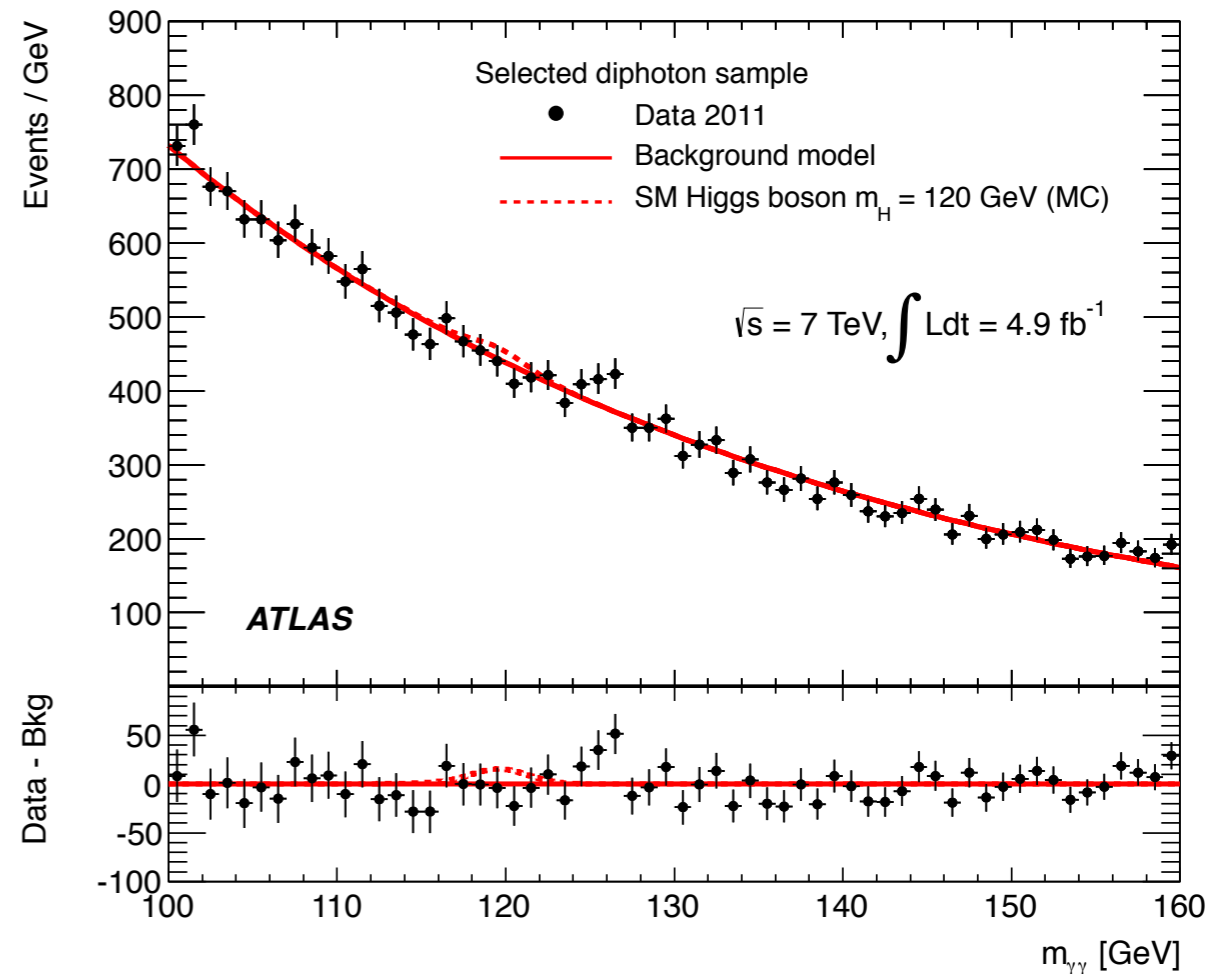
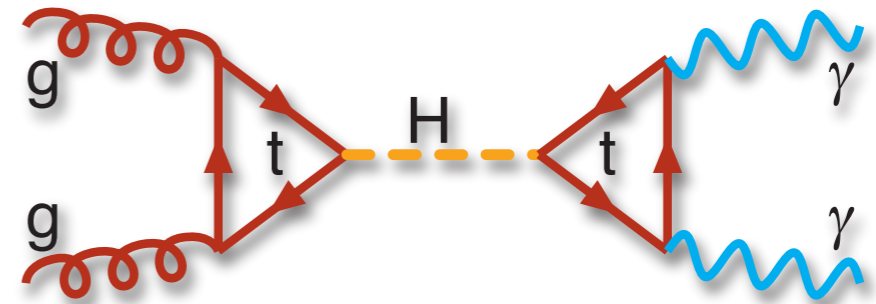


Higgs Boson Searches



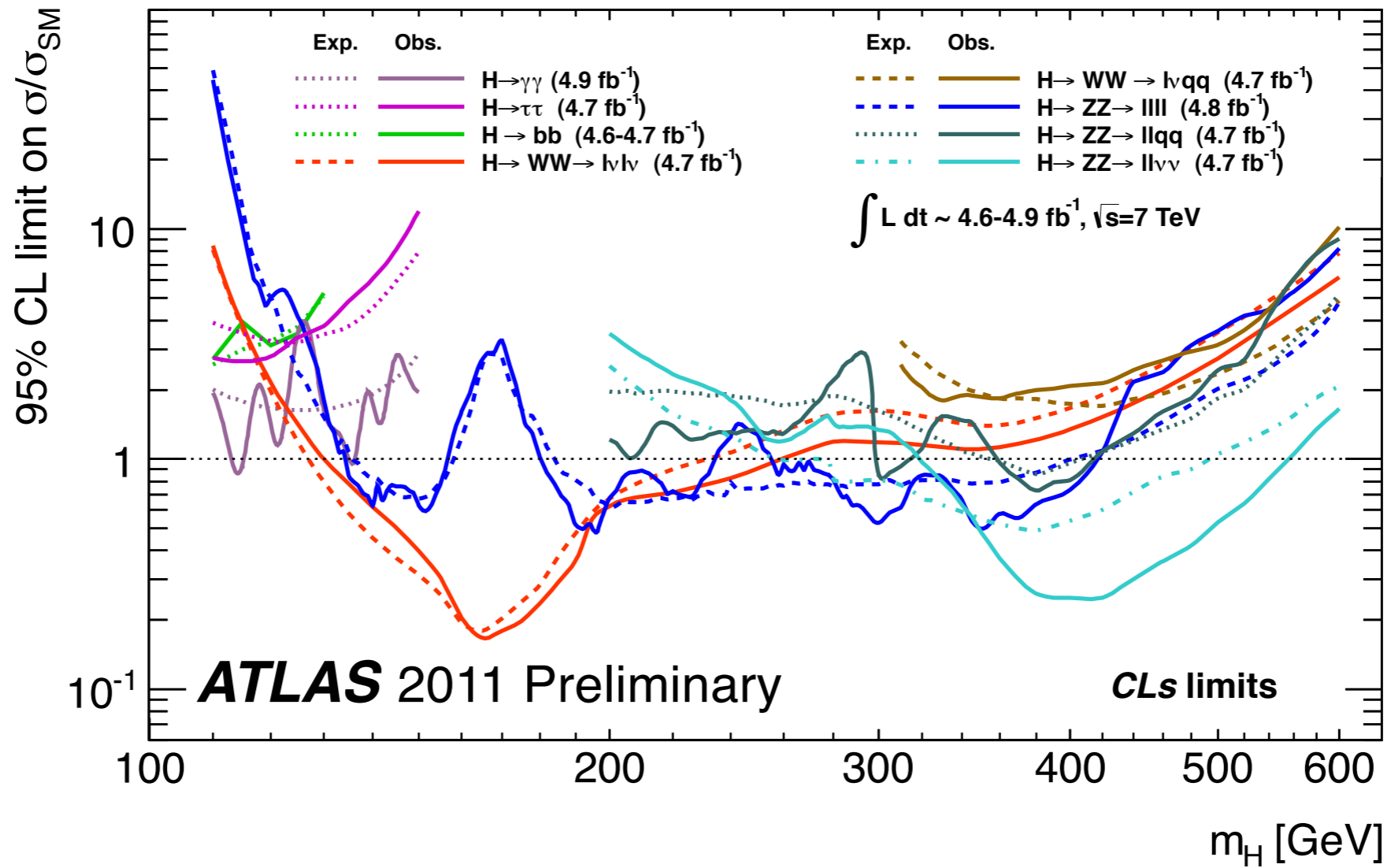
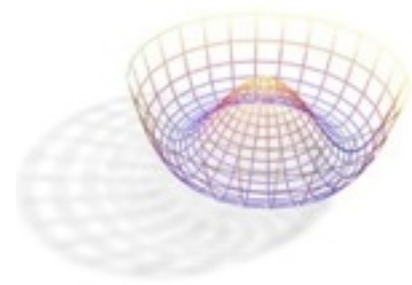
Many possible production and decay modes! Here, focus on channels relevant in the most “interesting” mass range:

- $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$: relatively large event rate but cannot reconstruct mass of event candidates due to escaping neutrinos
 - rely on shapes of kinematic variables
 - also substantial backgrounds
- $H \rightarrow ZZ^{(*)} \rightarrow l^+l^- l^+l^-$: precise mass reconstruction, rare but very pure
- $H \rightarrow \gamma\gamma$: precise mass reconstruction, modest rate but large background

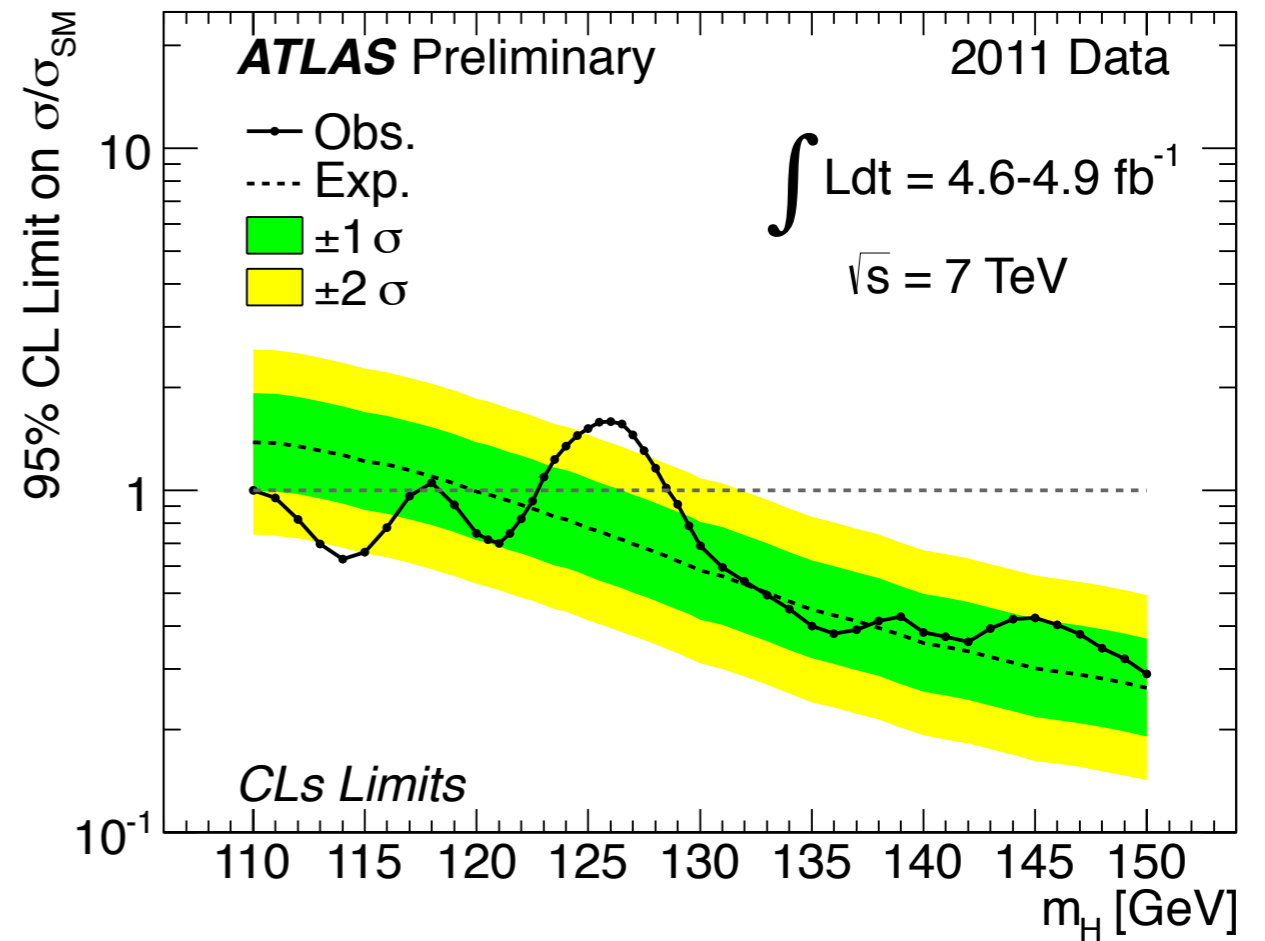
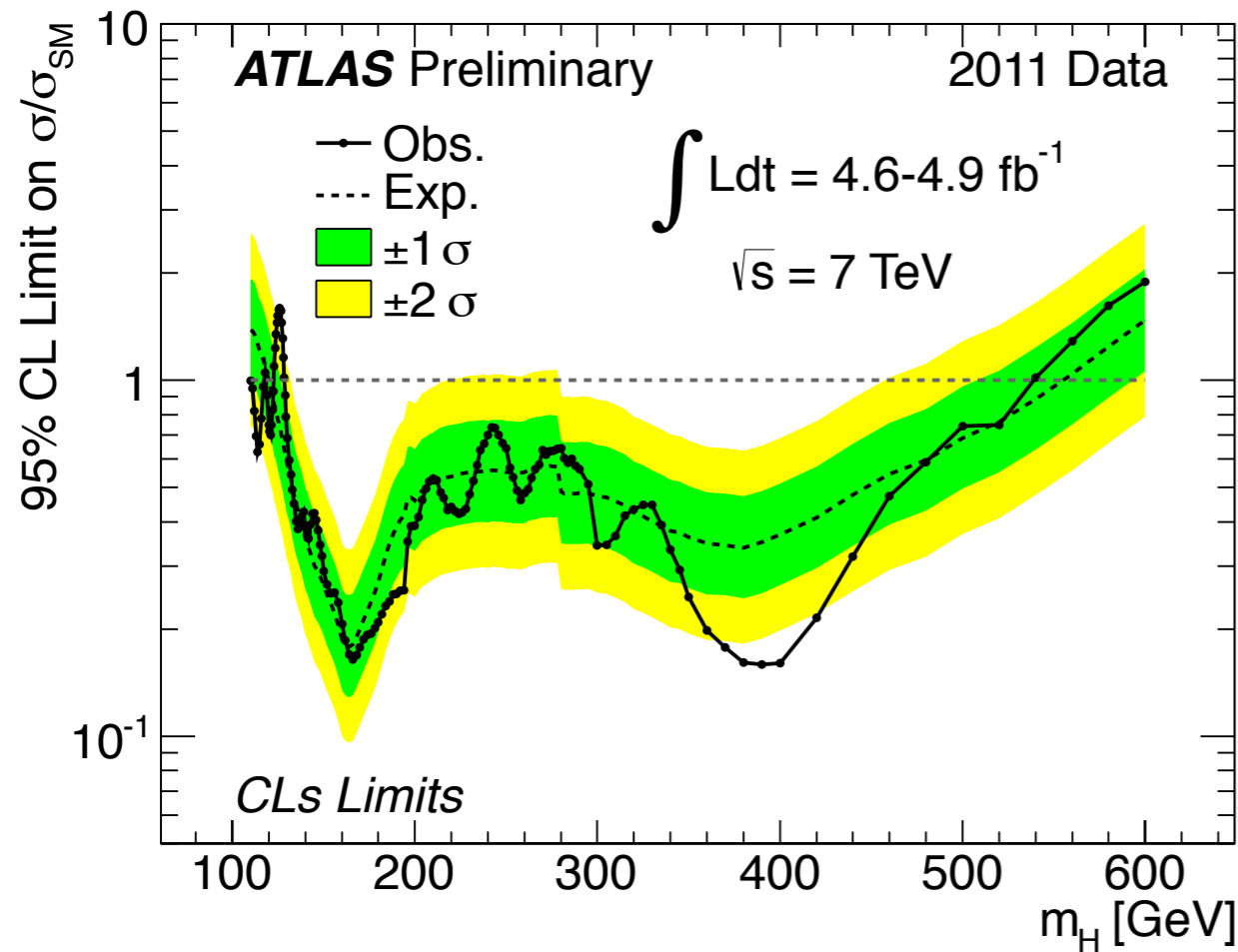
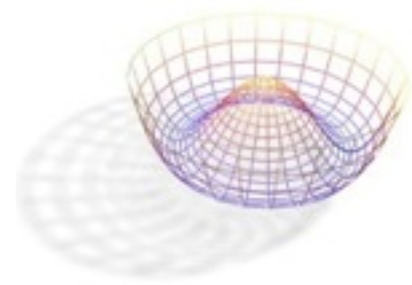


Note: Higgs boson couples to mass but most promising production modes involve massless gluons...

Combining It All



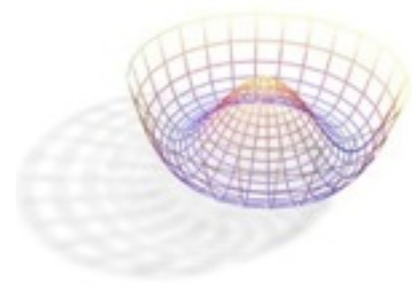
Combining It All



Very similar results by the CMS collaboration!

- exclude $127 \text{ GeV} < M_H < 600 \text{ GeV}$; see excess around 124 GeV (ATLAS excess is around 126 GeV)

Prospects

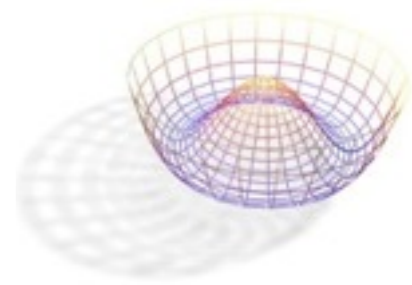


With full 2012 dataset, expect to multiply statistics by factor > 4

- even if we have to cope with more difficult data taking conditions (learning...)

Should be able either to discover a $M_H \sim 125$ GeV Higgs signal, or rule it out altogether!

Prospects



With full 2012 dataset, expect to multiply statistics by factor > 4

- even if we have to cope with more difficult data taking conditions (learning...)

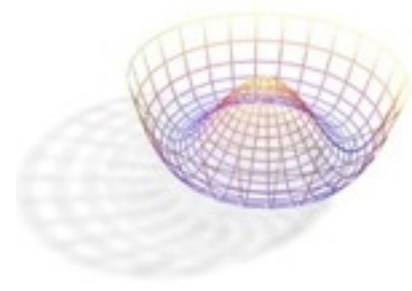
Should be able either to discover a $M_H \sim 125$ GeV Higgs signal, or rule it out altogether!

Exclusion:

- this (essentially) means **ruling out the Standard Model!**
- have to look for other ways to break electroweak symmetry
 - study vector boson scattering at high energy; look for direct signs of new physics

↑
See talk by
P. de Jong

Prospects



With full 2012 dataset, expect to multiply statistics by factor > 4

- even if we have to cope with more difficult data taking conditions (learning...)

Should be able either to discover a $M_H \sim 125$ GeV Higgs signal, or rule it out altogether!

Exclusion:

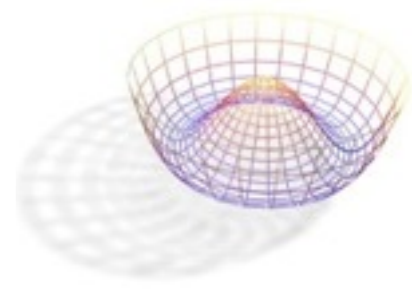
- this (essentially) means **ruling out the Standard Model!**
- have to look for other ways to break electroweak symmetry
 - study vector boson scattering at high energy; look for direct signs of new physics

Discovery:

- need to study its properties: spin & CP quantum numbers, coupling to other particles
- lacking direct signs of supersymmetry, the best way to distinguish the Standard Model from this alternative!

See talk by
P. de Jong

Prospects



With full 2012 dataset, expect to multiply statistics by factor > 4

- even if we have to cope with more difficult data taking conditions (learning...)

Should be able either to discover a $M_H \sim 125$ GeV Higgs signal, or rule it out altogether!

Exclusion:

- this (essentially) means **ruling out the Standard Model!**
- have to look for other ways to break electroweak symmetry
 - study vector boson scattering at high energy; look for direct signs of new physics

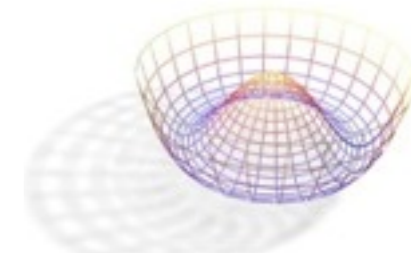
Discovery:

- need to study its properties: spin & CP quantum numbers, coupling to other particles
- lacking direct signs of supersymmetry, the best way to distinguish the Standard Model from this alternative!

See talk by
P. de Jong

Particle physics has an exciting time ahead!

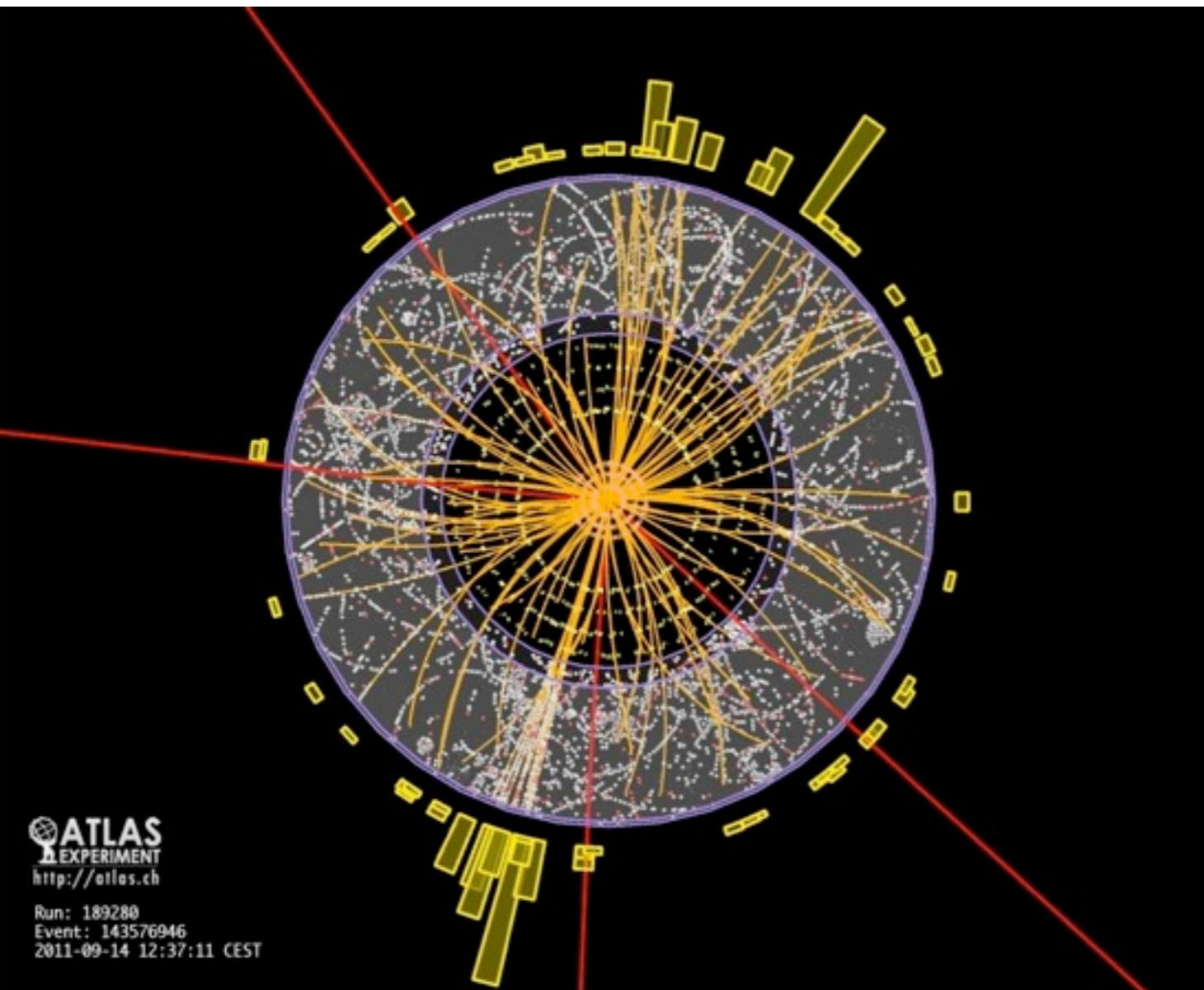
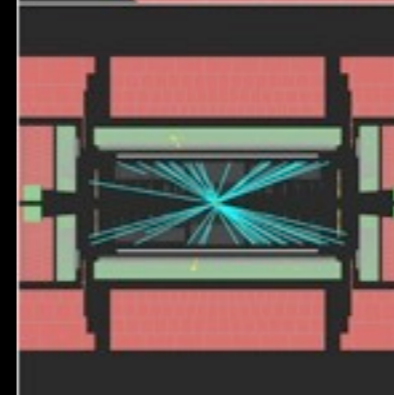
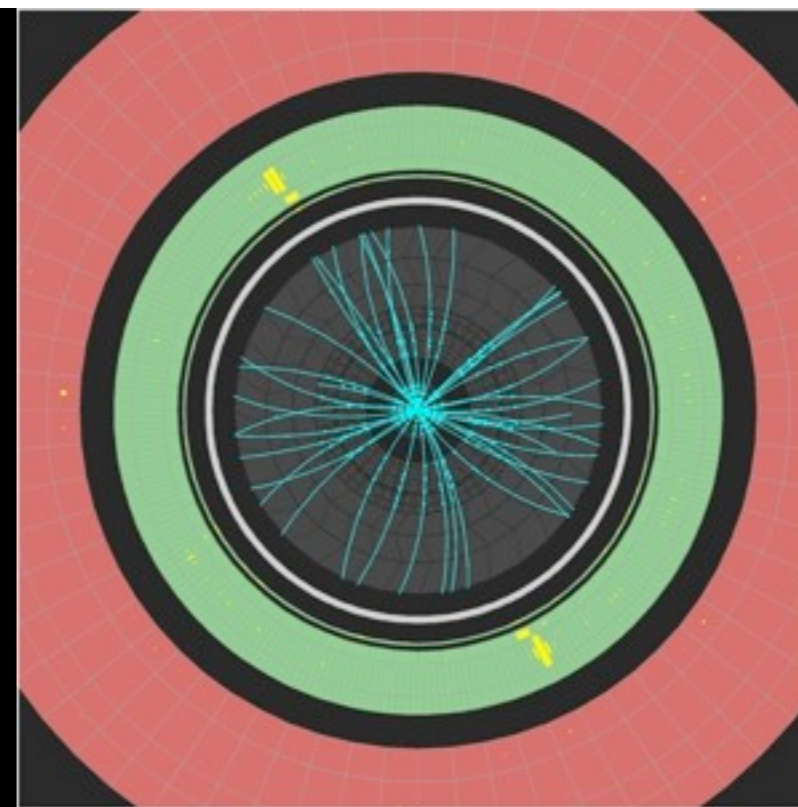
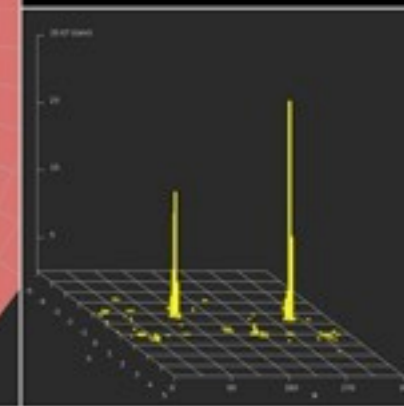
Thank you!



ATLAS
EXPERIMENT

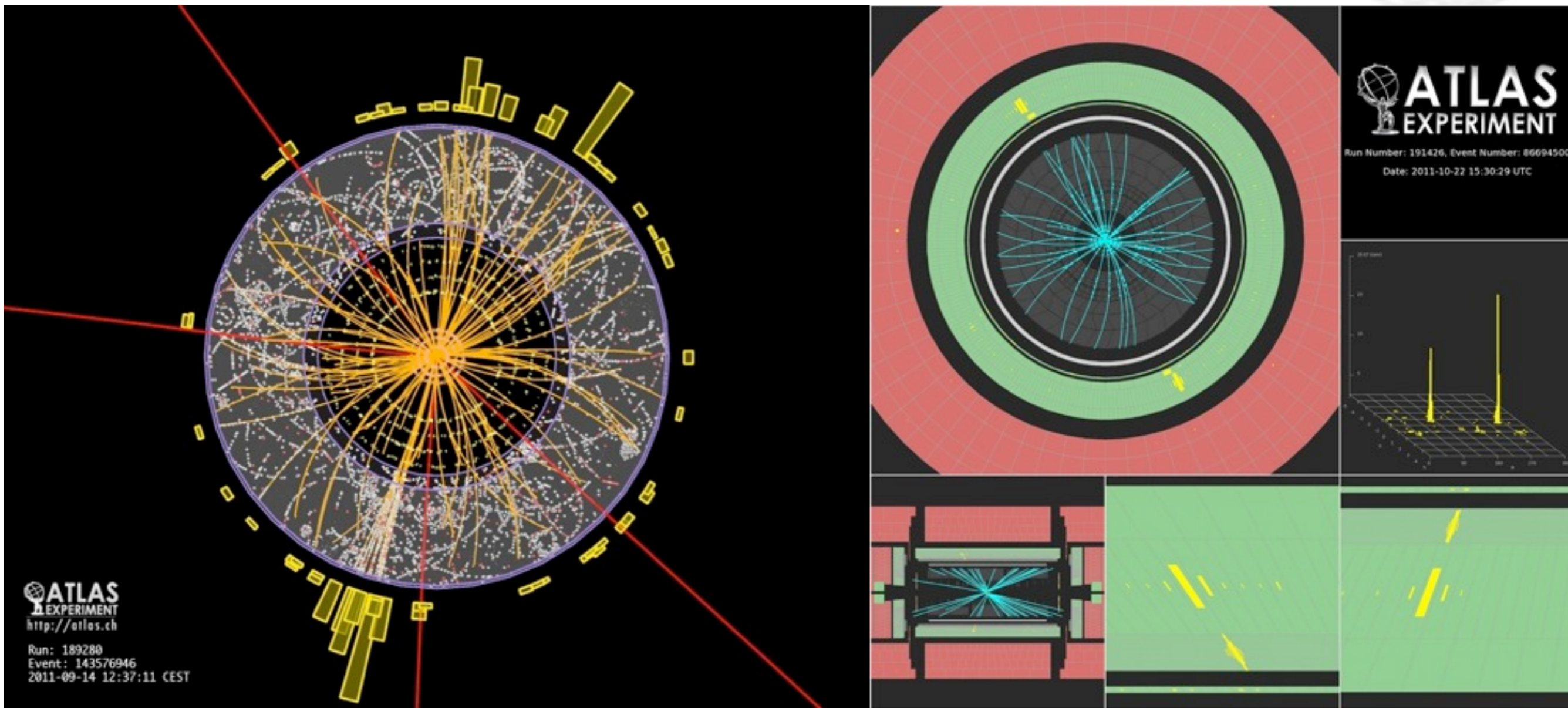
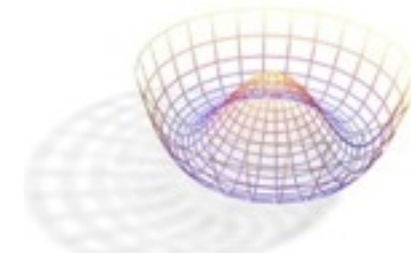
Run Number: 191426, Event Number: 86694500

Date: 2011-10-22 15:30:29 UTC



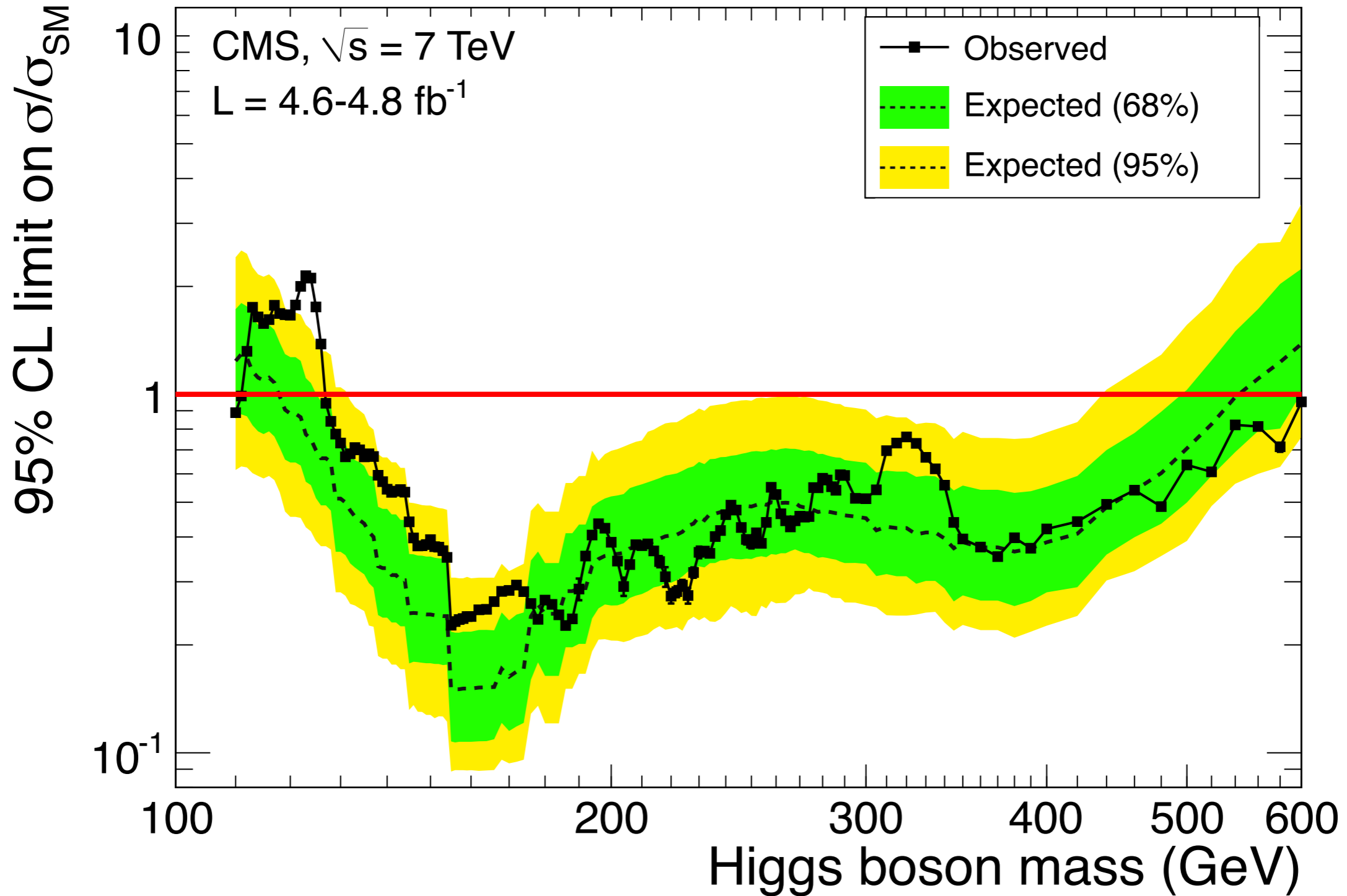
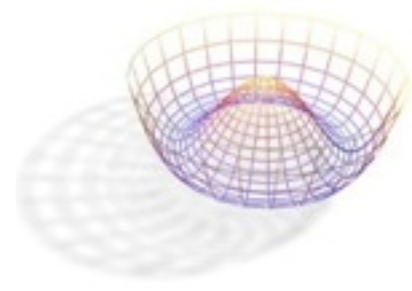
ATLAS $H \rightarrow ZZ(*) \rightarrow \mu^+\mu^-\mu^+\mu^-$ and $H \rightarrow \gamma\gamma$ candidates

Thank you!

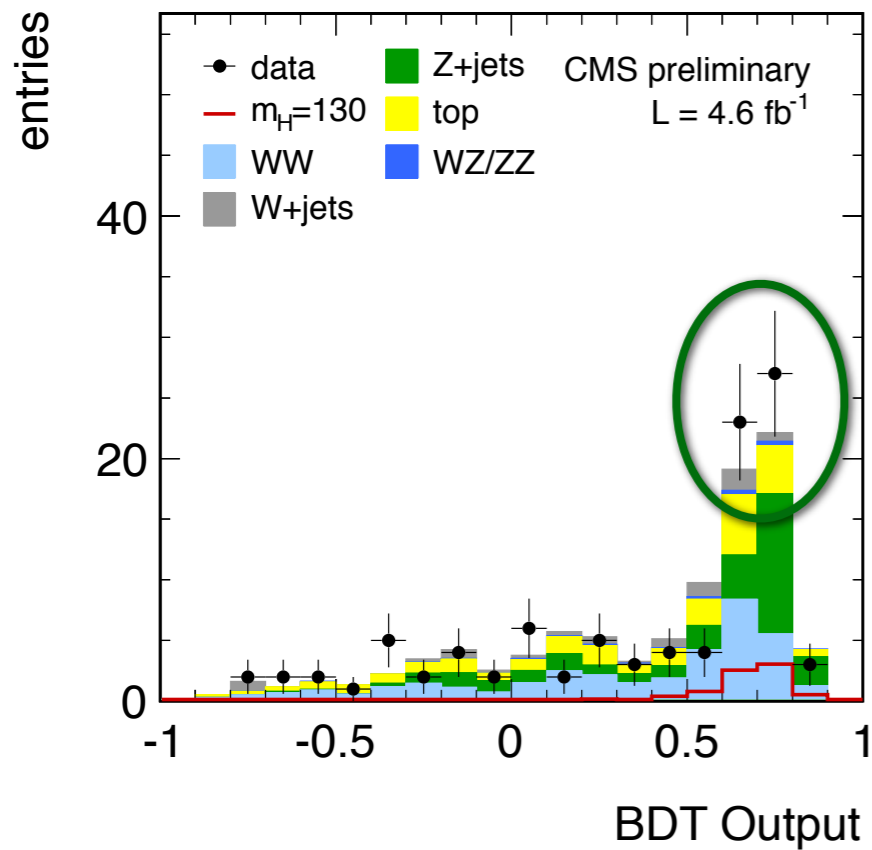
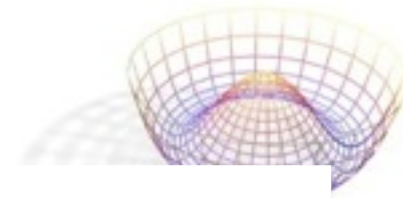


ATLAS $H \rightarrow ZZ(*) \rightarrow \mu^+\mu^-\mu^+\mu^-$ and $H \rightarrow \gamma\gamma$ candidates

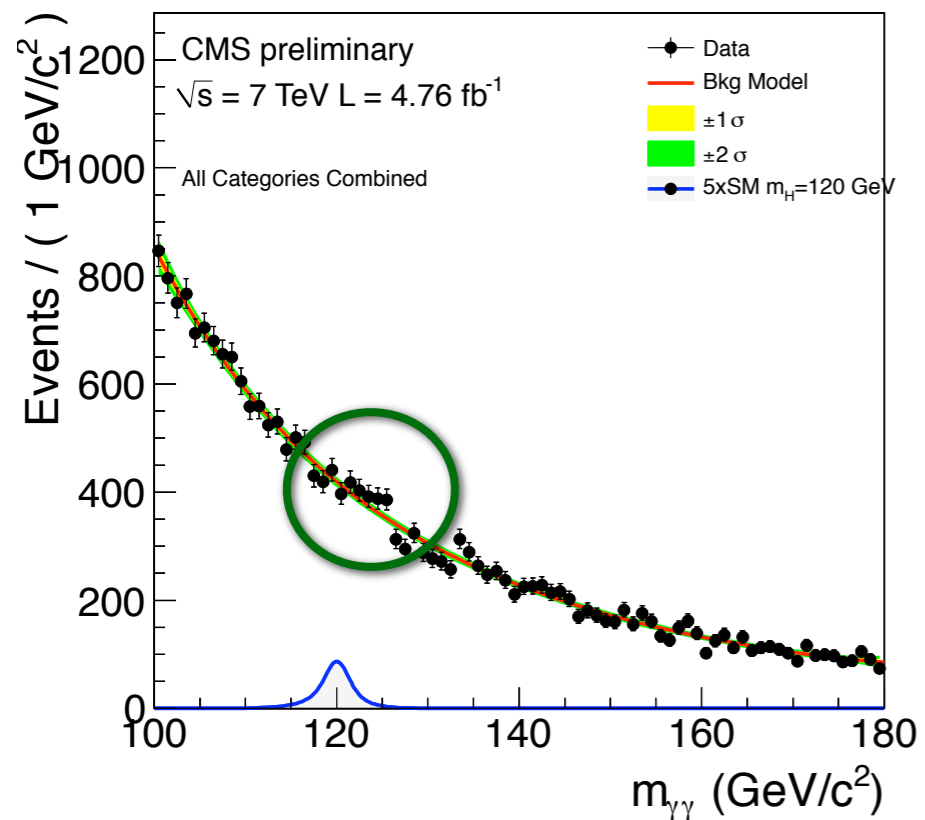
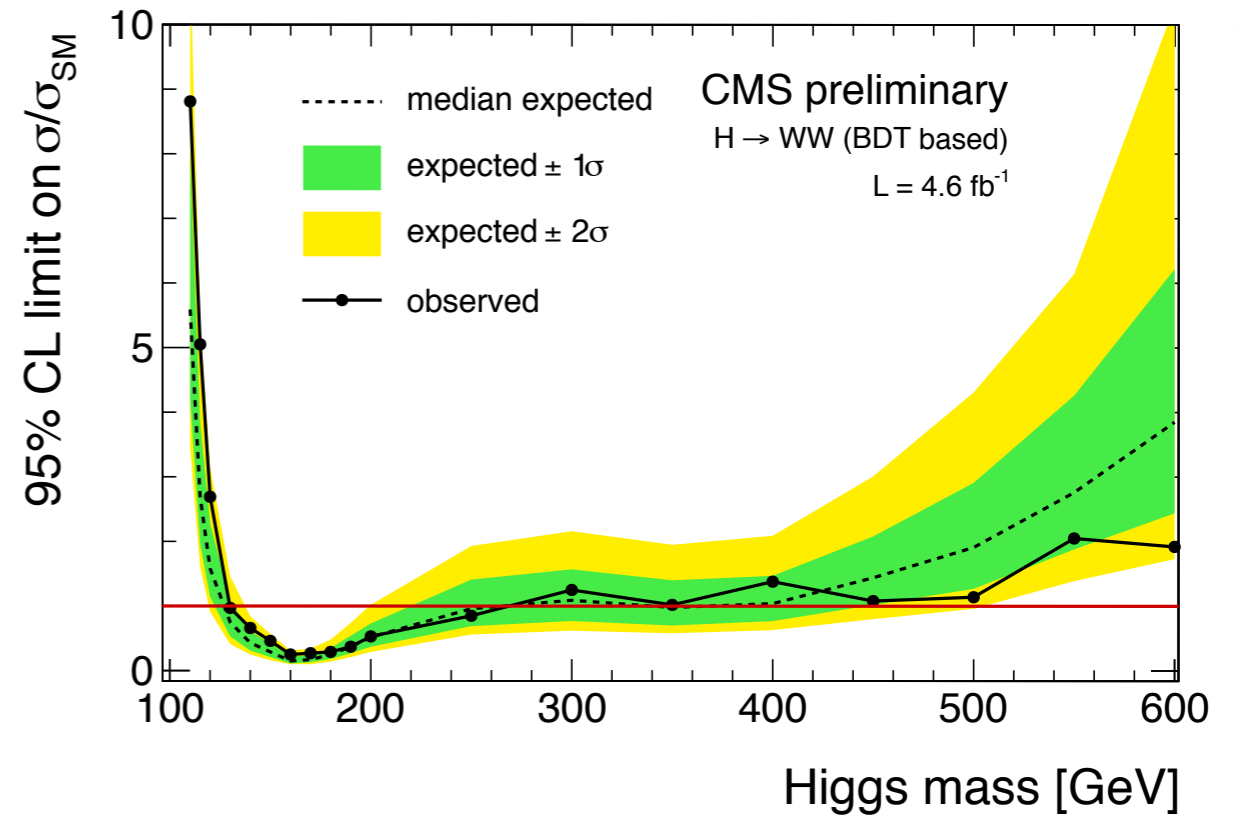
Results by the CMS Experiment



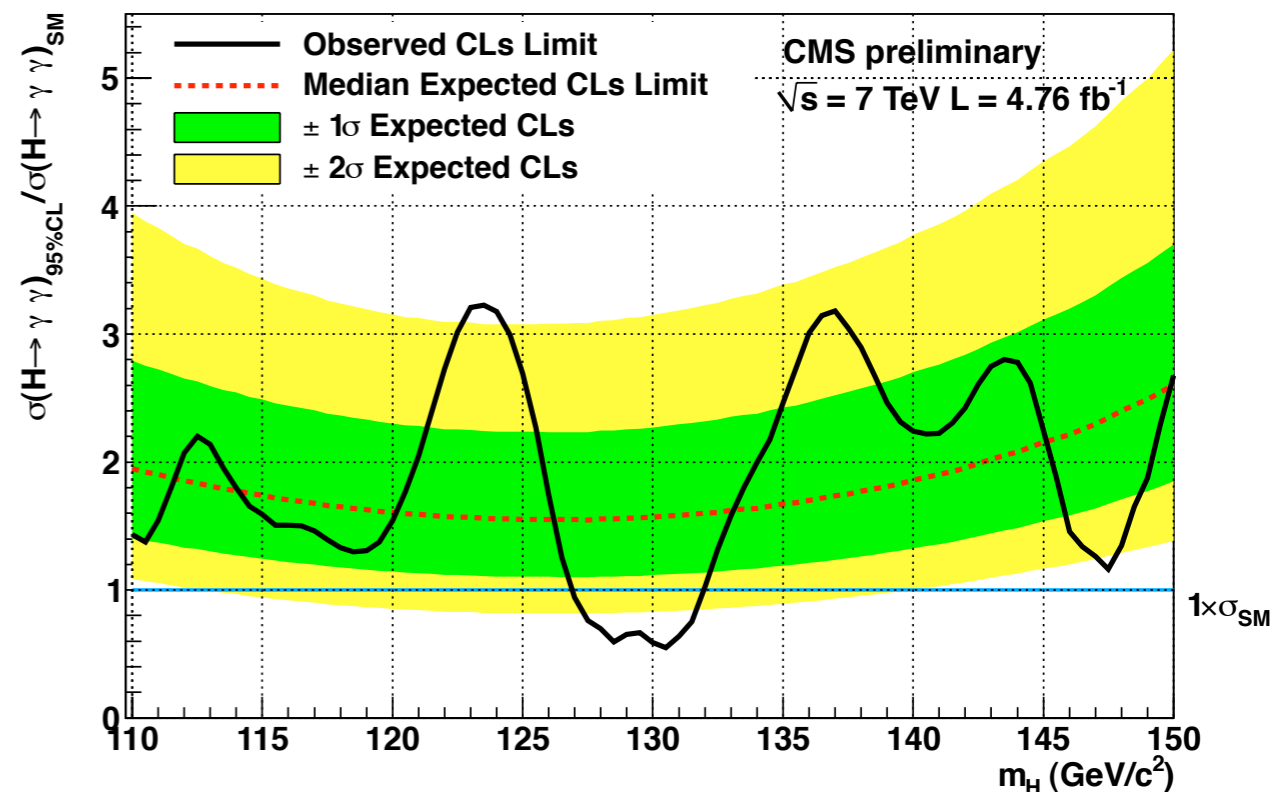
Results by the CMS Experiment



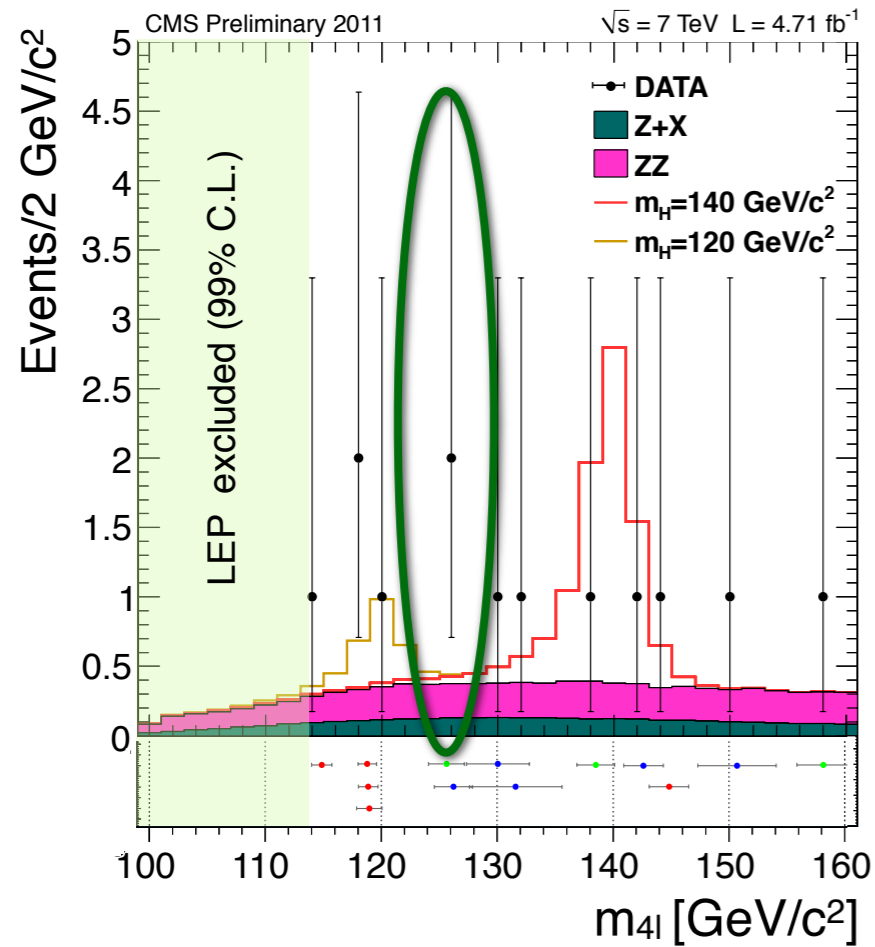
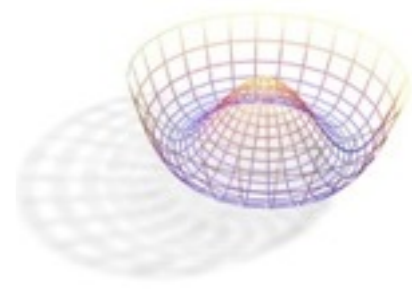
$H \rightarrow W^+W^-$



$H \rightarrow \gamma\gamma$



Results by the CMS Experiment

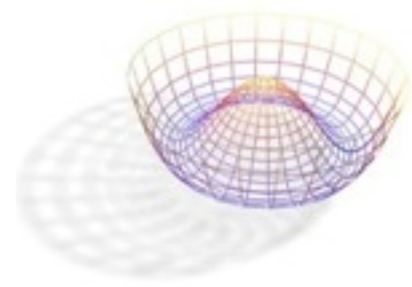


$H \rightarrow ZZ$

Found 2 candidate events near 126 GeV

- 1 in $e^+e^-e^+e^-$
- 1 in $e^+e^-\mu^+\mu^-$

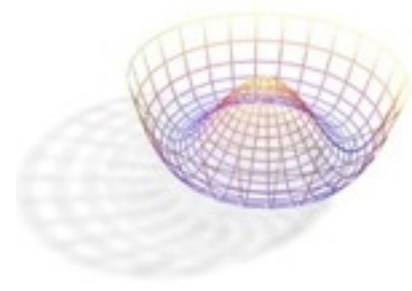
Comparison



ATLAS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (124 GeV)
- **excess** in $\gamma\gamma$ final state (126 GeV)

Comparison



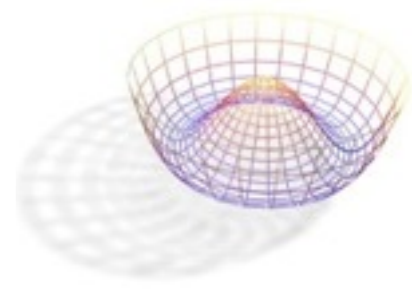
ATLAS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (124 GeV)
- **excess** in $\gamma\gamma$ final state (126 GeV)

CMS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (126 GeV)
- **excess** in $\gamma\gamma$ final state (123 GeV)

Comparison



ATLAS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (124 GeV)
- **excess** in $\gamma\gamma$ final state (126 GeV)

CMS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (126 GeV)
- **excess** in $\gamma\gamma$ final state (123 GeV)

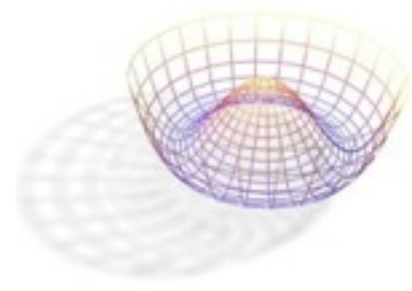
Caveat emptor!

- each individual excess not statistically significant
- masses in $\gamma\gamma$, ZZ are close **but do not match** \Rightarrow questions:
 - are the energy calibrations as well understood as we think?
 - is this just a statistical fluctuation after all?

Time (and additional investigation) will tell

But one way or the other, we expect to make a much more definite statement within a year

Comparison



ATLAS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (124 GeV)
- **excess** in $\gamma\gamma$ final state (126 GeV)

CMS:

- **excess** in W^+W^- final states: broad but compatible with low-mass Higgs boson
- **excess** in ZZ final state (126 GeV)
- **excess** in $\gamma\gamma$ final state (123 GeV)

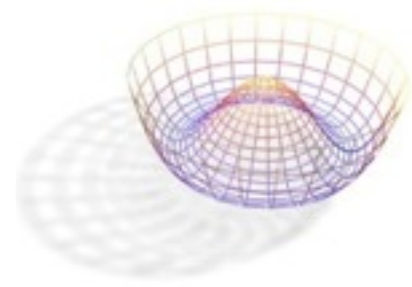
Caveat emptor!

- each individual excess not statistically significant
- masses in $\gamma\gamma$, ZZ are close **but do not match** \Rightarrow questions:
 - are the energy calibrations as well understood as we think?
 - is this just a statistical fluctuation after all?

Time (and additional investigation) will tell

Either we find the Higgs particle or we rule out the Standard Model!

Finally...

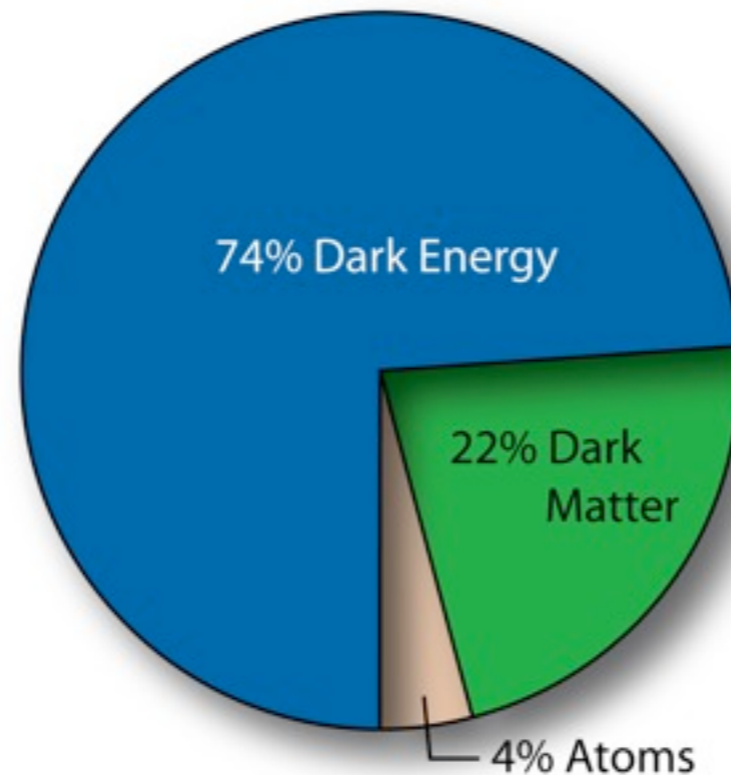


Finding the Higgs boson does not mean particle physics is finished!

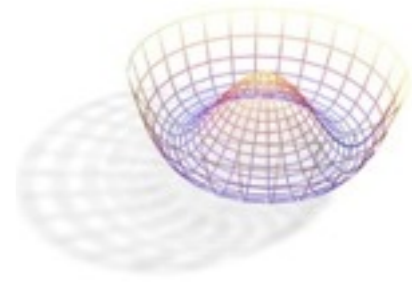
The Standard Model cannot incorporate **gravity** in a consistent way

The Higgs boson's mass is not stable against **radiative corrections**

The Standard Model does not explain **Dark Matter / Dark Energy**

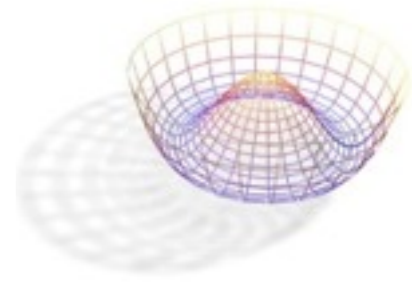


Outlook

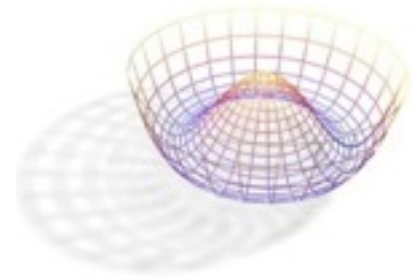


The ghost you're trying to reach is currently unavailable.
Please leave a message after the beep.

Outlook

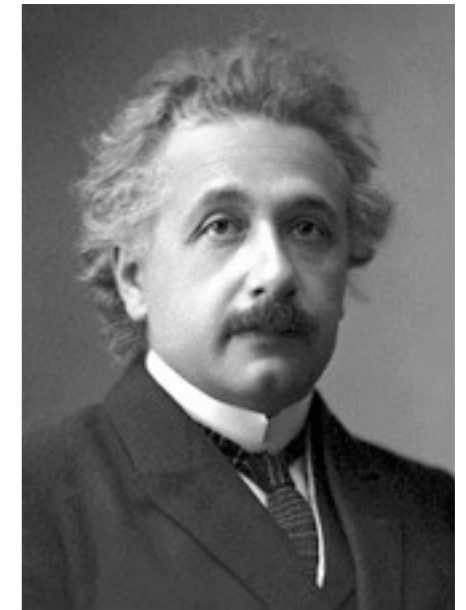
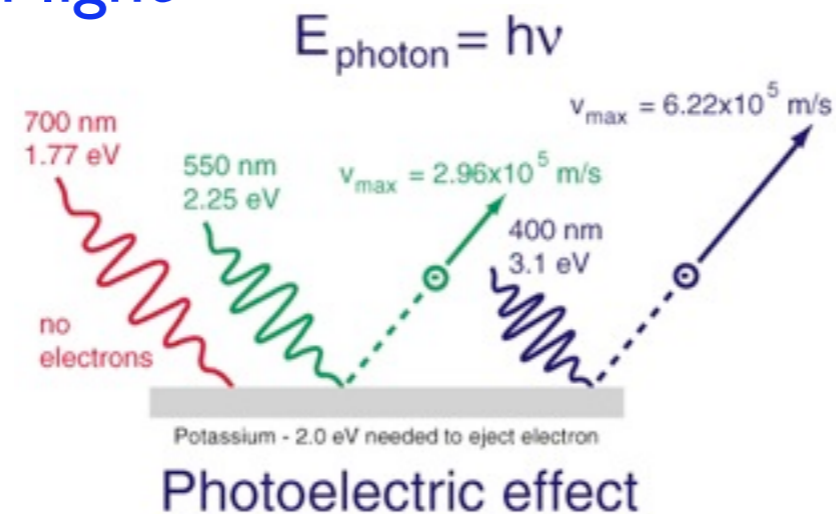


Quantum Electrodynamics

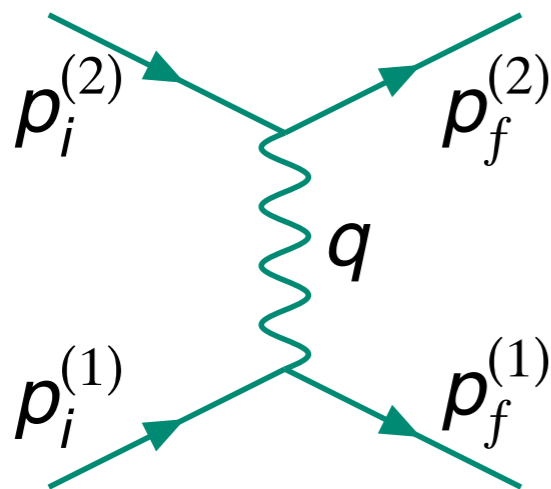


Einstein (1905): photo-electric effect

➔ particle nature of light



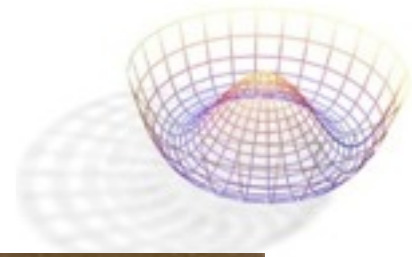
Paradigm change! Electromagnetic interaction described as photon exchange



Graphical representation:
Feynman diagrams
(intuitive way to compute outcome of scattering processes in QM)



Going beyond the naked eye



Antoni van Leeuwenhoek, 1632-1723:

- invention of the microscope
- discovery first bacteria (“kleine beestjes”), 0.5 - 500 μm

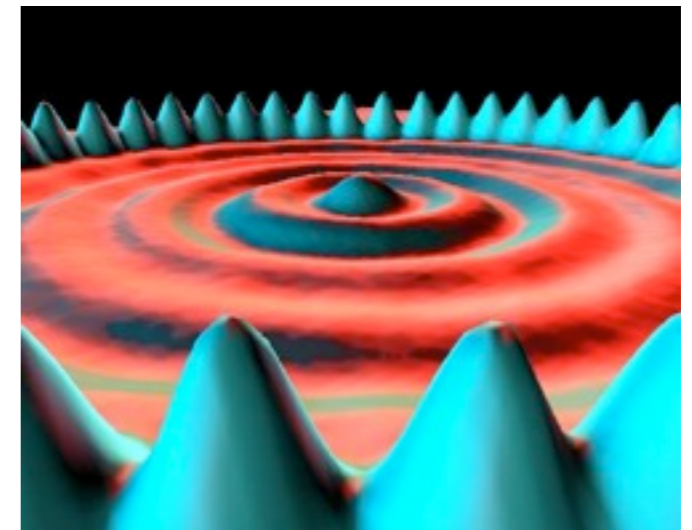


E. coli (size $\sim 1 \mu\text{m}$)

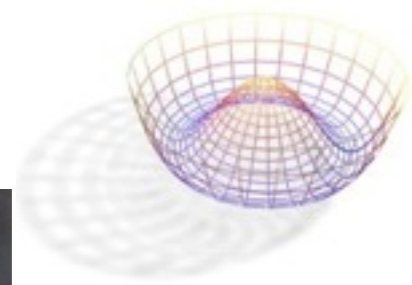


Minimum discernible dimensions $\sim \lambda$

- limit when using visible light: 0.5 μm
- improvement to $\sim 1 \text{ \AA}$ possible using **STM, AFM**



The atom “cracked”



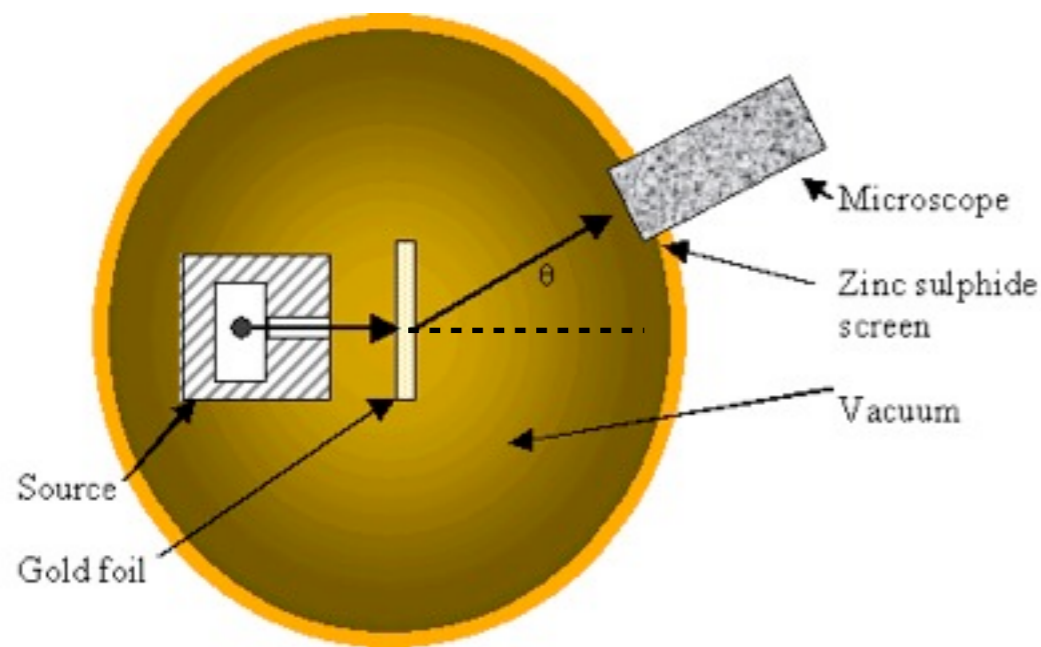
Idea: use particles to “see” smaller structures

- Rutherford: scattering of α -particles (${}^4\text{He}$ nuclei, $E_\alpha \approx 3 \text{ MeV}$) off a gold foil
- Quantum mechanical translation: de Broglie wavelength $\lambda \sim h/p$



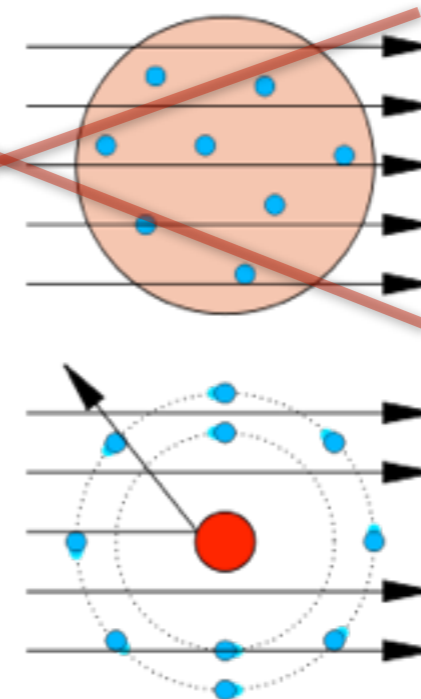
Planck's constant

projectile momentum



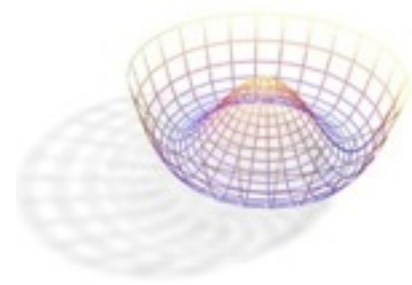
~~diffuse charge distribution (Thomson)~~

charge distribution with nucleus (Rutherford)



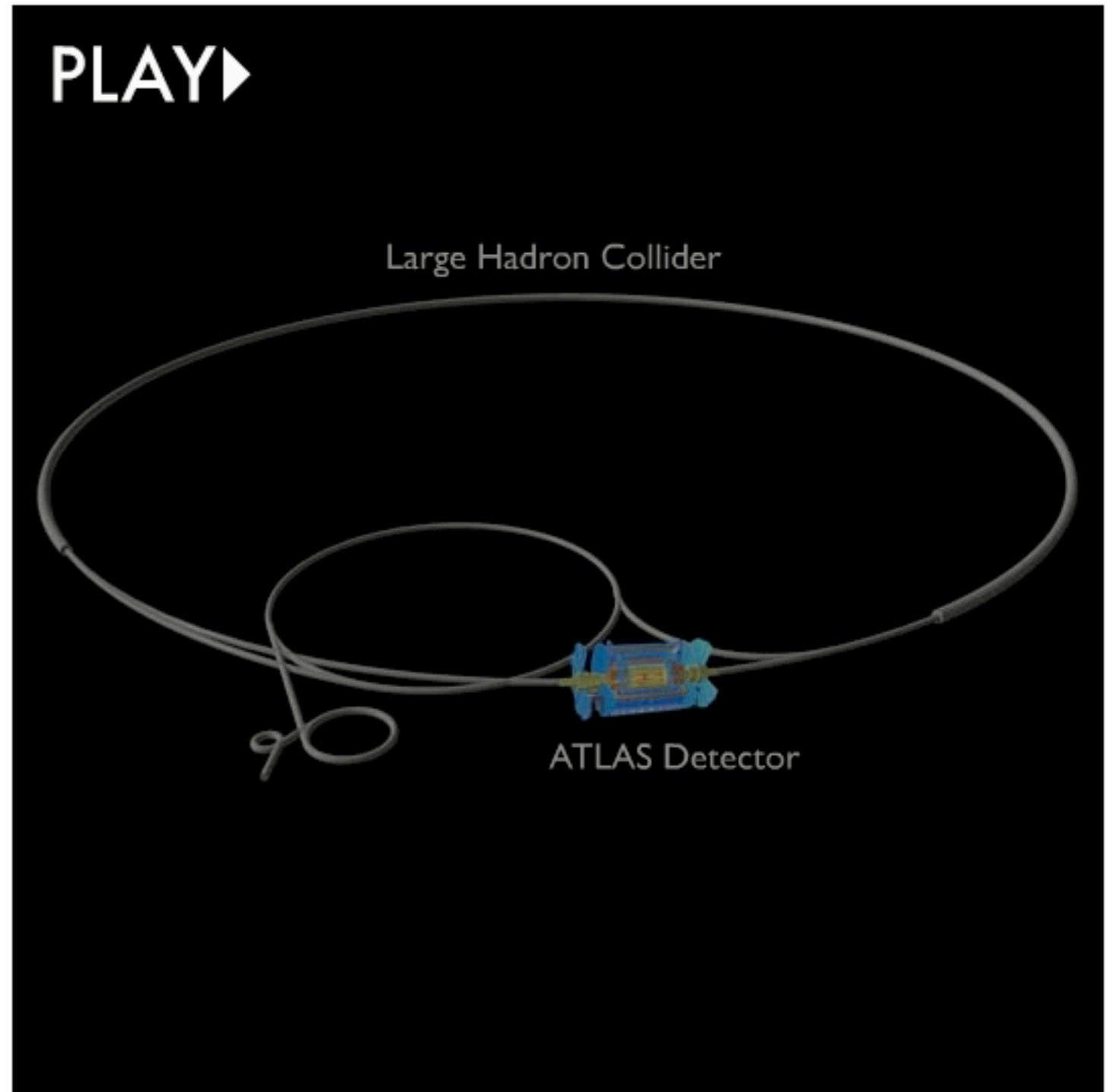
Repeated in the 60's with scattering of 180 MeV electrons on protons
➡ the proton ($r \sim 1 \text{ fm}$) contains further sub-structure (quarks)!

State of the art

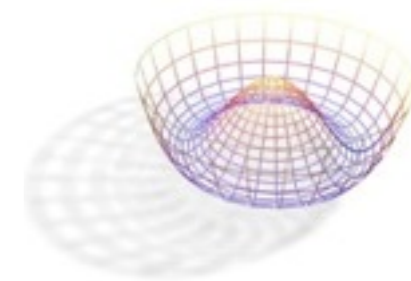


Present scheme in CERN's Large Hadron Collider:

- accelerate proton beams to energies of 3.5 TeV per proton
 - $v/c \approx 0.999999996$ (energy to be doubled in 2014: 6→8)
 - in both directions!
- make them collide in the centres of the detectors
- experiments analyze outcome of collisions and select “interesting” events
- stochastic process, no control over outcome of individual collision
 - ▮ can only select after the fact

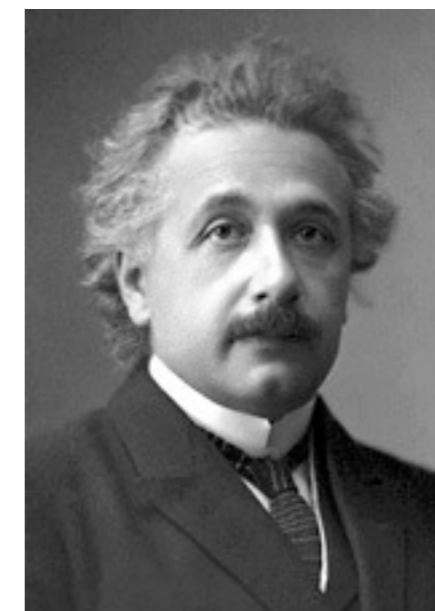
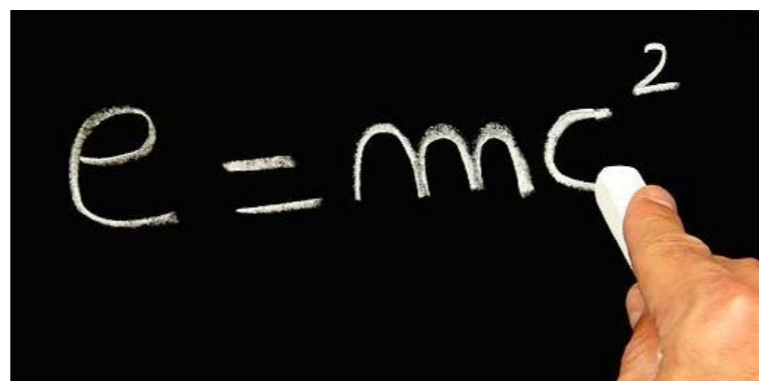


Techniques



High energy allows for the creation of other, usually short-lived particles

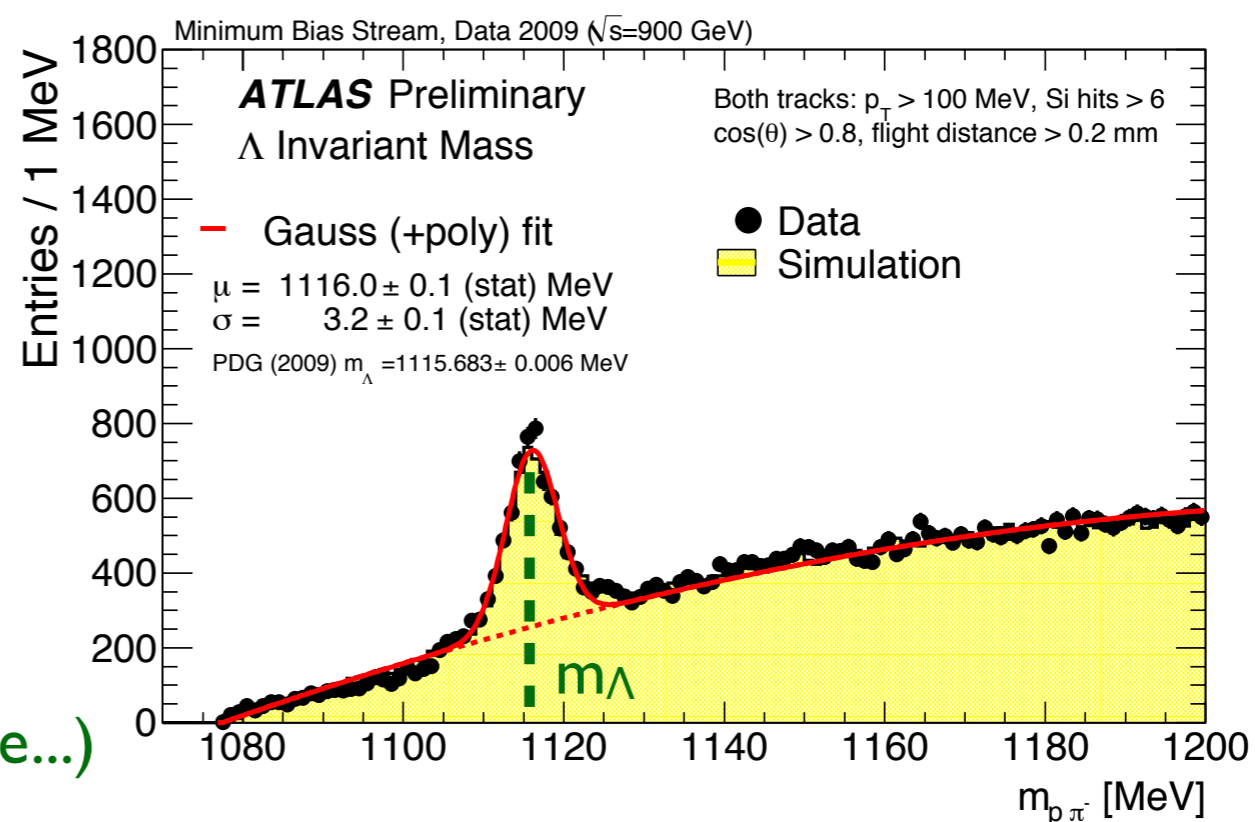
- $\tau < 10^{-22}$ s for “interesting” particles



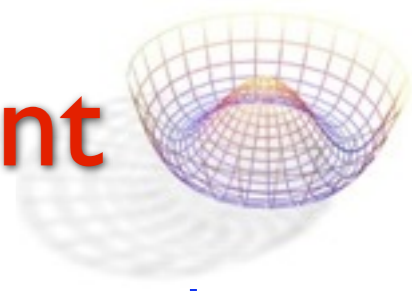
- **in collisions:** convert kinetic energy into mass

- **in decay processes:** reconstruct mass of the decaying particle (if all decay products are measured)

(not the Higgs particle...)



The Electron's Magnetic Dipole Moment



Well-known system: interaction of magnetic dipole moments with external magnetic field

$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

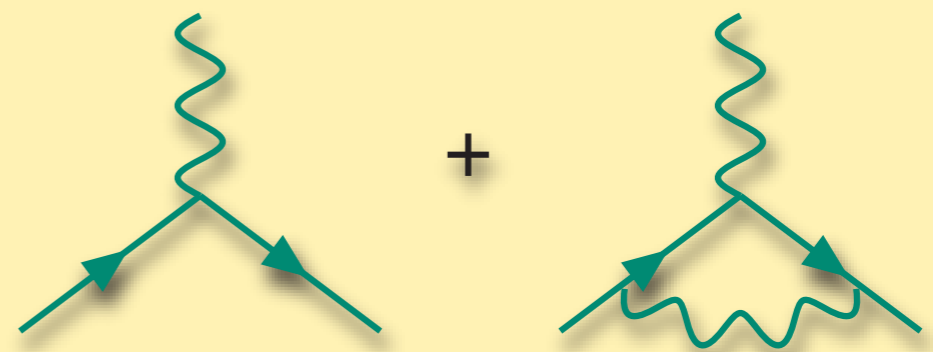
- **Zeeman splitting** of (atomic) energy levels
- Spin precession around B-field axis, **Larmor frequency** $\omega = \gamma B$

Unlike regular QM, QED provides a **prediction** for g !

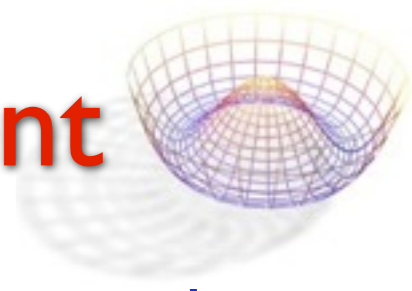
- Applying the gauge principle to the **Dirac equation** (relativistic equation of motion for spin-1/2 particles): $g=2$
- Computing quantum corrections: expansion in powers (up to fifth power) of fine structure constant

$$\alpha \equiv \frac{e^2}{4\pi}$$

contributions at lowest orders



The Electron's Magnetic Dipole Moment



Well-known system: interaction of magnetic dipole moments with external magnetic field

$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

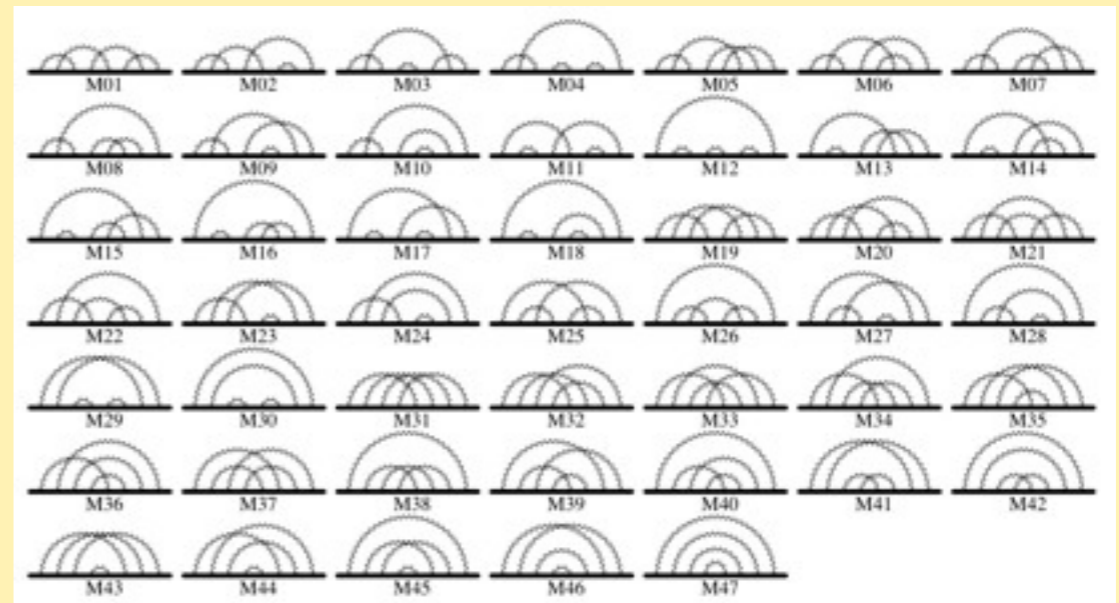
- Zeeman splitting of (atomic) energy levels
- Spin precession around B-field axis, Larmor frequency $\omega = \gamma B$

Unlike regular QM, QED provides a prediction for g !

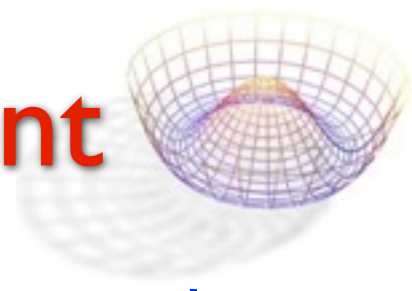
- Applying the gauge principle to the Dirac equation (relativistic equation of motion for spin-1/2 particles): $g=2$
- Computing quantum corrections: expansion in powers (up to fifth power) of fine structure constant

$$\alpha \equiv \frac{e^2}{4\pi}$$

subset of contributions at 5th order



The Electron's Magnetic Dipole Moment



Well-known system: interaction of magnetic dipole moments with external magnetic field

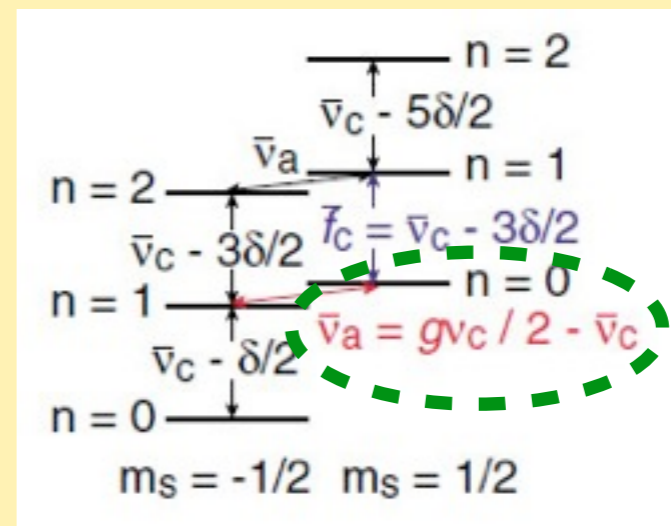
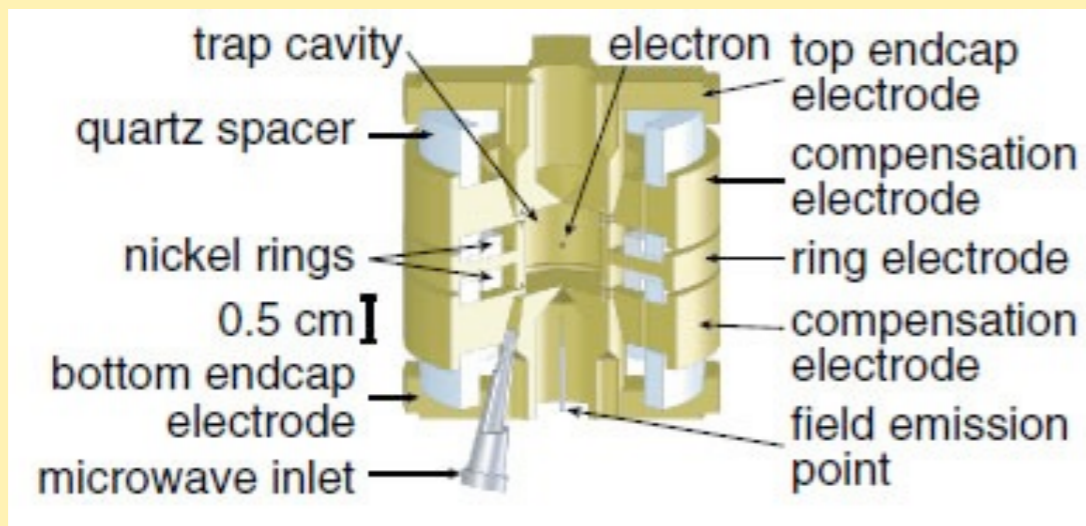
$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

- Zeeman splitting of (atomic) energy levels
- Spin precession around B-field axis, Larmor frequency $\omega = \gamma B$

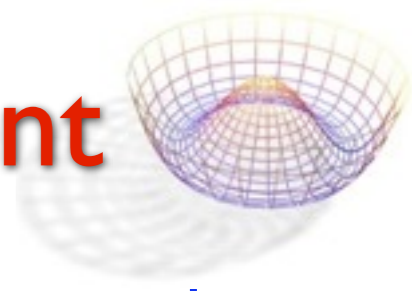
Unlike regular QM, QED provides a prediction for g !

G. Gabrielse et al., 2008
Cylindrical Penning trap

Observe a single electron for months



The Electron's Magnetic Dipole Moment



Well-known system: interaction of magnetic dipole moments with external magnetic field

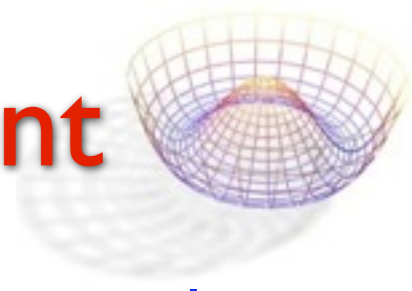
$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

- Zeeman splitting of (atomic) energy levels
- Spin precession around B-field axis, Larmor frequency $\omega = \gamma B$

Unlike regular QM, QED provides a prediction for g !

The comparison:

The Electron's Magnetic Dipole Moment



Well-known system: interaction of magnetic dipole moments with external magnetic field

$$H = -\vec{\mu} \cdot \vec{B}, \quad \vec{\mu} = \gamma \vec{S} \equiv g \left(\frac{q}{2m} \right) \vec{S}$$

- Zeeman splitting of (atomic) energy levels
- Spin precession around B-field axis, Larmor frequency $\omega = \gamma B$

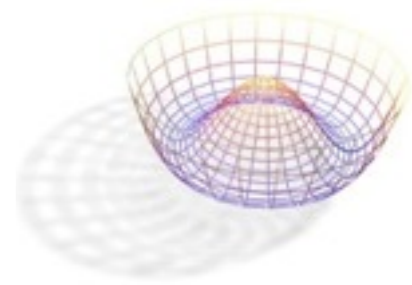
Unlike regular QM, QED provides a prediction for g !

The comparison:

$$g/2 = \begin{cases} 1.001\,159\,652\,180\,73(28) & \text{(experiment)} \\ 1.001\,159\,652\,180\,85(76) & \text{(theory)} \end{cases}$$

A triumph for QED!

A Colourful Interaction



Three quarks forming baryons (and quark-antiquark pairs forming mesons):
 a new symmetry (and interaction), colour

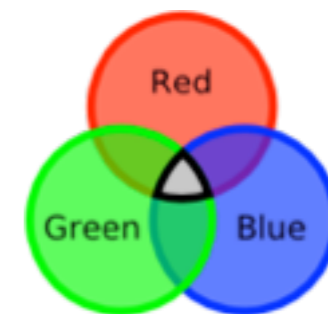
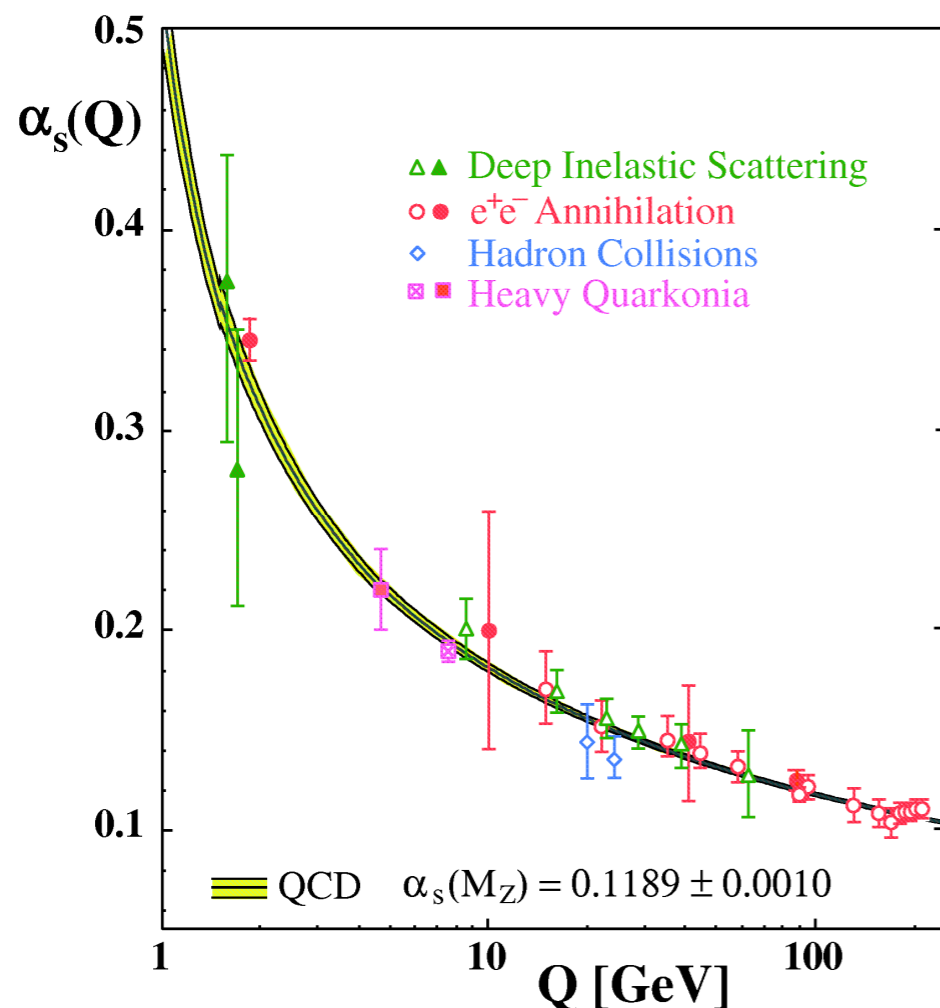
- “gauge principle” interaction with gluons:

$$q \rightarrow \begin{pmatrix} q_r \\ q_g \\ q_b \end{pmatrix}$$

- quarks change identity (colour) under exchange of a gluon!

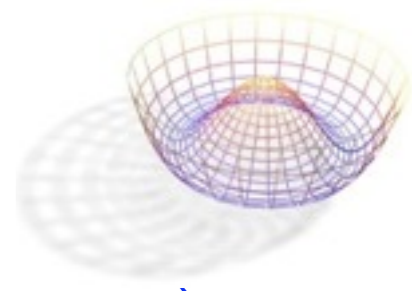
$$p \rightarrow p + \frac{g_s}{2} \sum_a T_a G_a$$

$$\alpha_s = \frac{g_s^2}{4\pi}$$



Quark confinement at low energy

A Colourful Interaction



Three quarks forming baryons (and quark-antiquark pairs forming mesons):
 a new symmetry (and interaction), colour

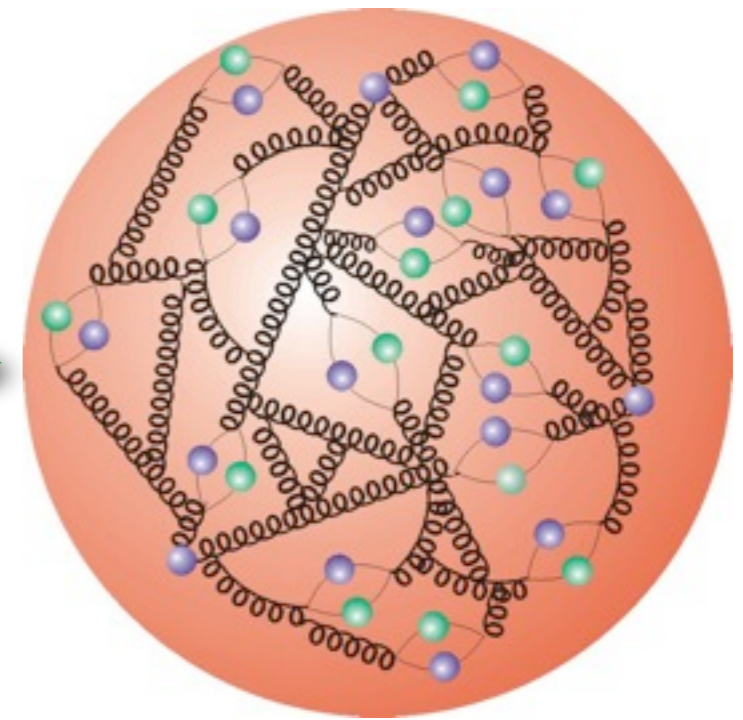
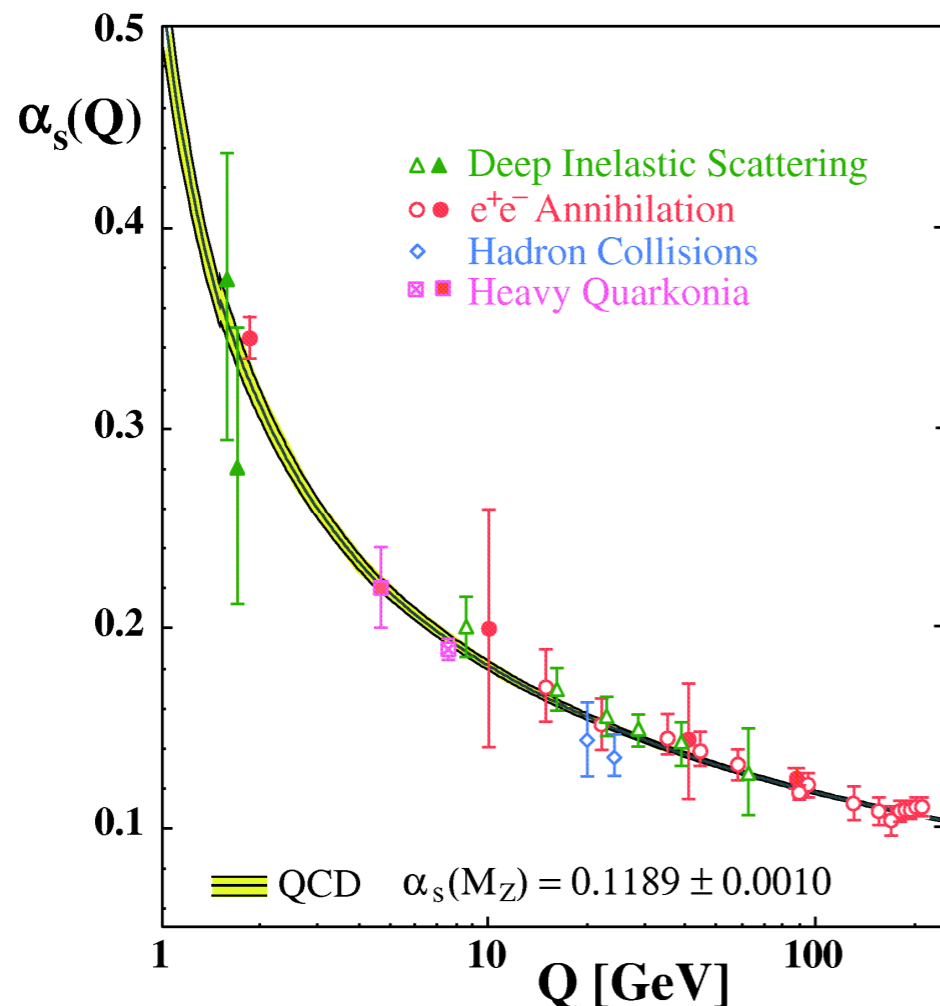
- “gauge principle” interaction with gluons:

$$q \rightarrow \begin{pmatrix} q_r \\ q_g \\ q_b \end{pmatrix}$$

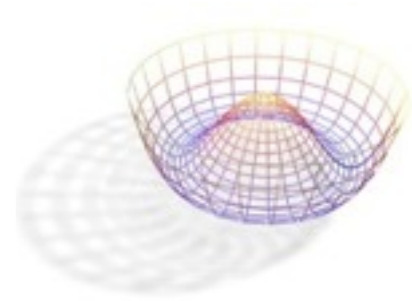
- quarks change identity (colour) under exchange of a gluon!

$$p \rightarrow p + \frac{g_s}{2} \sum_a T_a G_a$$

$$\alpha_s = \frac{g_s^2}{4\pi}$$



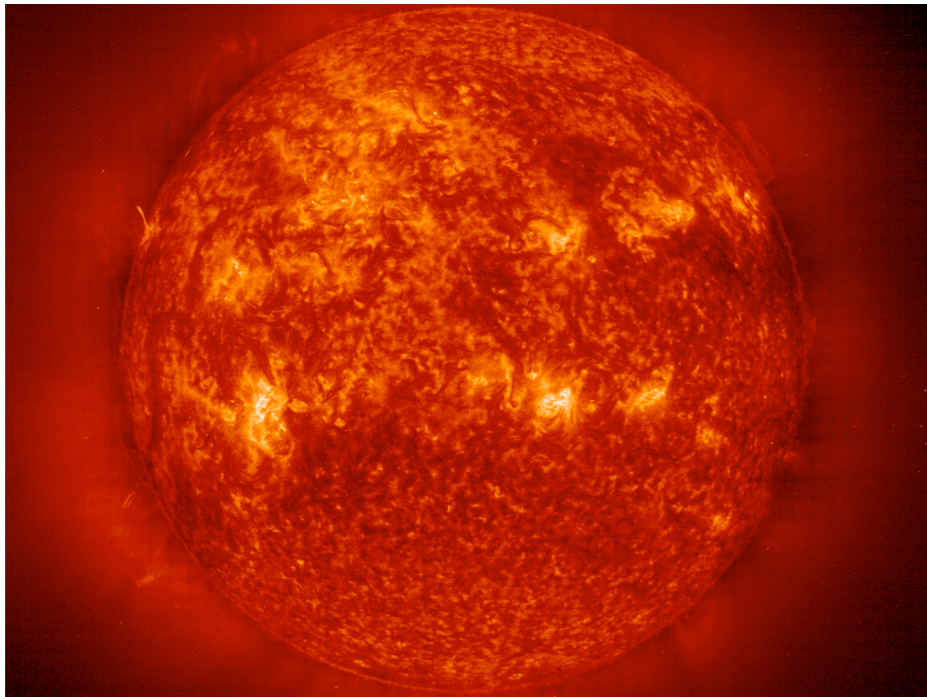
Quark confinement at low energy



The Weak Interaction

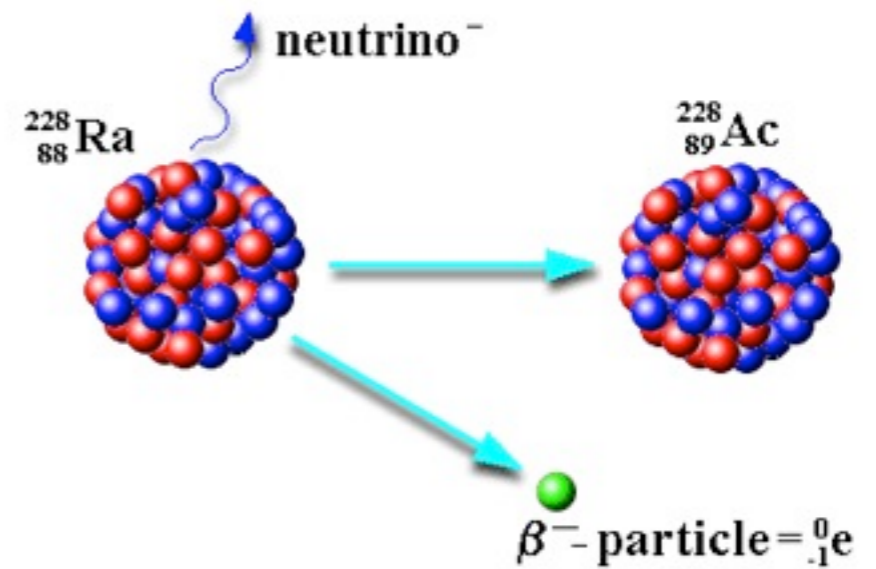
responsible for all nucleonic transmutations

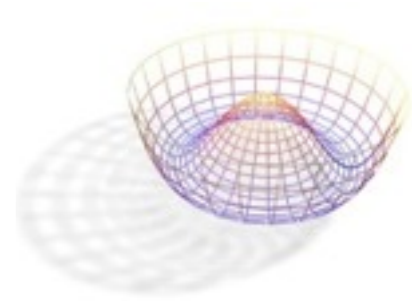
fusion



radioactivity (β decay)

beta minus decay



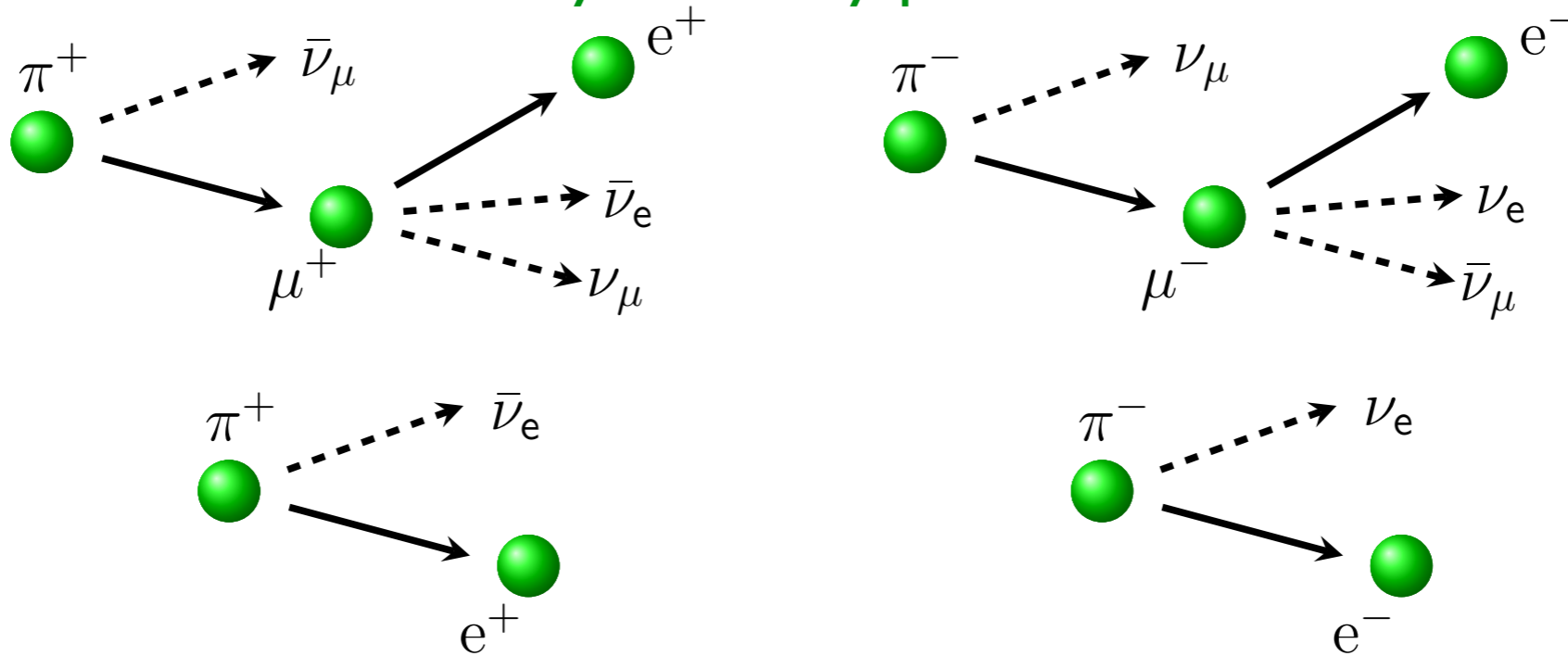


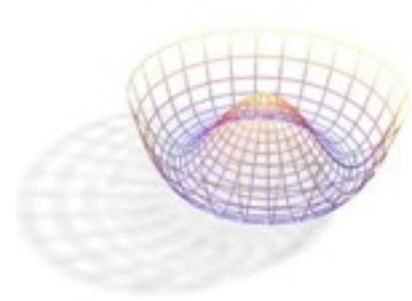
The Weak Interaction

responsible for all nucleonic transmutations

and particle decays

Decays of heavy particles



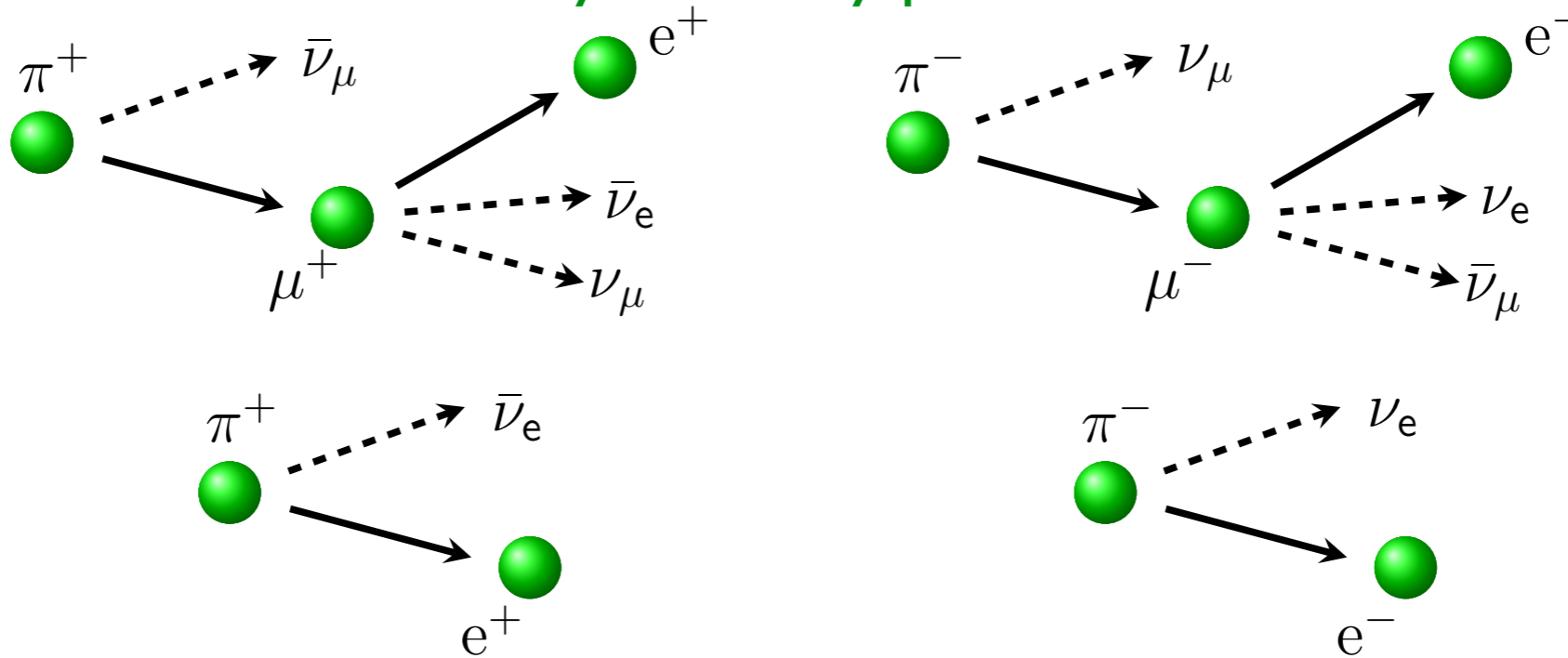


The Weak Interaction

responsible for all nucleonic transmutations

and particle decays

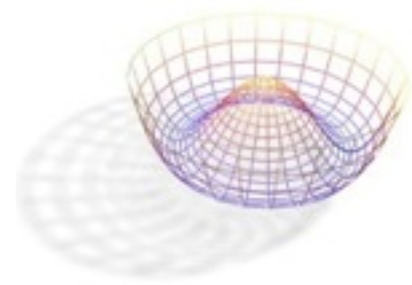
Decays of heavy particles



- Truly a weak interaction:

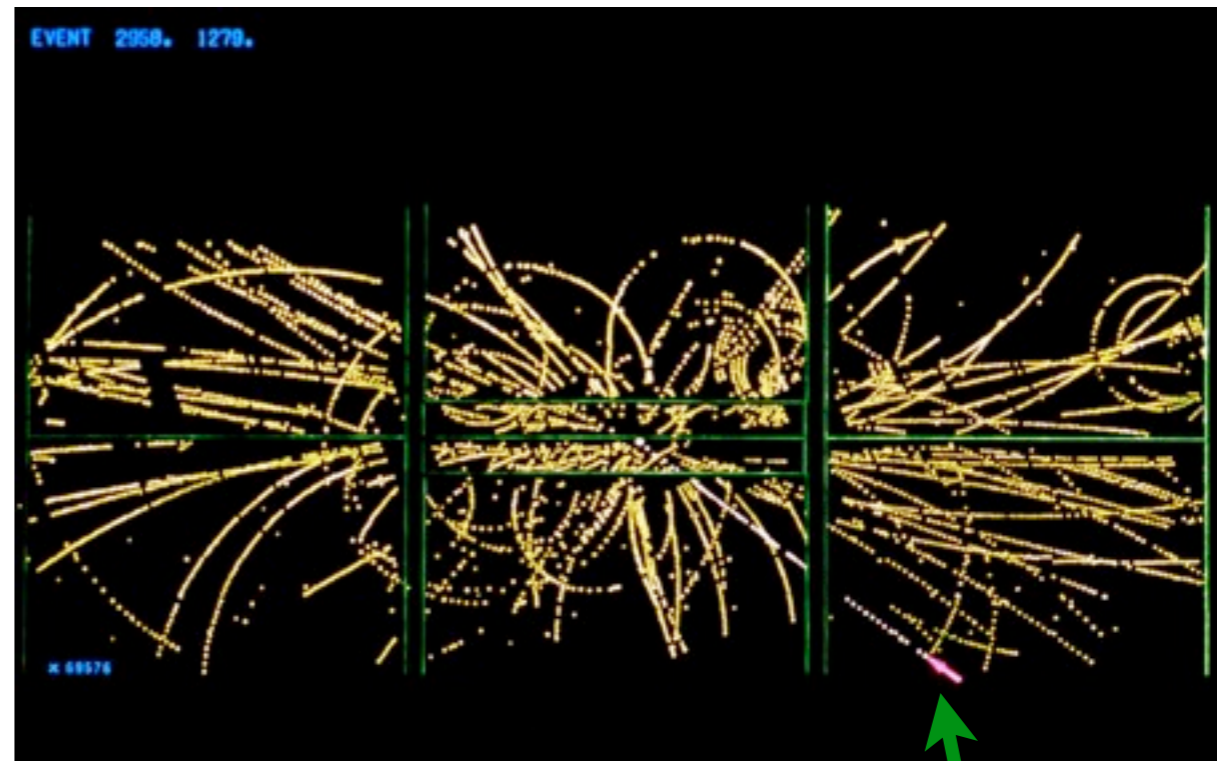
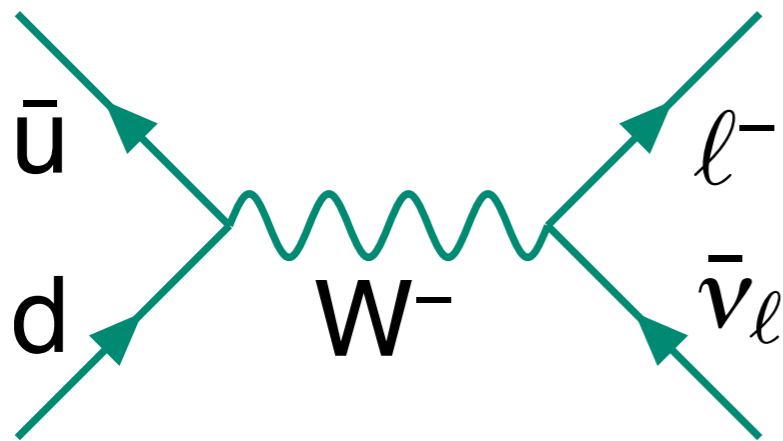
- solar ν flux on Earth: $\sim 6 \cdot 10^{14} \text{ m}^{-2} \text{ s}^{-1}$
- during your lifetime, **at most a few** will interact with your body at all!

The Weak Interaction



Exchange / production of **heavy** particles!

W-boson production
(and decay)



e^+

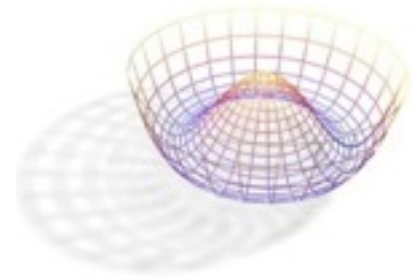
W and Z particles are heavy!

- $M_W = 80.398(25)$ GeV (\sim Sr, Kr)
- $M_Z = 92.188(2)$ GeV (\sim Ru)

Discovered in $p\bar{p}$ collisions, $E_{CM} = 630$ GeV

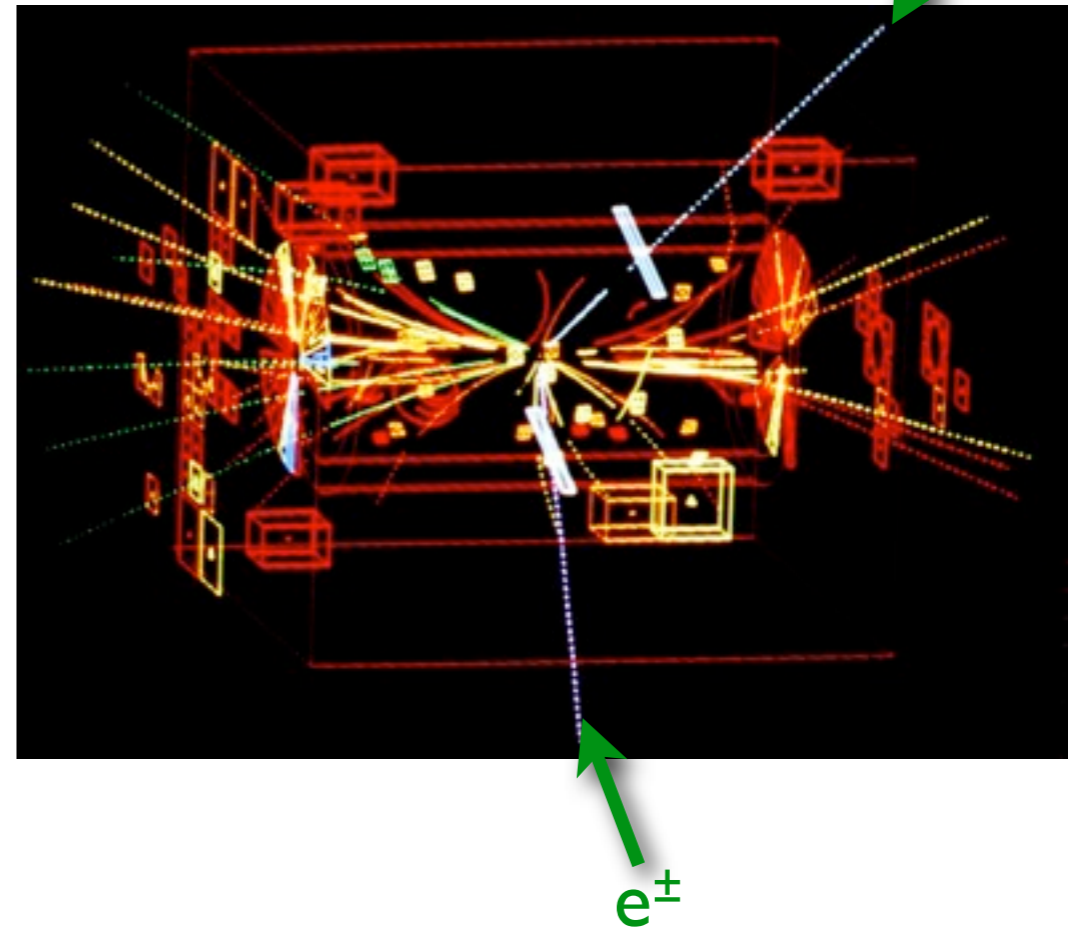
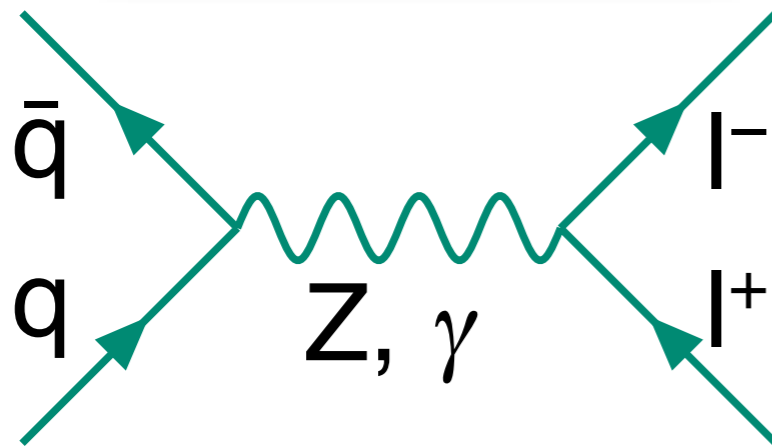
Most collisions between protons involve the strong interaction \Rightarrow look for **leptons** (only EM and weak interactions)

The Weak Interaction



Exchange / production of **heavy** particles!

Z-boson production and decay



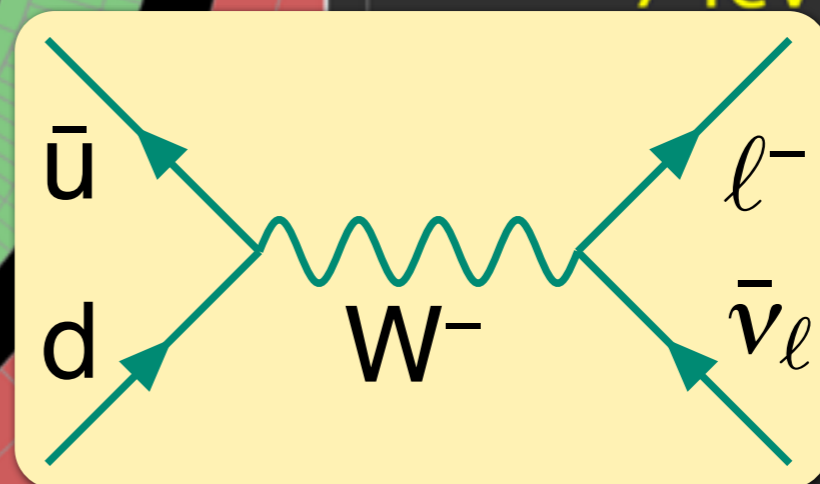
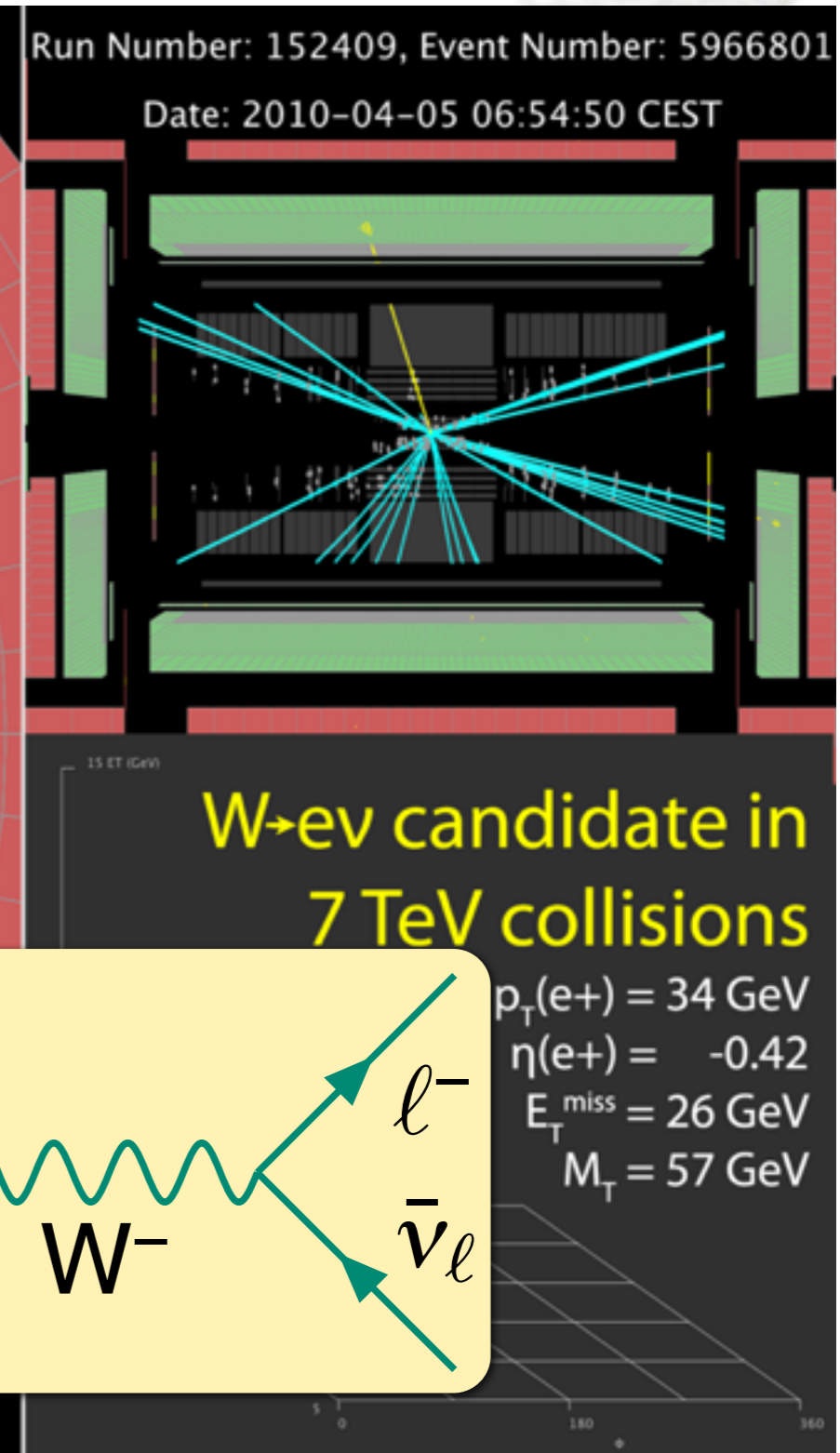
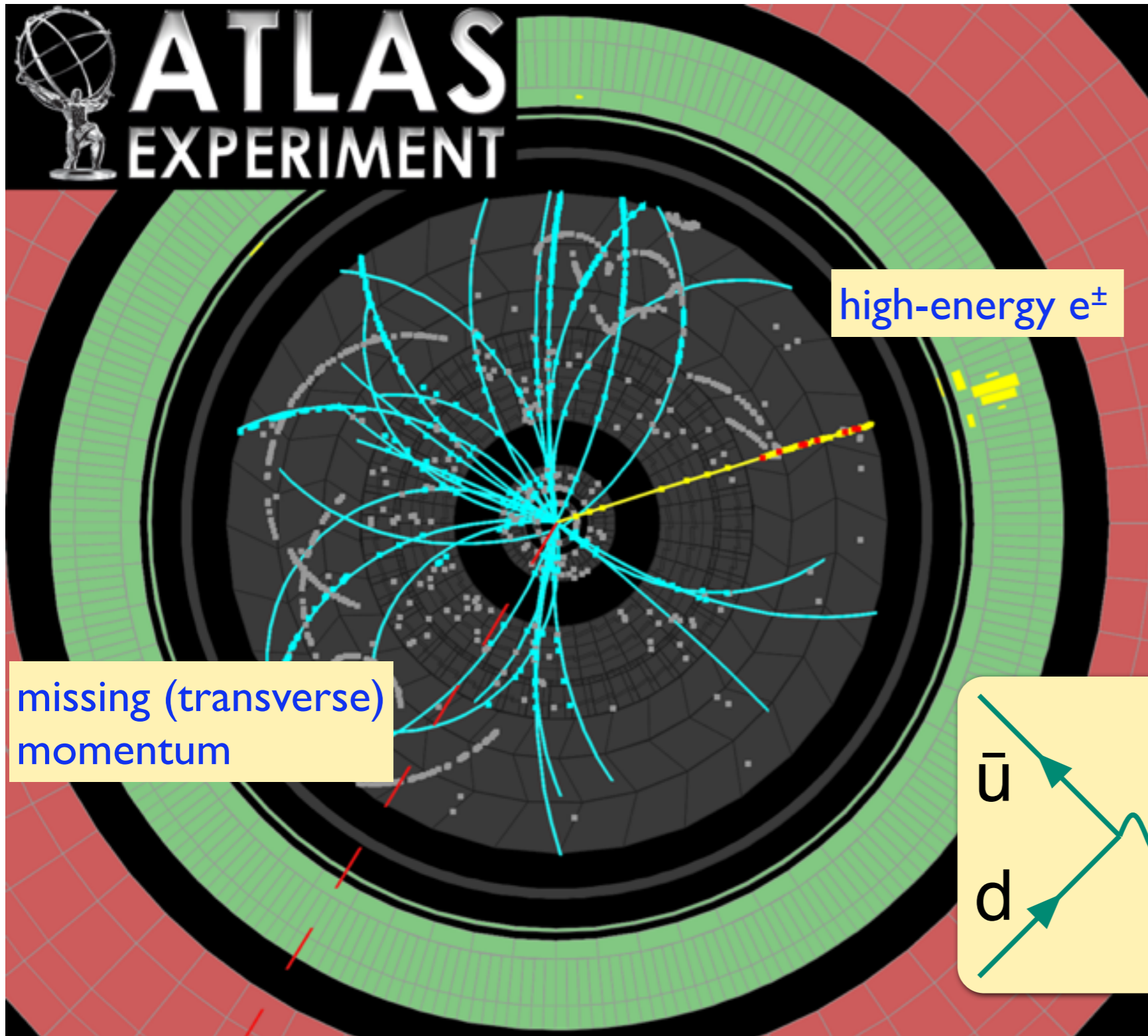
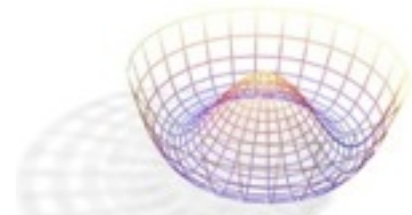
W and Z particles are heavy!

- $M_W = 80.398(25) \text{ GeV}$ ($\sim \text{Sr, Kr}$)
- $M_Z = 92.188(2) \text{ GeV}$ ($\sim \text{Ru}$)

Discovered in $p\text{-}\bar{p}$ collisions, $E_{\text{CM}} = 630 \text{ GeV}$

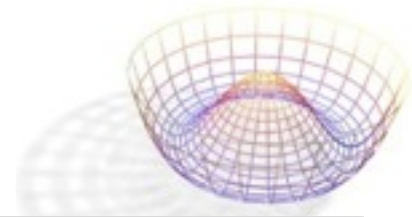
Most collisions between protons involve the strong interaction \Rightarrow look for **leptons** (only EM and weak interactions)

The Weak Interaction: W Boson



$M_W = 80.387 \pm 0.017 \text{ GeV} \text{ (~ krypton)}$

The Weak Interaction: Z Boson



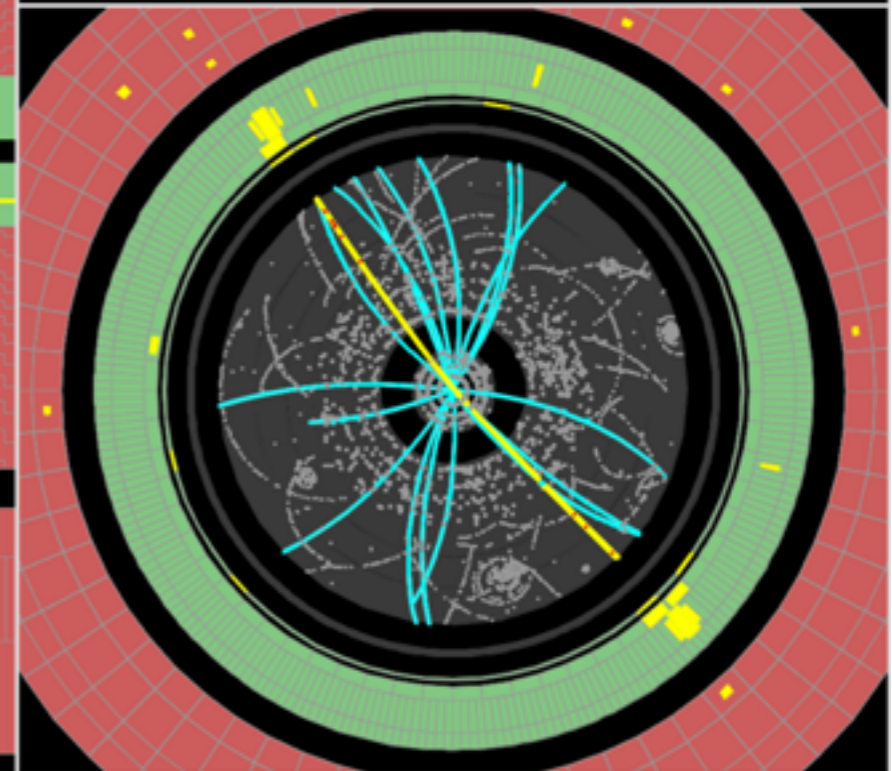
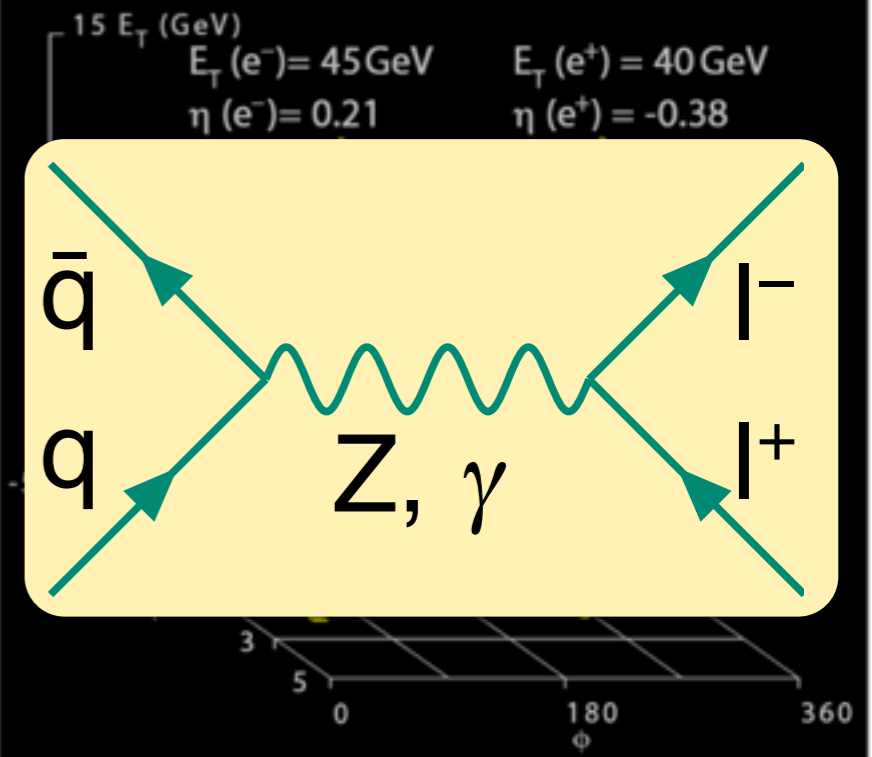
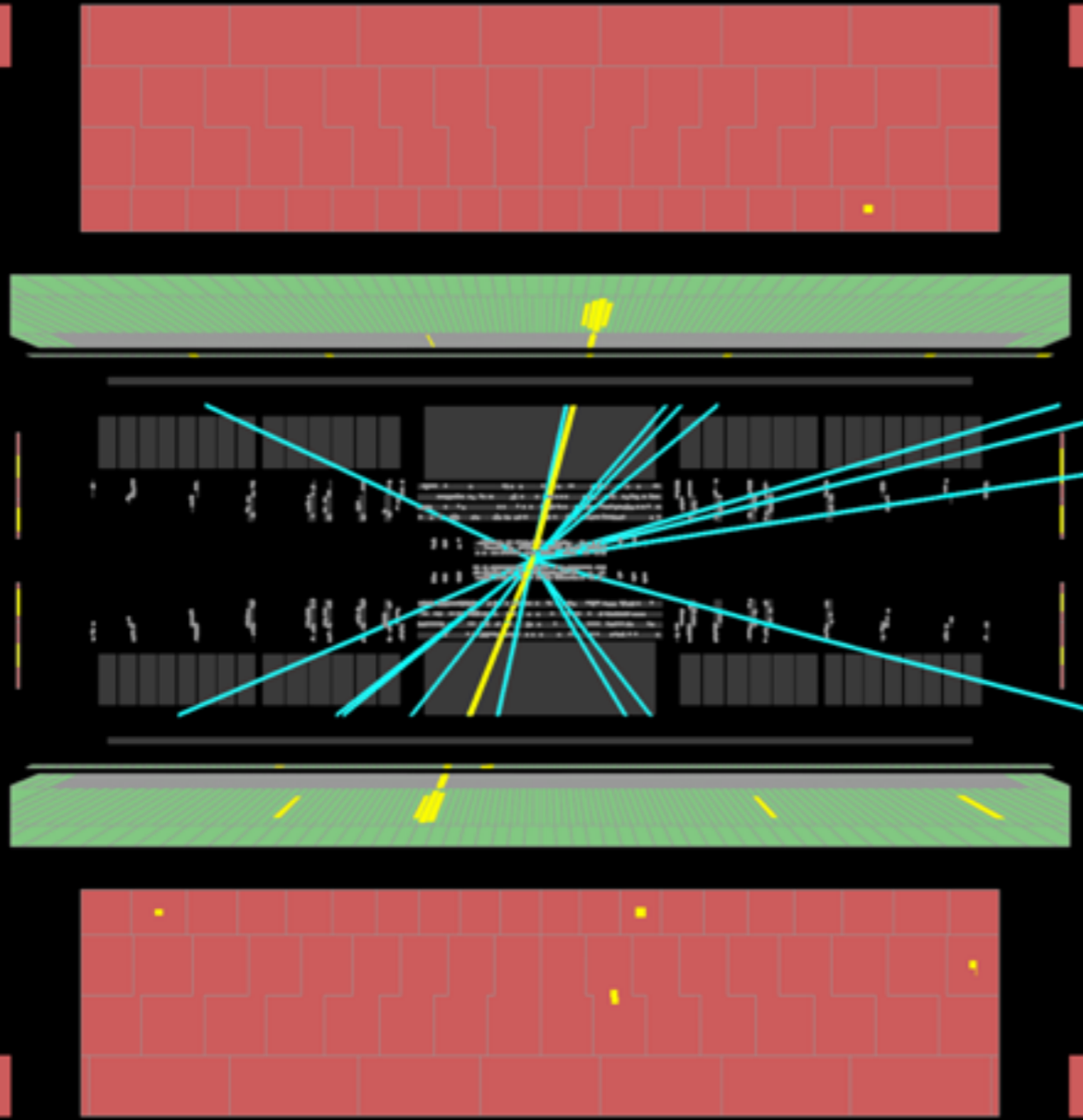
ATLAS
EXPERIMENT

Run Number: 154817, Event Number: 968871

Date: 2010-05-09 09:41:40 CEST

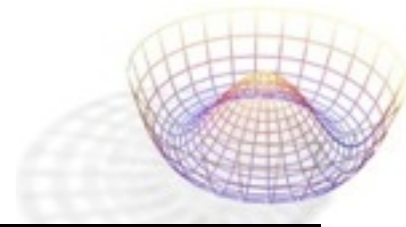
$M_{ee} = 89 \text{ GeV}$

$Z \rightarrow ee$ candidate in 7 TeV collisions



$$M_Z = 91.188 \pm 0.002 \text{ GeV} \quad (\sim \text{rubidium})$$

The Particle Family



“Leptons”??

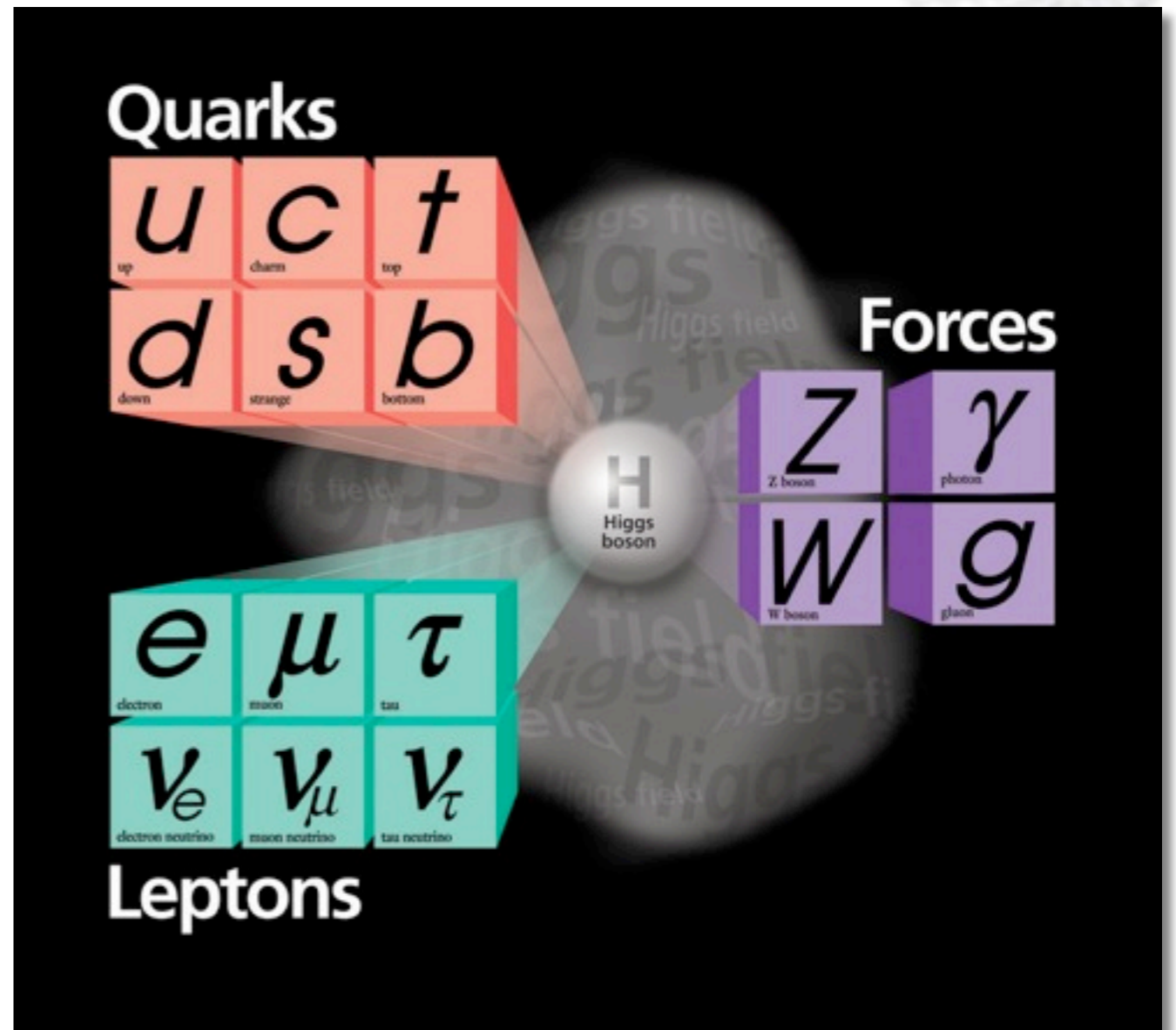
- particles not involved in the strong interaction (only weak, EM)
- also heavier counterparts of the electron and its neutral partner, ν_e

Quarks:

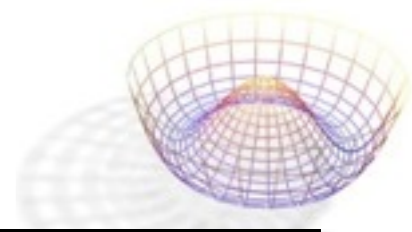
- particles also susceptible to the strong interaction
- again, 3 “generations” involving heavier partners than the **u, d** that are constituents of the proton

Force carriers:

- photon (EM), W/Z (weak interaction), gluon(s) (strong interaction)



The Particle Family



“Leptons”??

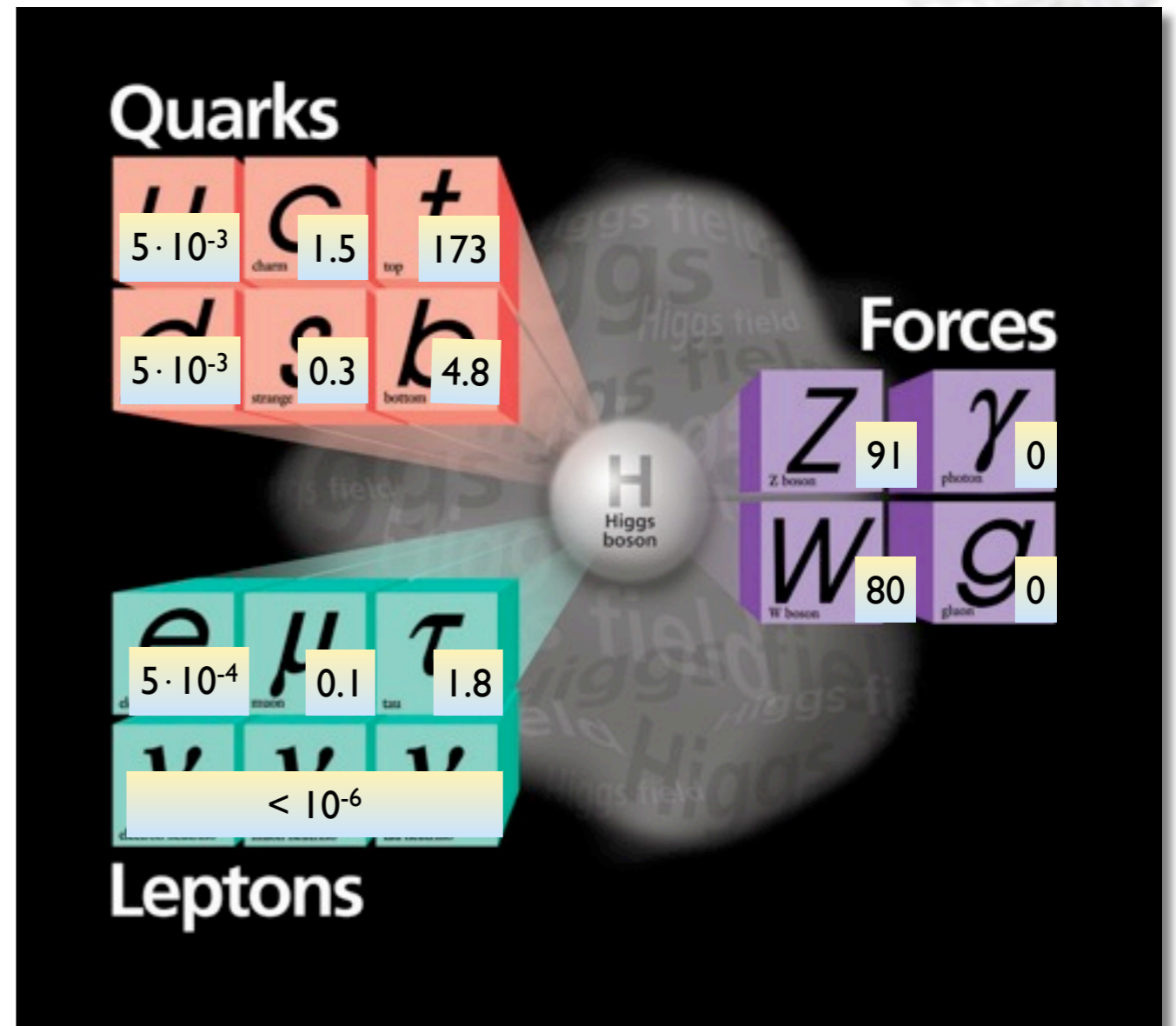
- particles not involved in the strong interaction (only weak, EM)
- also heavier counterparts of the electron and its neutral partner, ν_e

Quarks:

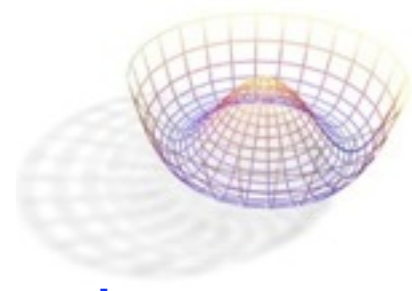
- particles also susceptible to the strong interaction
- again, 3 “generations” involving heavier partners than the **u, d** that are constituents of the proton

Force carriers:

- photon (EM), W/Z (weak interaction), gluon(s) (strong interaction)



The Higgs Mechanism



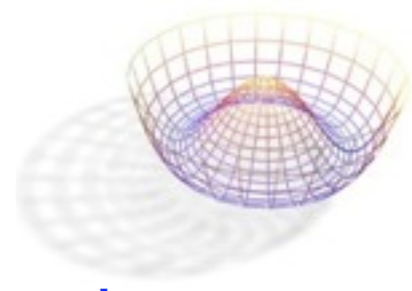
Particles become “effectively” massive by means of their interaction with the Higgs field!



More physical analogy: refractive index

- caused by **different speed of light** in medium
- caused by **forward scattering** of light by the medium

The Higgs Mechanism



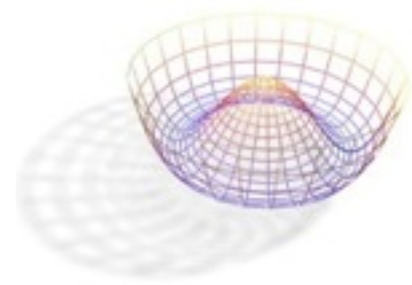
Particles become “effectively” massive by means of their interaction with the Higgs field



More physical analogy: refractive index

- caused by **different speed of light** in medium
- caused by **forward scattering** of light by the medium

The Higgs Mechanism



Particles become “effectively” massive by means of their interaction with the Higgs field



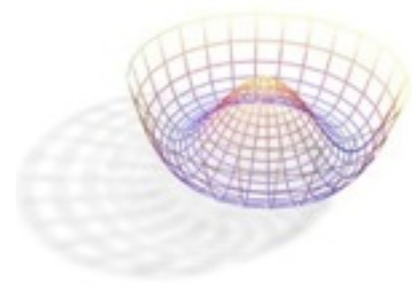
More physical analogy: refractive index

- caused by **different speed of light** in medium
- caused by **forward scattering** of light by the medium

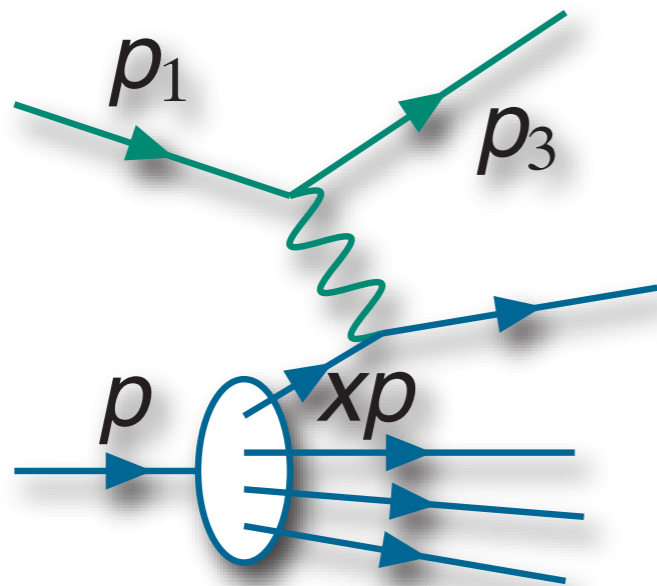
Photons in the medium are effectively massive



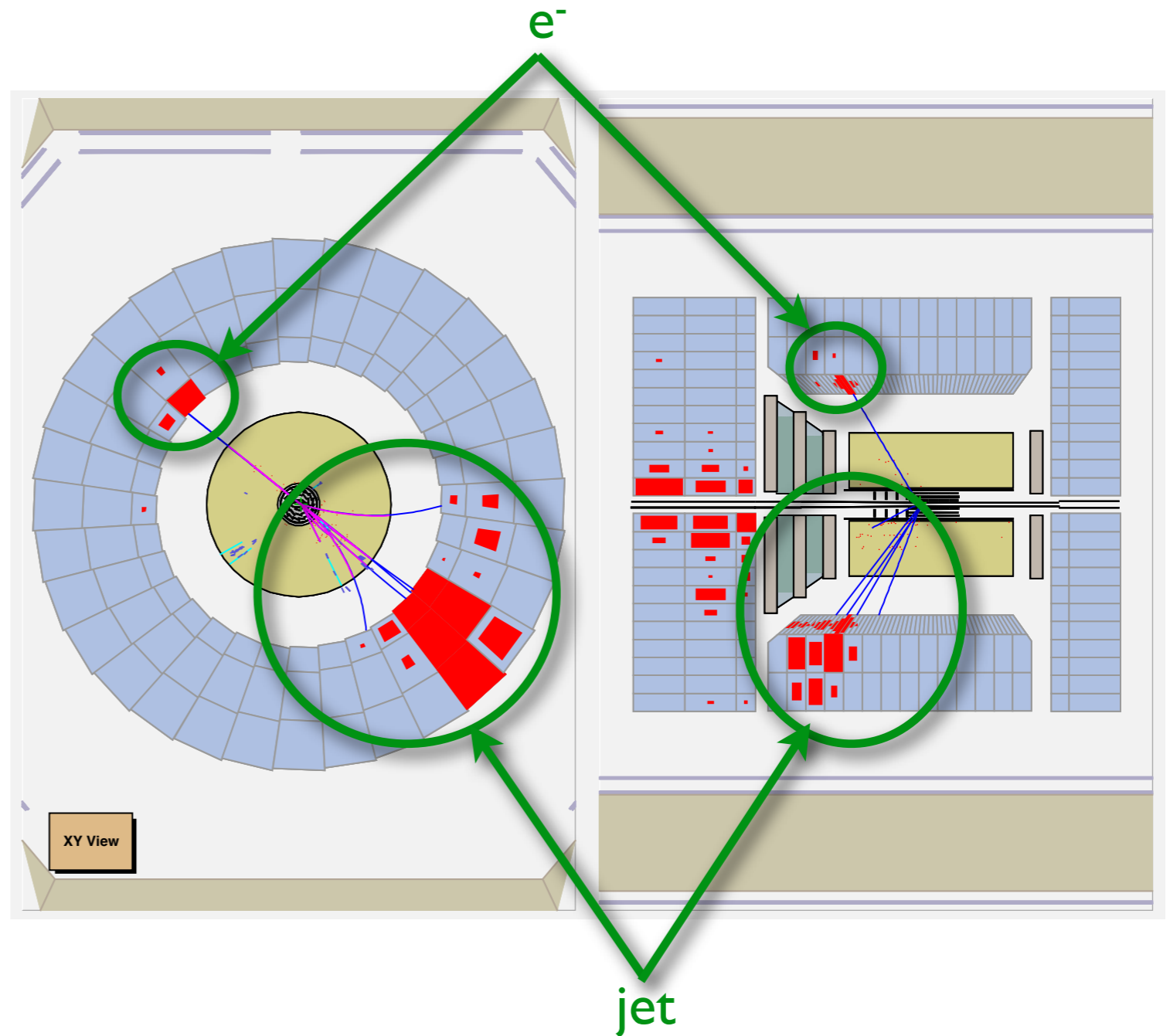
QCD at High Energies



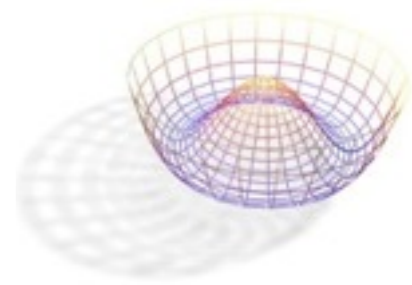
At high energies, quarks and gluons do manifest themselves as “free” particles → hadron **jets**



electron-proton scattering:
27.5 GeV + 920 GeV



A Weighty Issue...

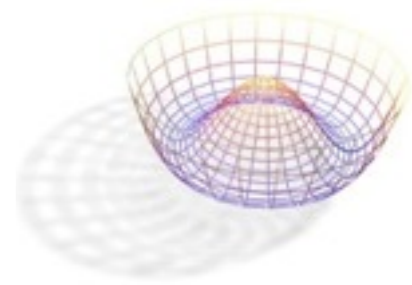


QED, QCD: photon & gluons are strictly massless

Weak interaction:

- massive W and Z bosons
- fermion masses: $m_\ell \neq m_{\nu_\ell}$ (and similarly for quarks)

A Weighty Issue...



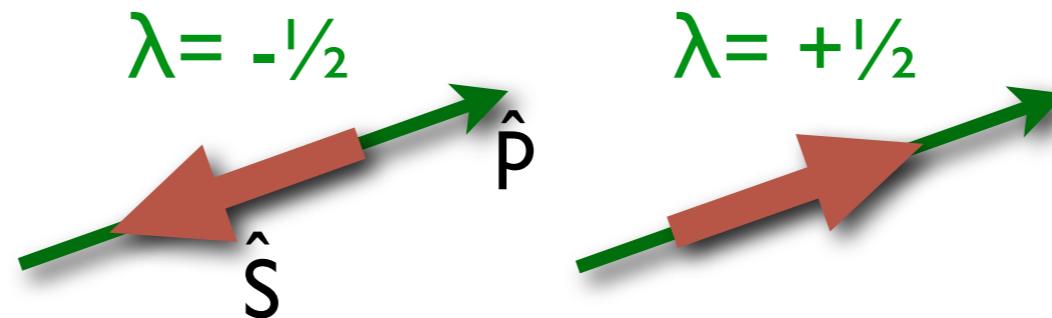
QED, QCD: photon & gluons are strictly massless

Weak interaction:

- massive W and Z bosons
- fermion masses: $m_e \neq m_{\nu_e}$ (and similarly for quarks)

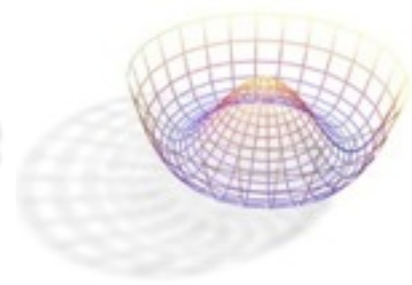
And worse!

- W-boson deals with **left-handed fermions** (right-handed anti-fermions) only



- left- and right-handed fermions should be **different particles**
- this requires them to be **strictly massless**

The Higgs Mechanism to the Rescue



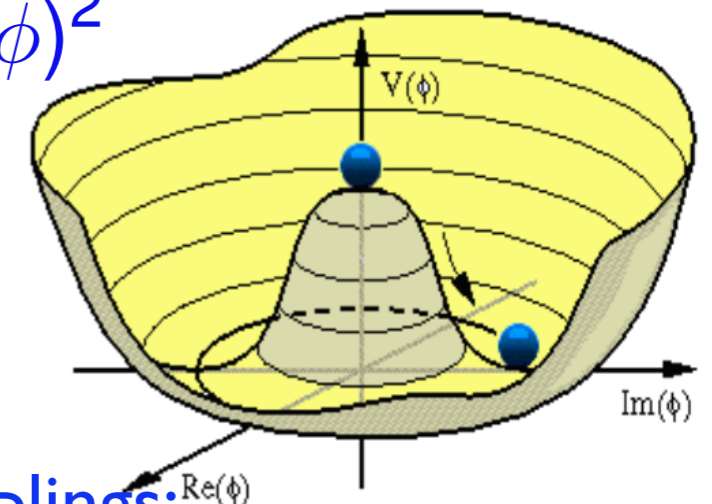
Required: a mechanism to break the EW symmetry **spontaneously**

- **Lagrangian** maintains full EW symmetry $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$
- but **the ground state does not!**

Achieved through the introduction of the (complex scalar) **Higgs field**

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} \quad V_\phi = \mu(\phi^\dagger \phi) + \lambda(\phi^\dagger \phi)^2$$

- With $\mu < 0$: minimum at $\phi \neq 0$

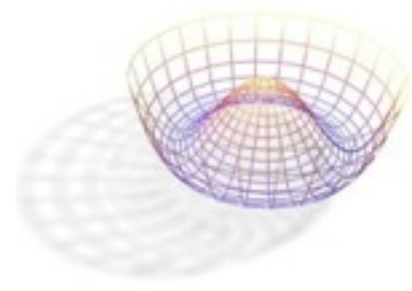


Generation of fermion masses through “Yukawa” couplings:

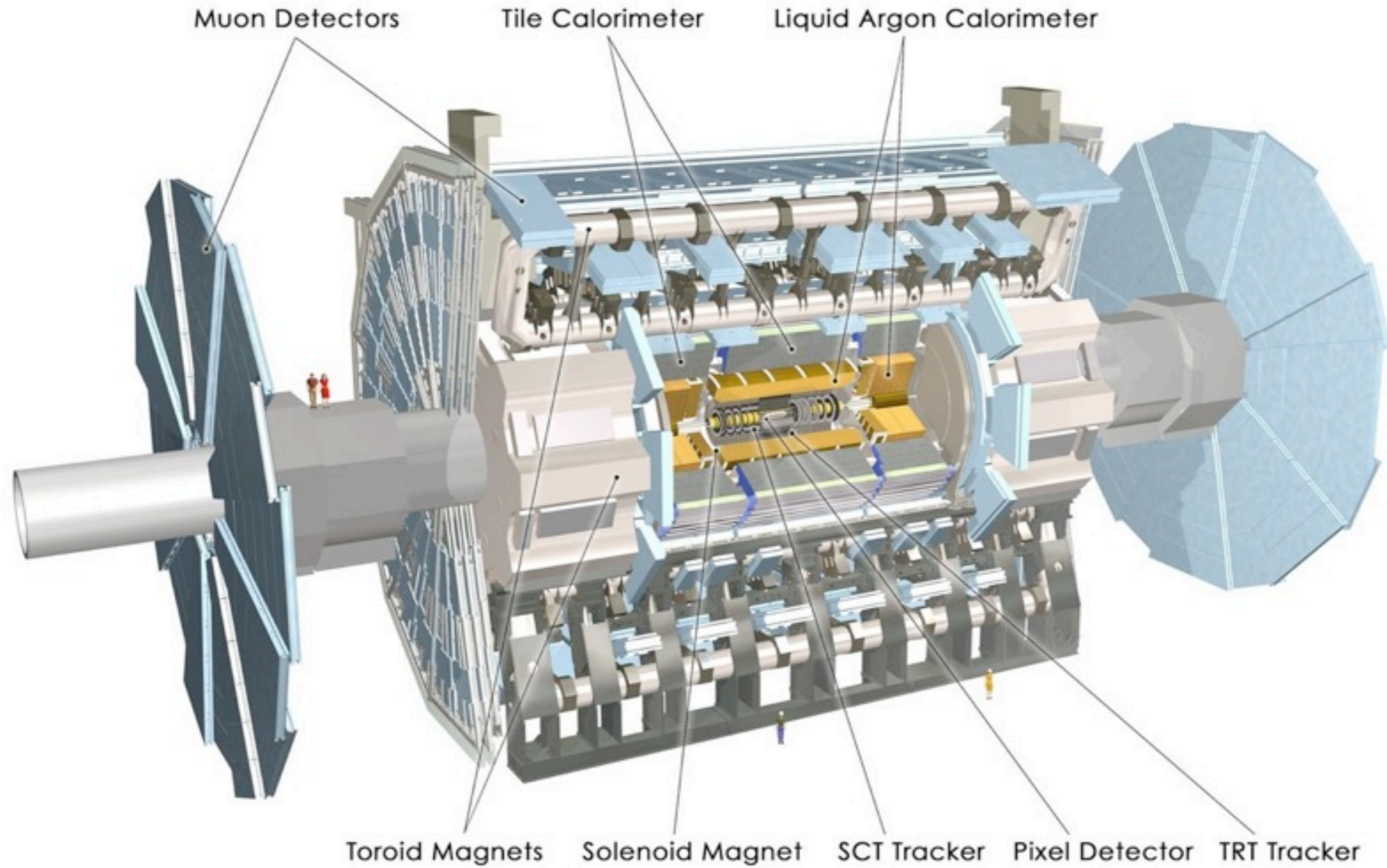
$$\mathcal{L}_Y = -g_e (\bar{e}_R \phi^\dagger \psi_L + \bar{\psi}_L \phi e_R)$$

$$(\langle \phi \rangle = \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix}) \rightarrow -\frac{g_e v}{\sqrt{2}} (\bar{e}_R e_L + \bar{e}_L e_R) = -\frac{g_e v}{\sqrt{2}} \bar{e} e$$

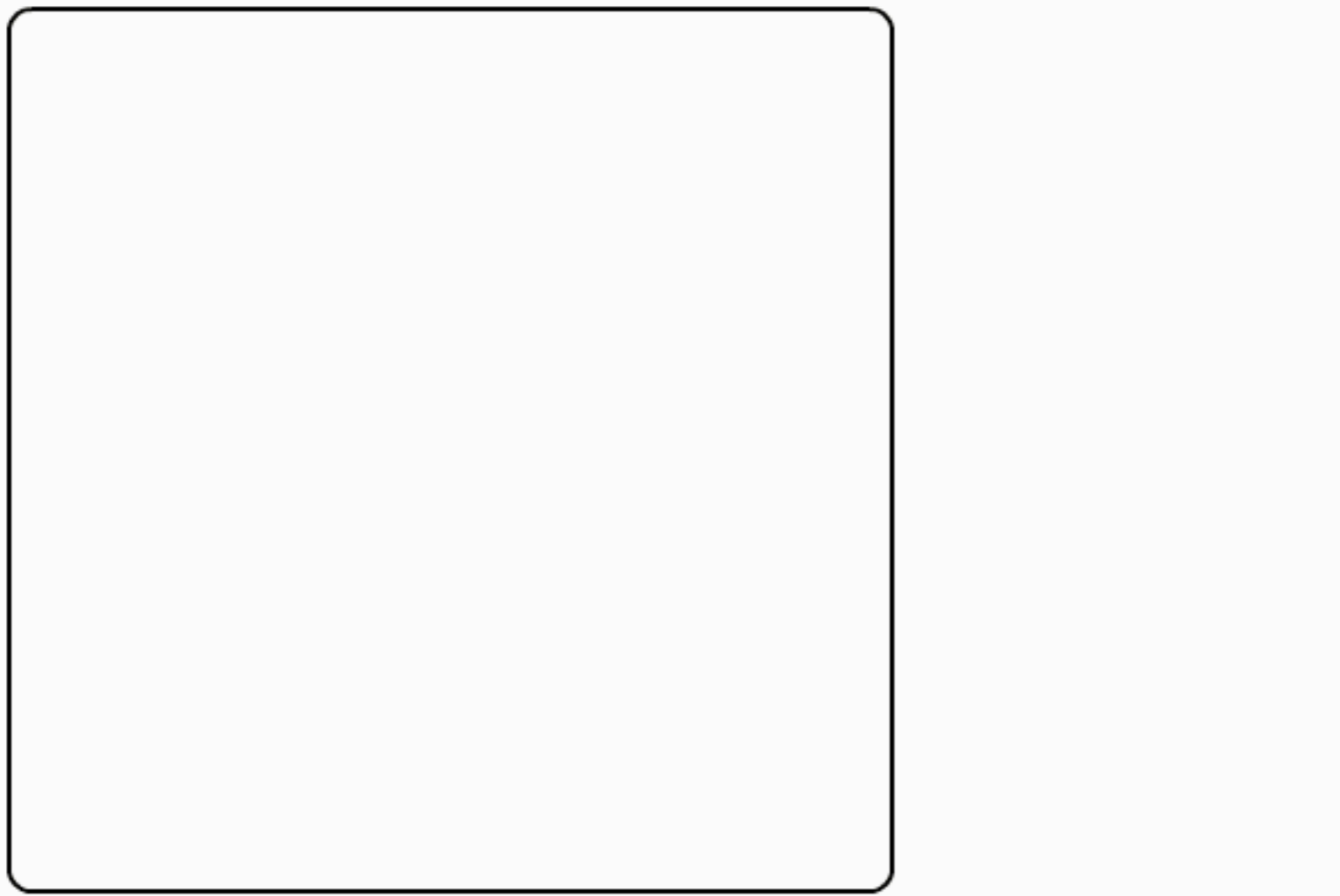
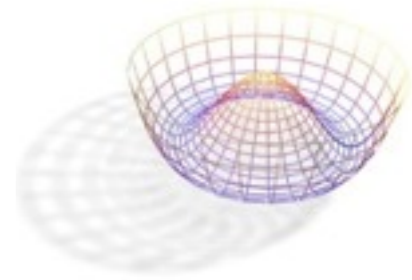
The Higgs Hunters



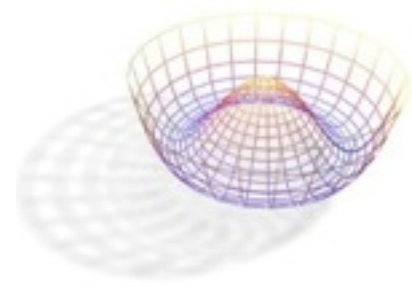
ATLAS... and CMS



Particle Detection



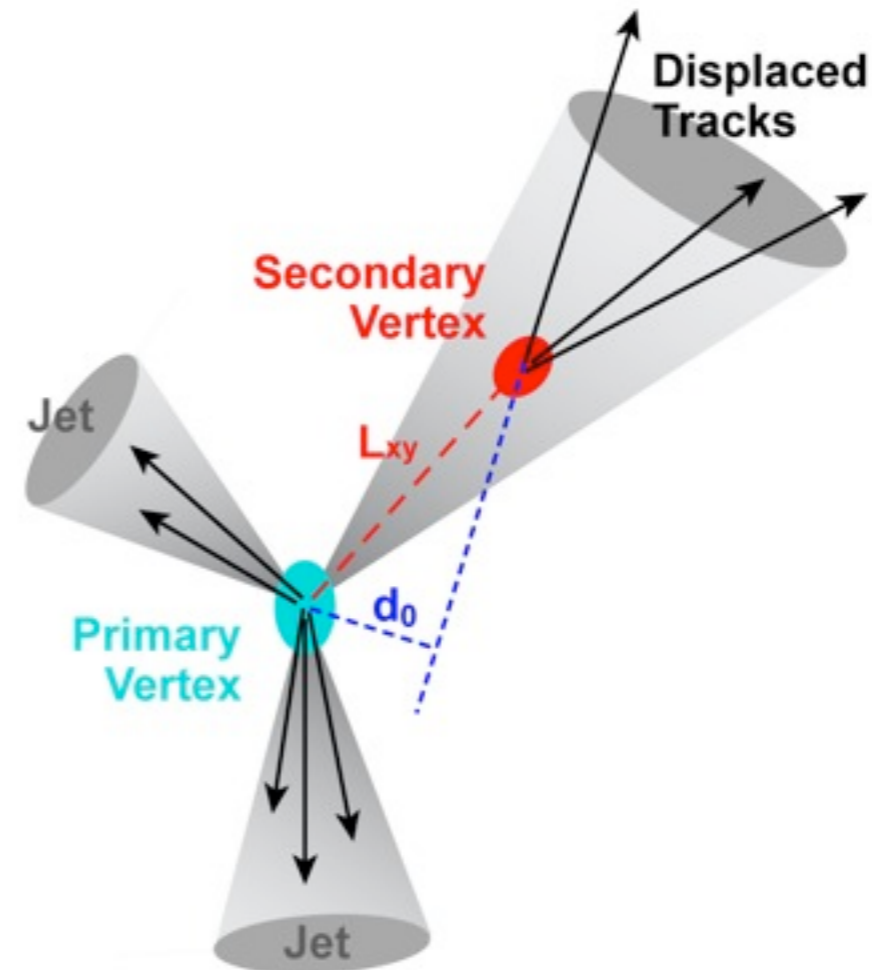
Particle Detection



In addition to individually observable particles:

- **neutrinos** (from apparent lack of momentum conservation)
- **hadron jets** (from calorimeter energy deposits/tracks)
- **τ leptons** (very narrow “hadronic jet”)
- **b-jets** (from hadronisation of b-quarks:
of B-hadrons, $\tau_B \approx 1.5$ ps)

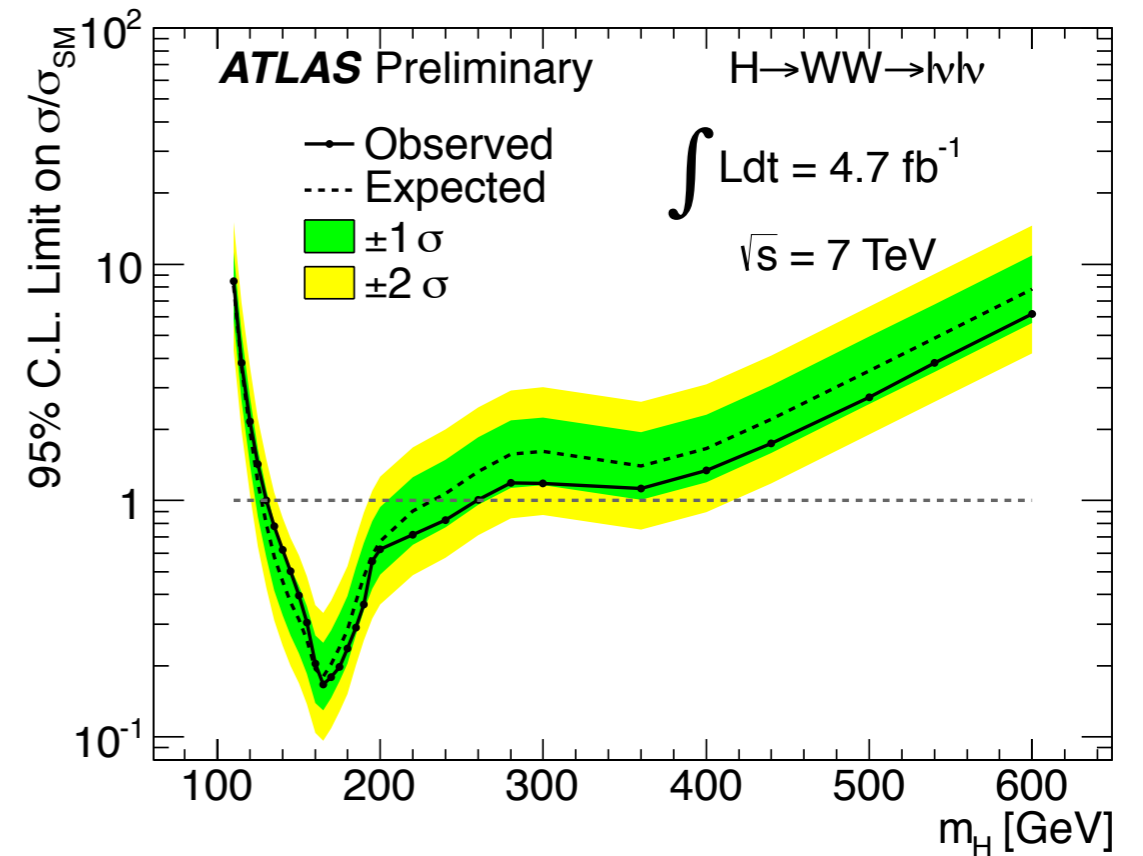
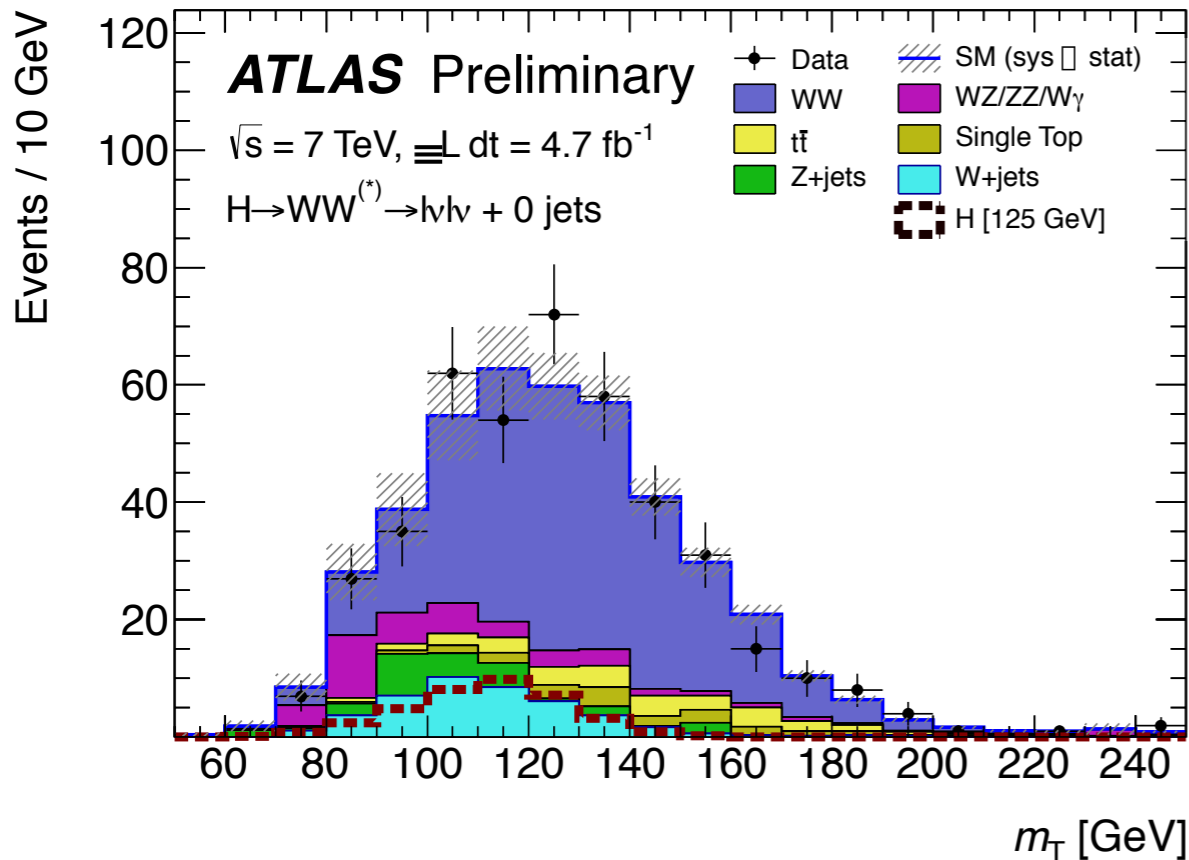
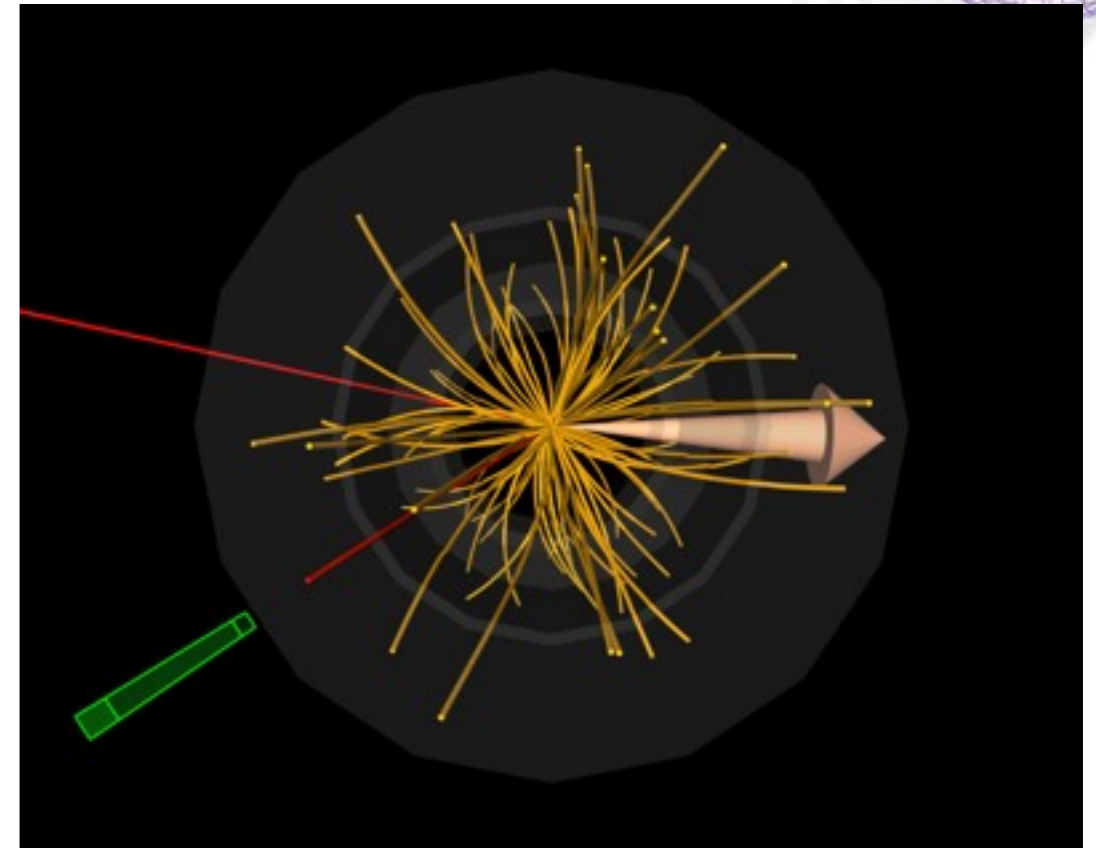
“long” lifetime



H \rightarrow W⁺W⁻

Relatively large event rate, but leptonic W boson decays lead to unobserved neutrinos

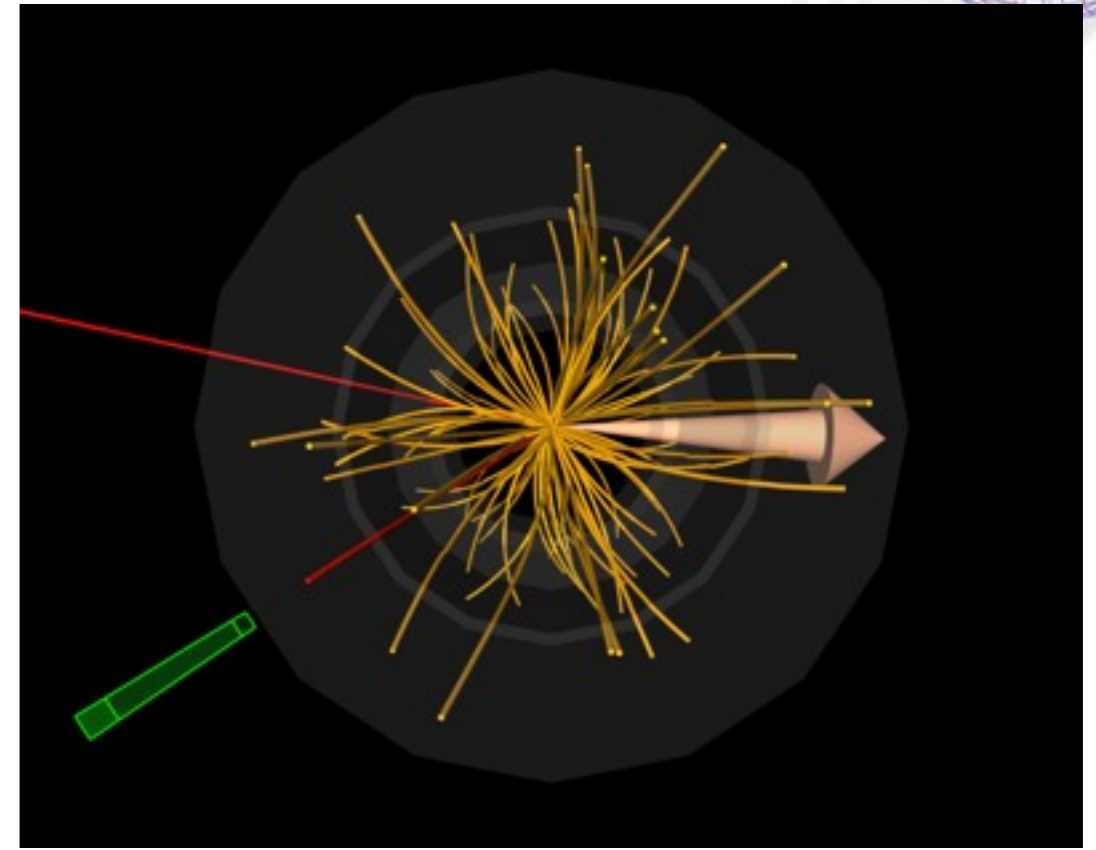
- cannot reconstruct mass of a system decaying to W⁺W⁻
- consider distribution of kinematic variables



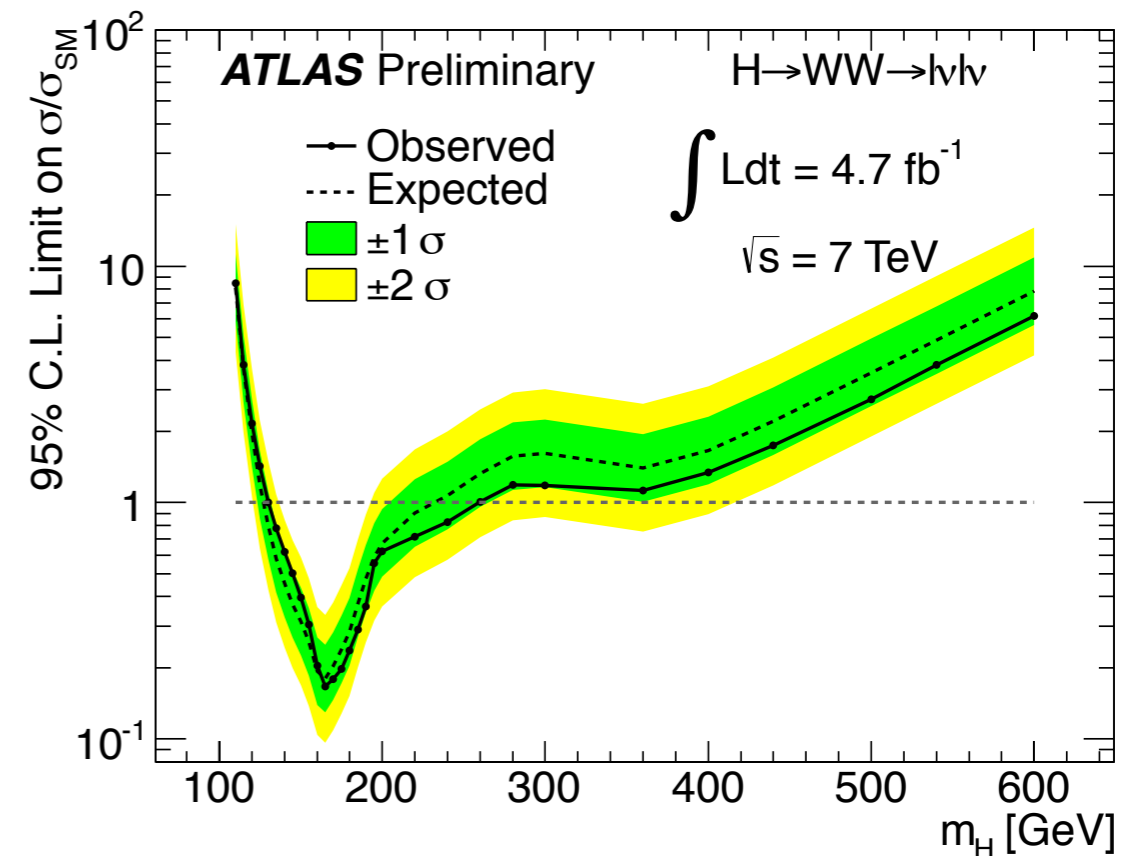


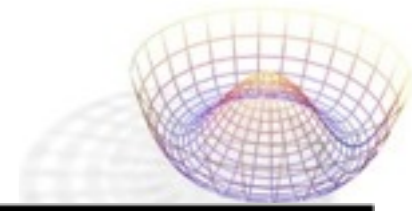
Relatively large event rate, but leptonic W boson decays lead to unobserved neutrinos

- cannot reconstruct mass of a system decaying to W^+W^-
- consider distribution of kinematic variables



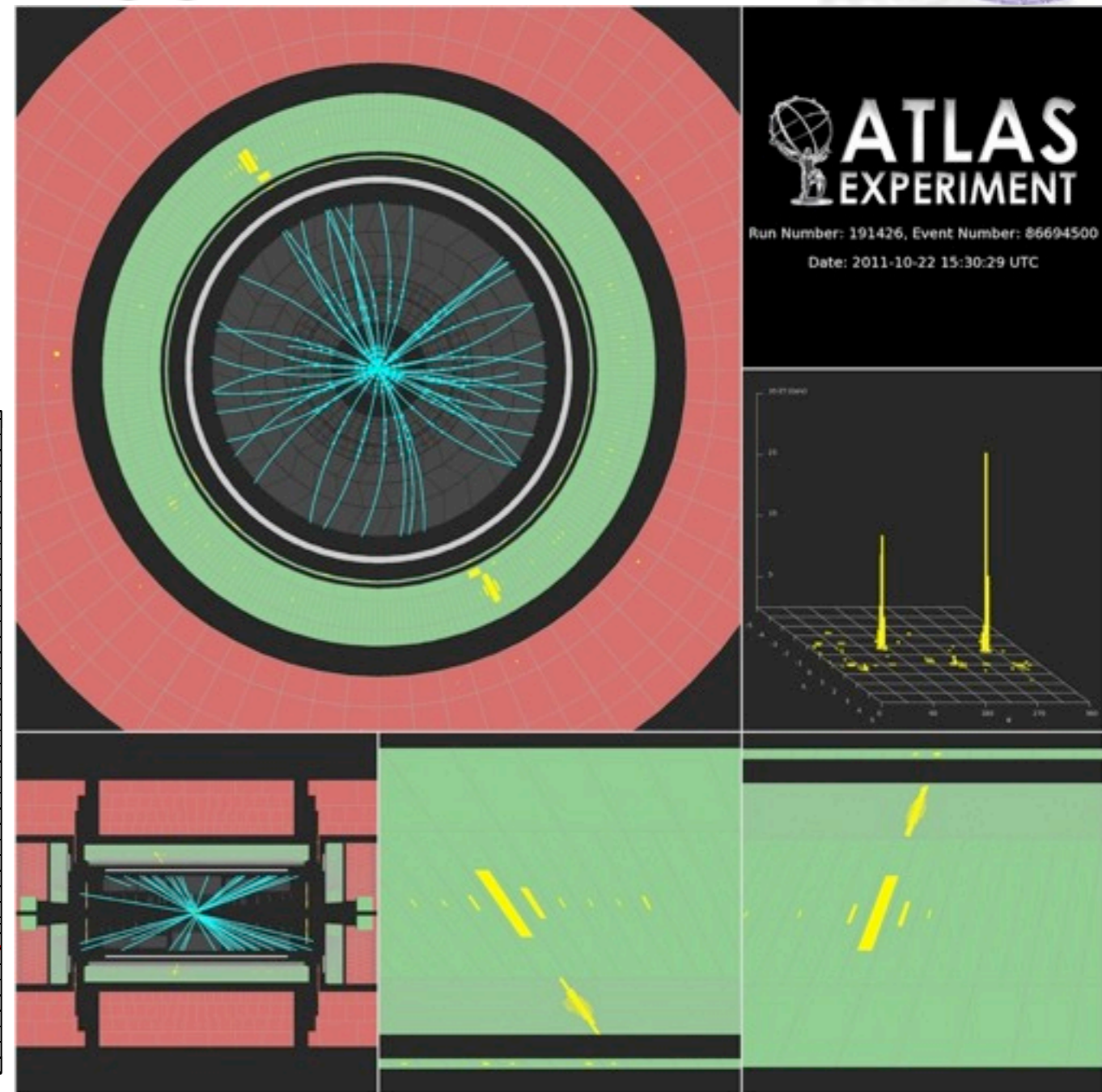
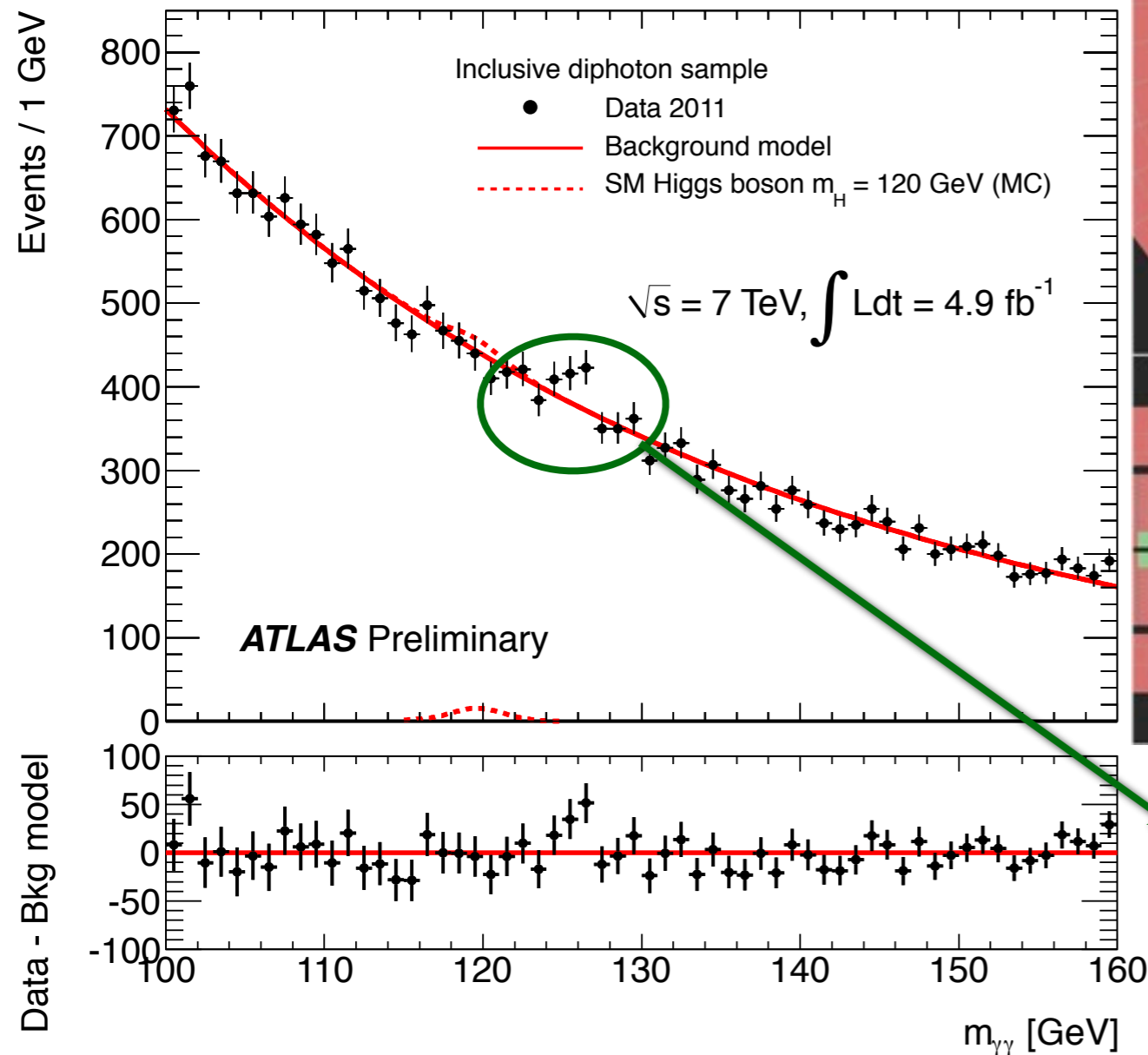
Exclude (at 95% confidence level) the Standard Model Higgs boson if $130 \text{ GeV} < M_H < 260 \text{ GeV}$



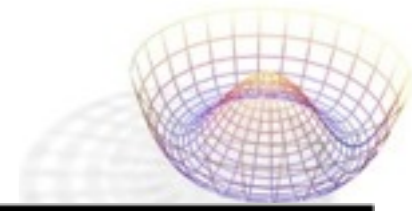


Requires excellent discrimination between single high-energy photons from hadrons

• but offers good energy resolution

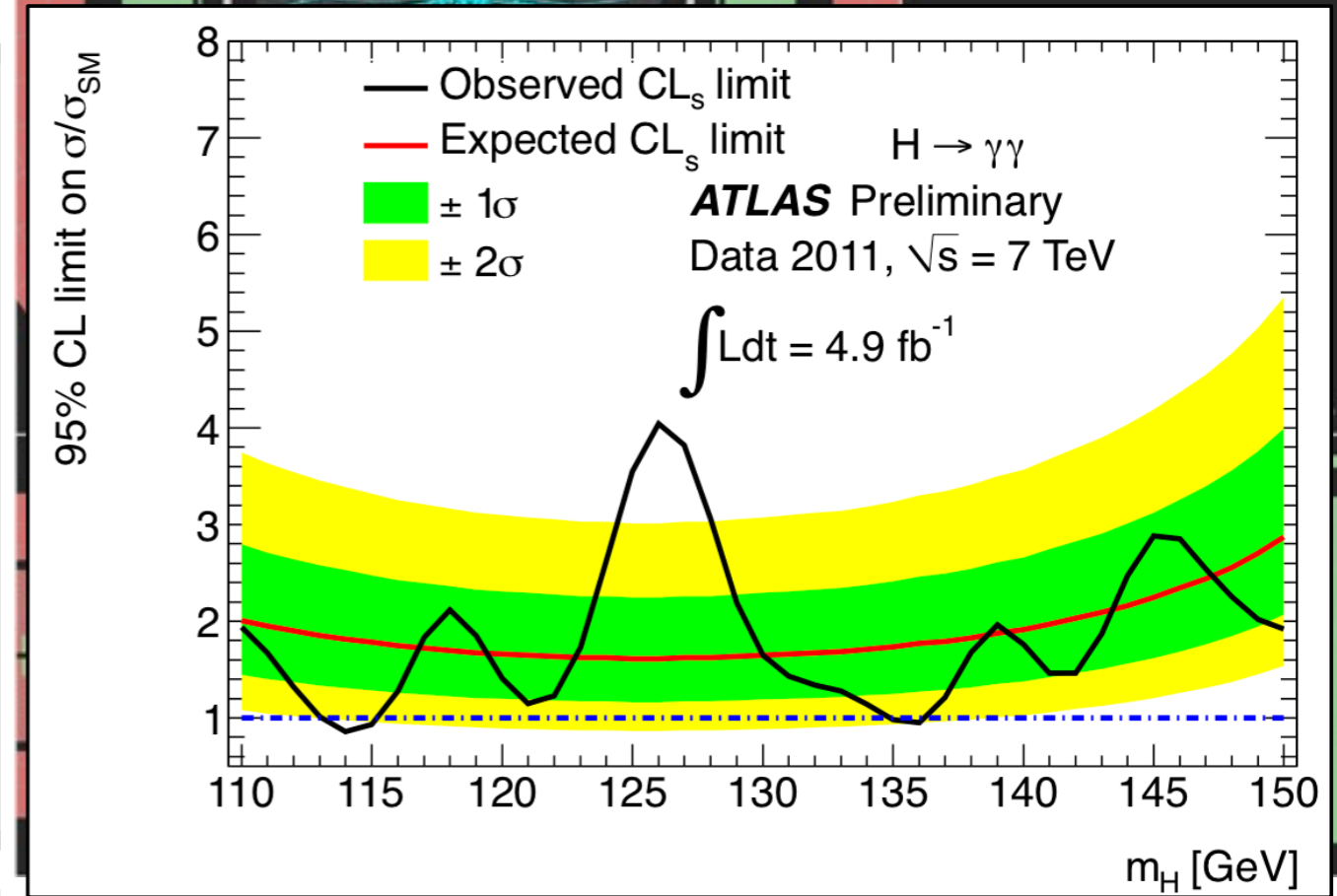
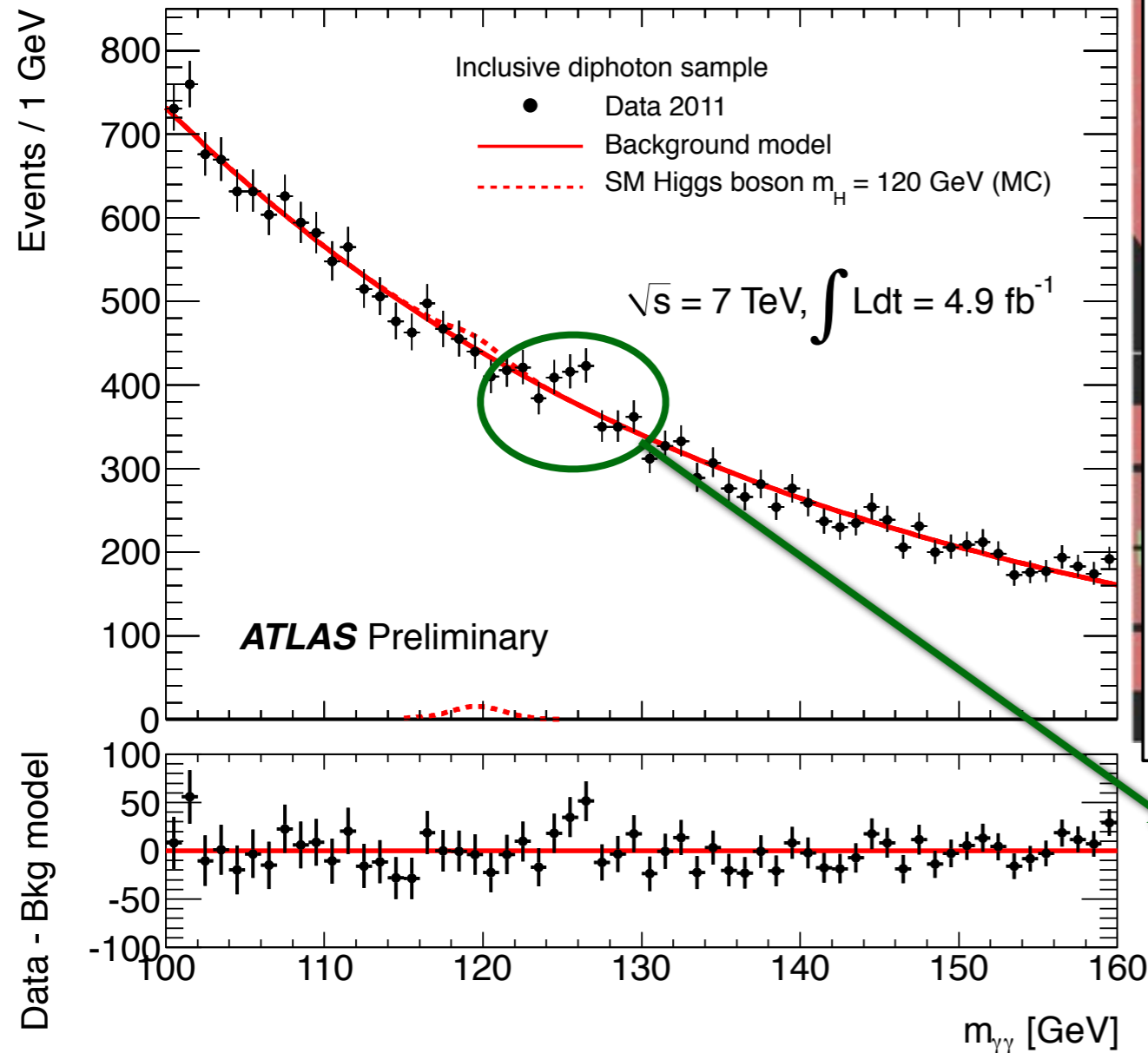
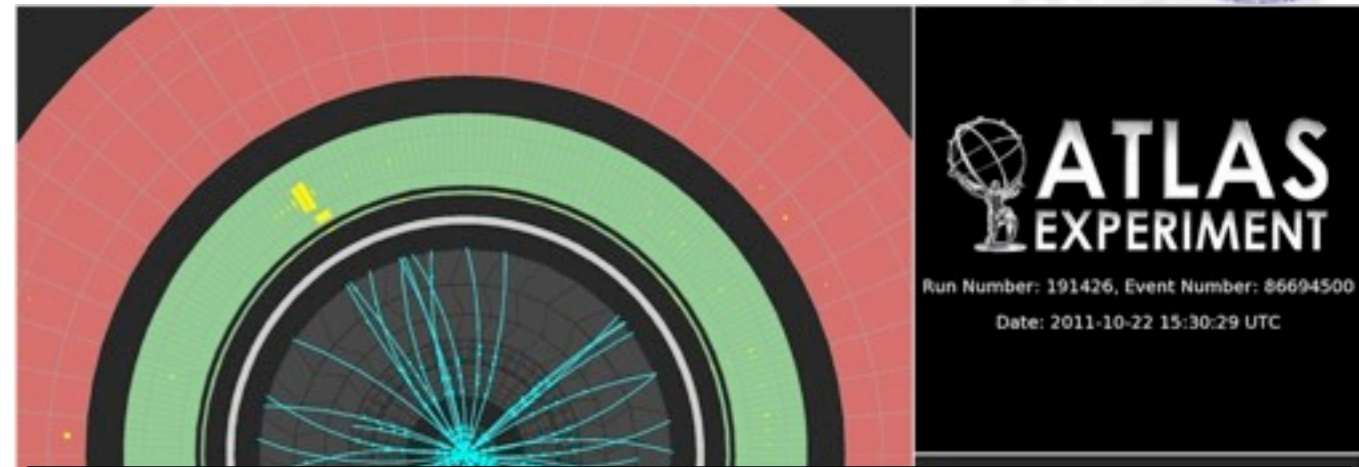


Looking for small excess on top of large (but smooth) background



Requires excellent discrimination between single high-energy photons from hadrons

• but offers good energy resolution

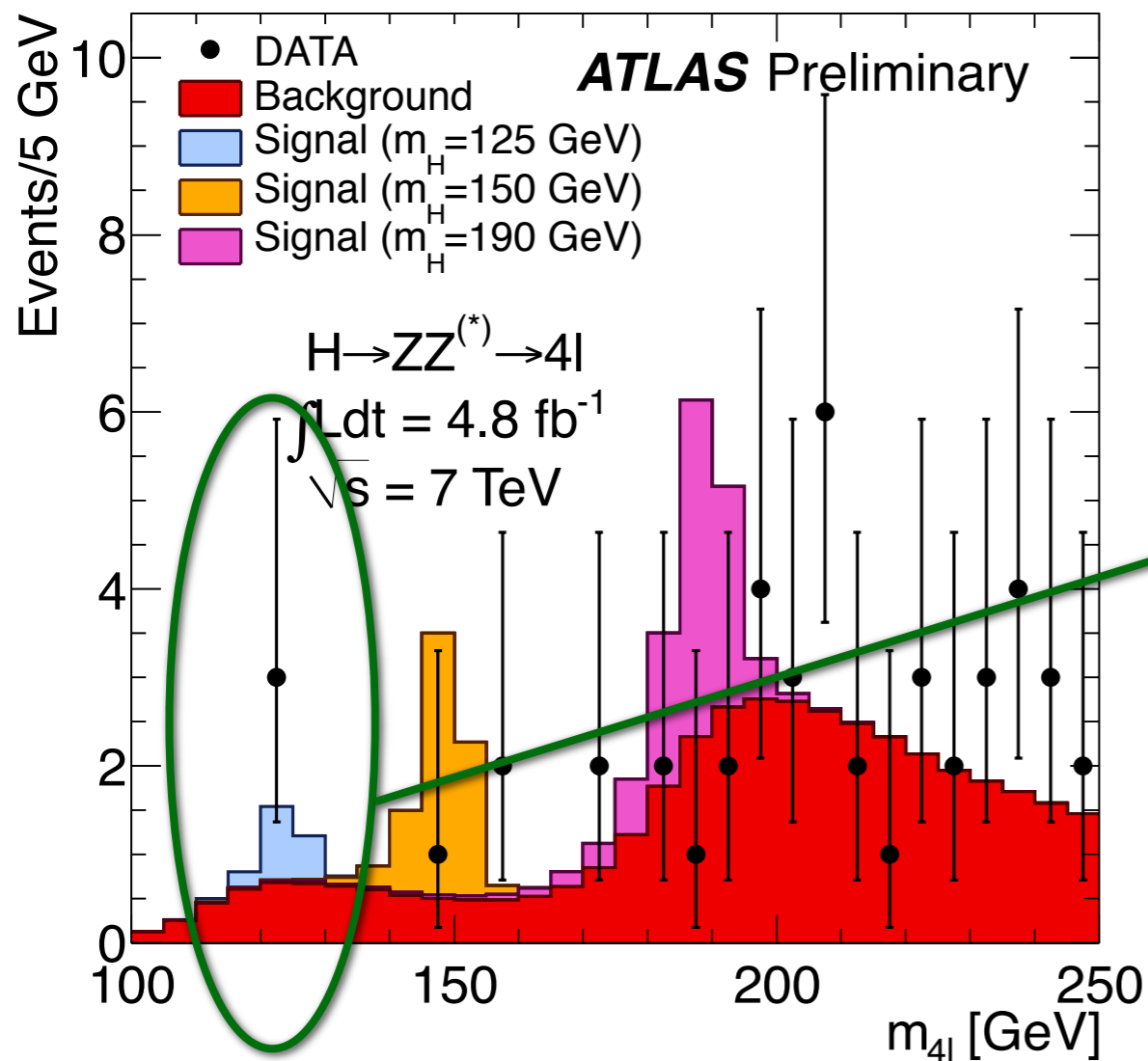
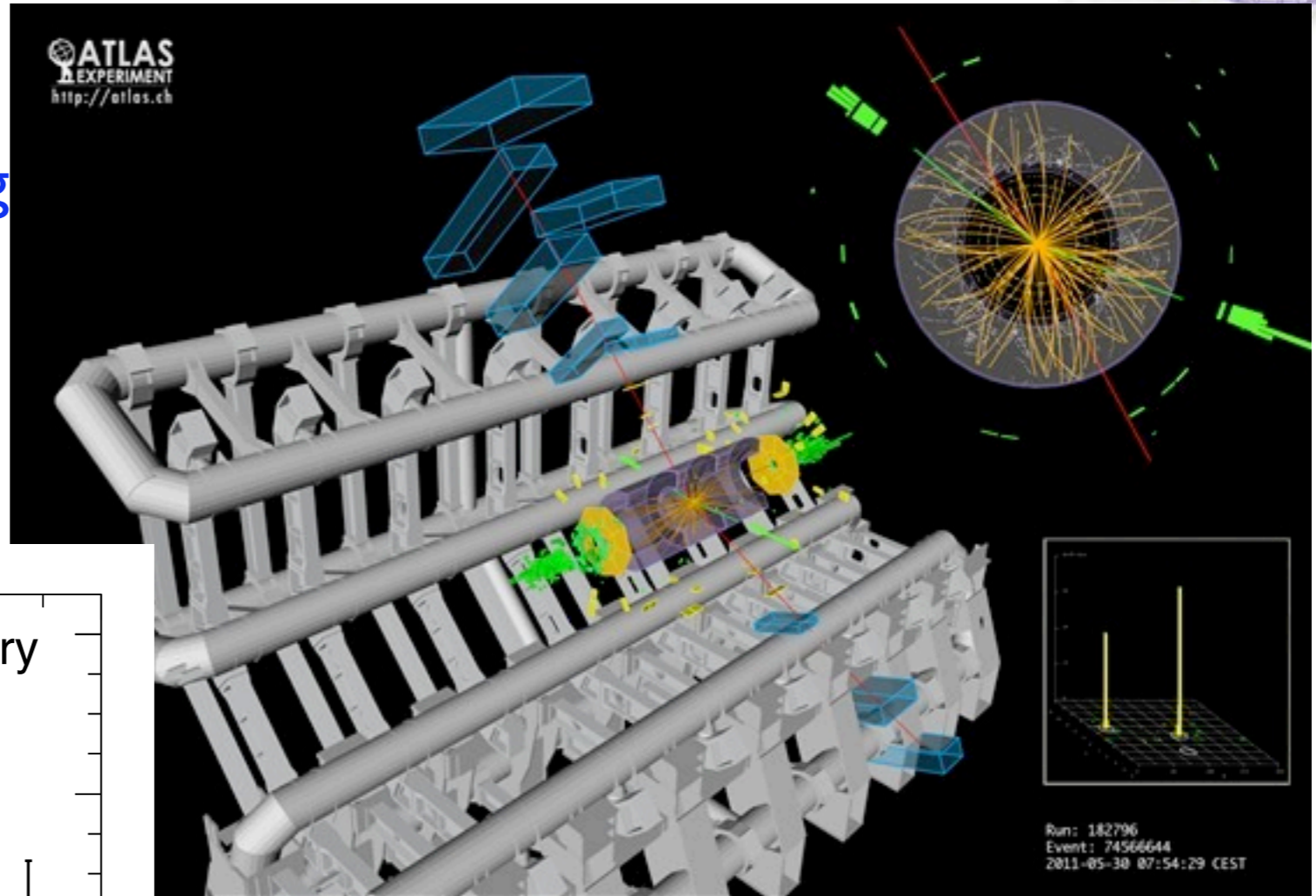


Looking for small excess on top of large (but smooth) background

H → ZZ

Very rare process, especially with both Z particles decaying to leptons

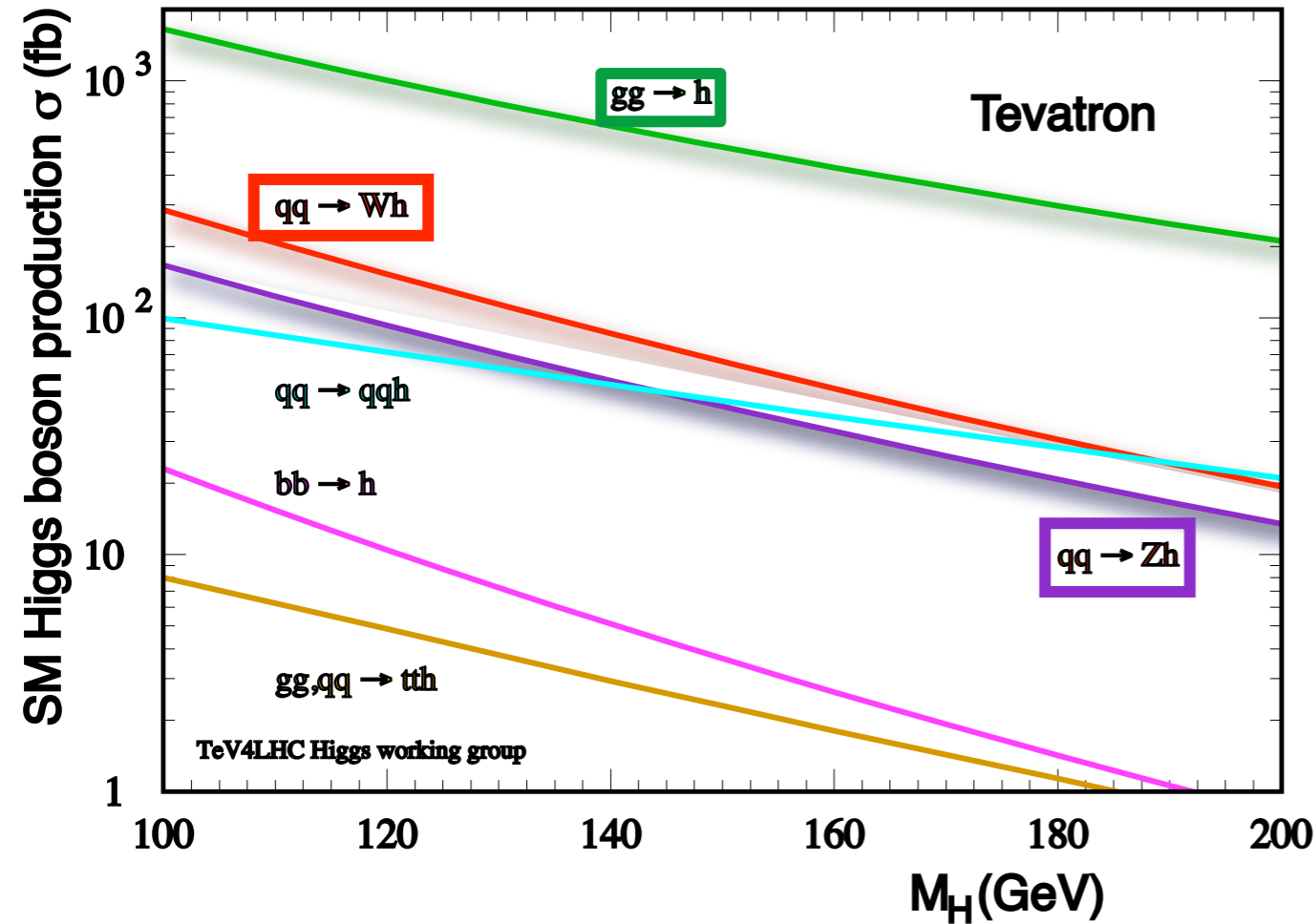
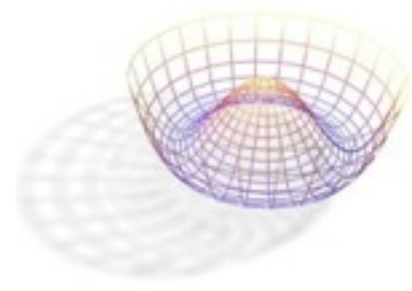
- but very clean, and with good mass resolution



Found 3 candidate events at low mass:

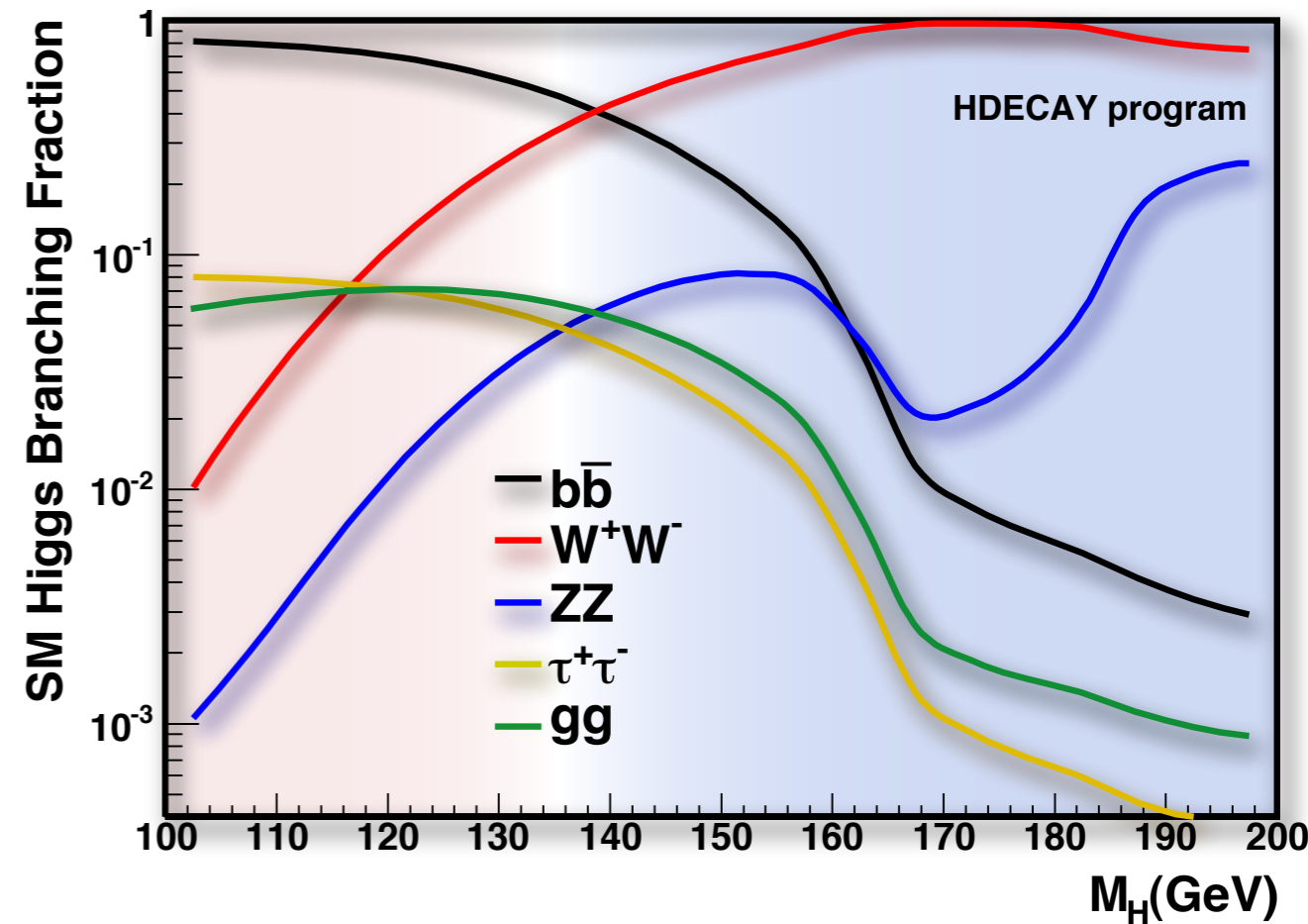
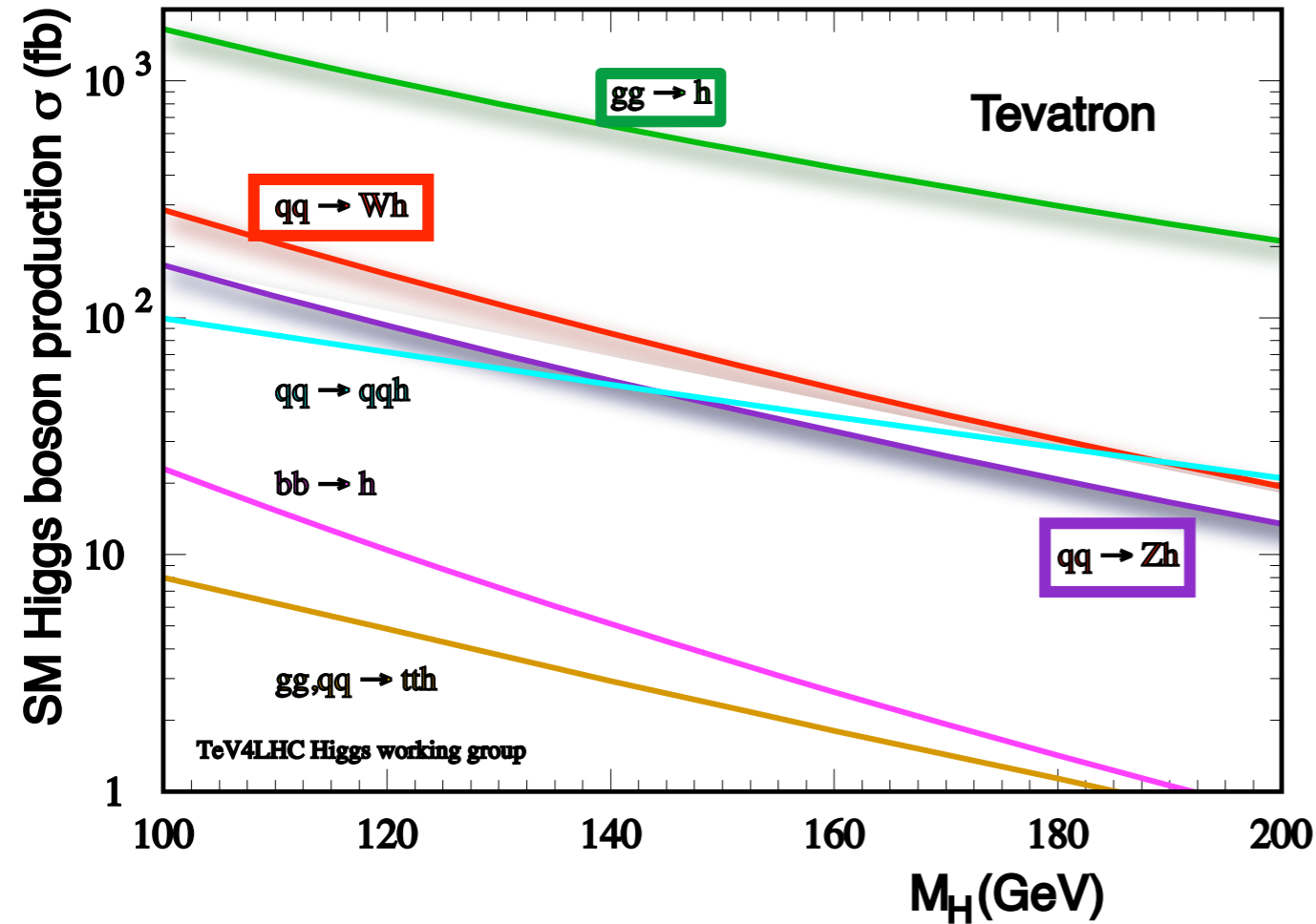
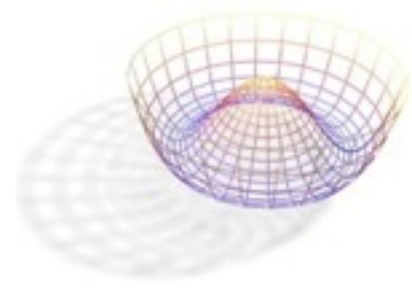
- 2 in $e^+e^-\mu^+\mu^-$ final state (124.3 GeV, 123.6 GeV)
- 1 in $\mu^+\mu^-\mu^+\mu^-$ final state (124.6 GeV)

Higgs Boson Production and Decay



- Total inelastic scattering cross section (strong interaction) ~ 60 mb:
- background suppression by 10-11 orders of magnitude required
 - use signatures not overwhelmed by the strong interaction

Higgs Boson Production and Decay

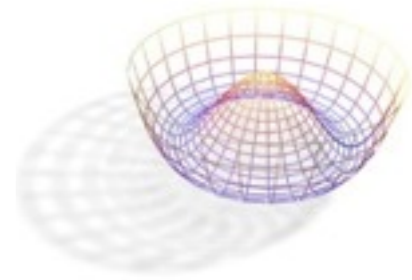


Strategy: use leptons!

- low M_H (≈ 135 GeV): VH associated production, leptonic V decay ($V=W,Z$)
- high M_H (≈ 135 GeV): $H \rightarrow W^+W^-$, both W bosons decaying leptonically

A straightforward strategy, but leading to a large number of final states

The Tevatron Collider

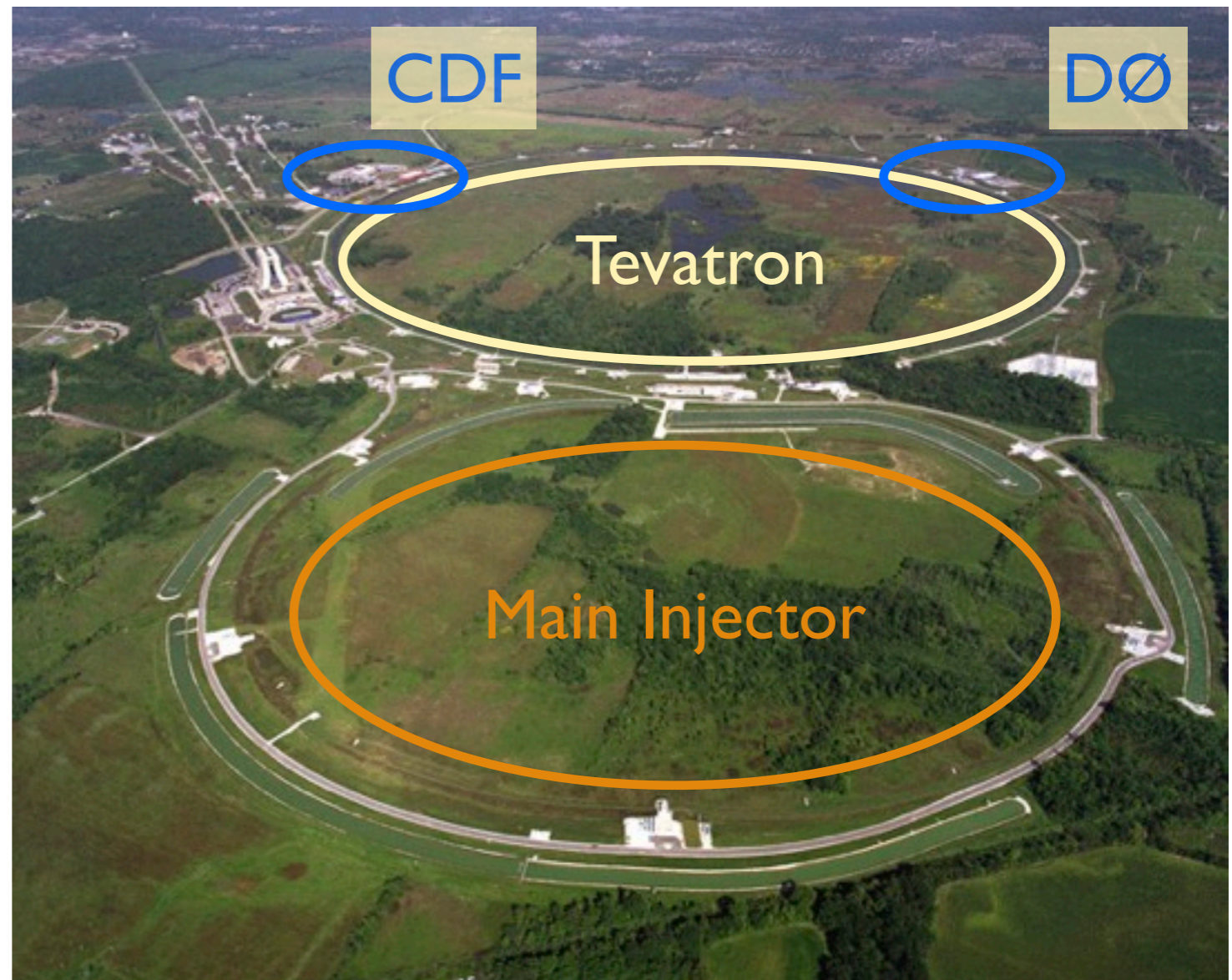


$p\bar{p}$ collisions,
1.96 TeV

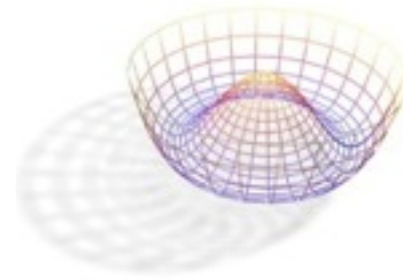
$\sqrt{s} =$

mature collider and
experiments

- running since 2001



The Tevatron Collider

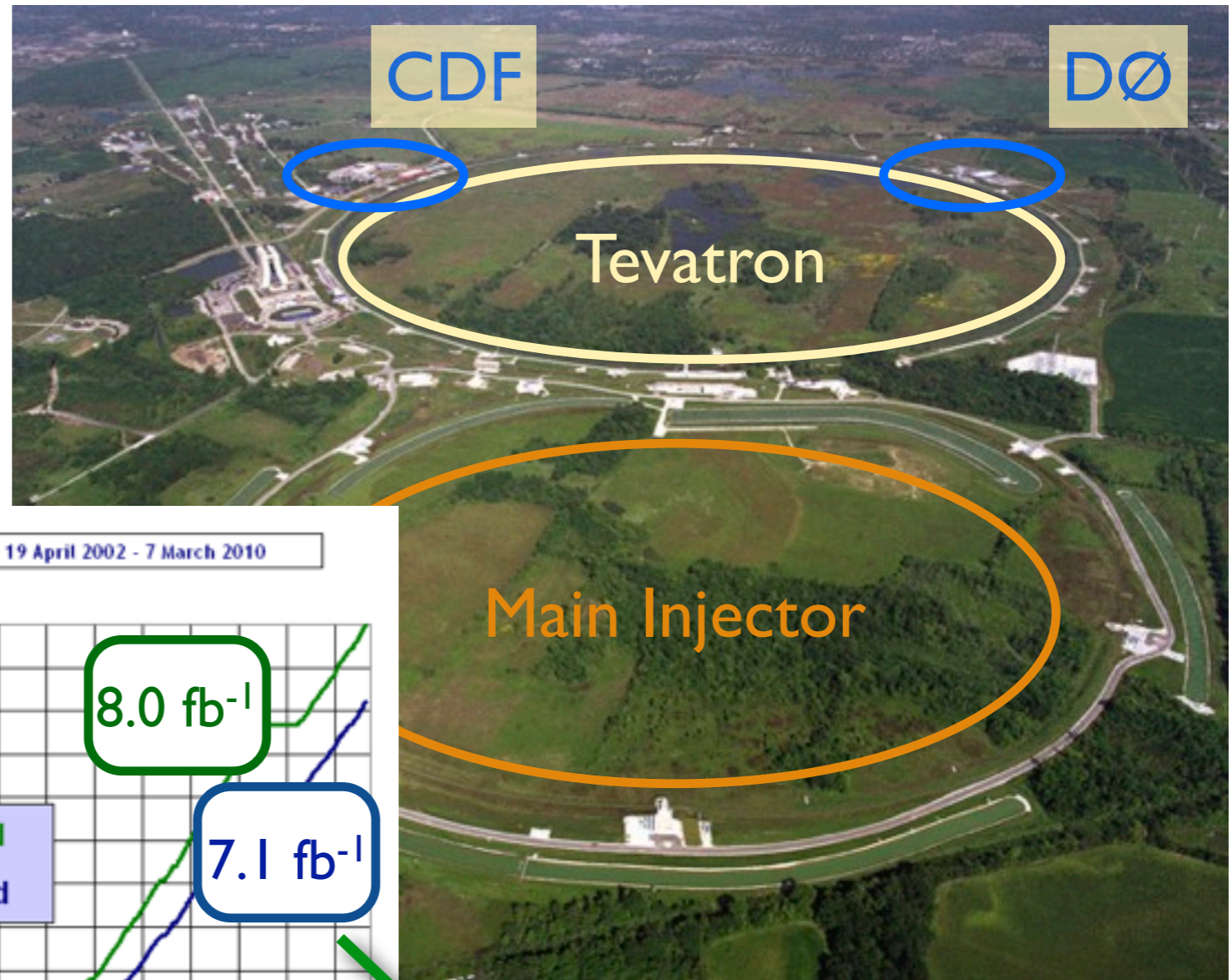


$p\bar{p}$ collisions,
1.96 TeV

$\sqrt{s} =$

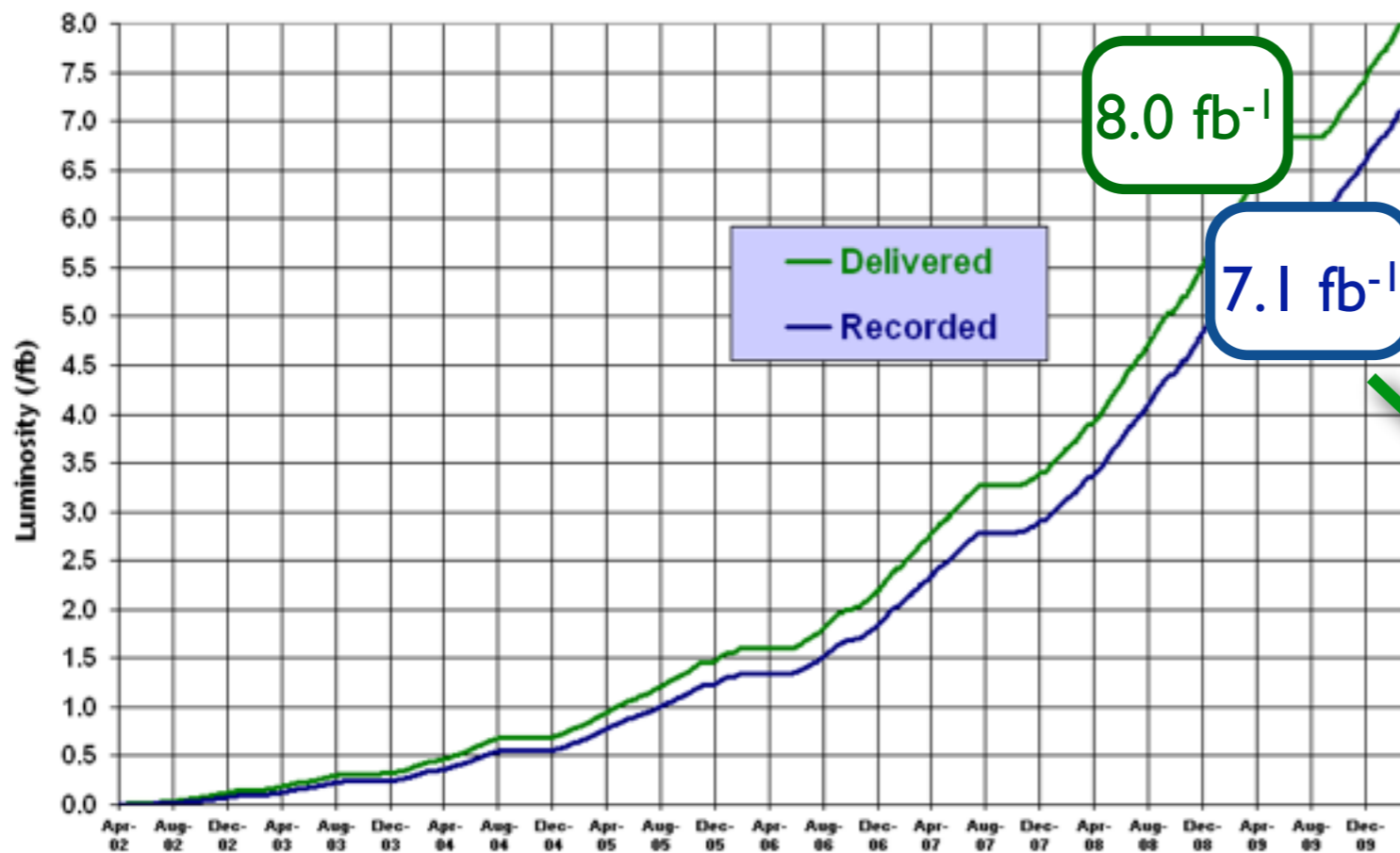
mature collider and
experiments

- running since 2001



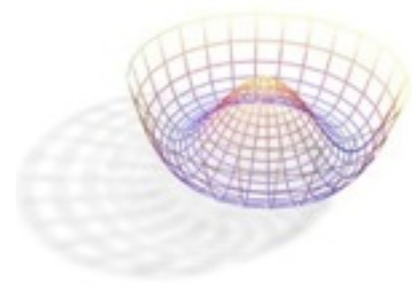
Run II Integrated Luminosity

19 April 2002 - 7 March 2010



- 1 fb = 10^{-43} cm²
- if $\sigma = 1$ fb: need $L = 1$ fb⁻¹ to produce one event
- many interesting processes have $\sigma \sim 100-1000$ fb

Limits



No significant signal-like excess observed... \Rightarrow set limits

Procedure:

- Compare data compatibility with $s+b$ / b -only hypotheses (each M_H)

$$Q = \frac{\mathcal{L}(s+b|m_H)}{\mathcal{L}(b)} = \prod_{i \in \text{bins}} \frac{e^{-(s_i+b_i)} (s_i+b_i)^{n_i}}{n_i!} / \frac{e^{-b_i} b_i^{n_i}}{n_i!}$$

- Calibrate outcome with **toy experiments**

- Compare resulting distributions with observed Q

$CL_{b/s+b} \equiv$ fraction of background-only/signal+bg experiments less signal-like than data

Reject $s+b$ hypothesis if $CL_{s+b} < 0.05$

