

Practice exercise for Quantum Mechanics 3

Extra 1: Polarization and elastic scattering

Consider elastic collisions between spin-1/2 particles (such as nucleons) and spin-0 particles (such as pions). The corresponding scattering potential is in general spin dependent. This implies that the scattering amplitude becomes a scattering operator \hat{F} that acts as a 2×2 matrix F on the spin-1/2 spin space. For a given pure initial spin state $|\chi^i\rangle$ of the spin-1/2 beam particles and a pure final spin state $|\chi^f\rangle$ of the scattered spin-1/2 particles, the differential cross section is given by

$$\frac{d\sigma_{i \rightarrow f}}{d\Omega} = |\langle \chi^f | \hat{F} | \chi^i \rangle|^2 .$$

Scenario 1: sometimes the polarization of the scattered spin-1/2 particles is not measured. As a result, the differential cross section can be summed over the two independent final spin states $|\chi^f\rangle = |\chi_{\frac{1}{2}, \frac{1}{2}}\rangle$ and $|\chi^f\rangle = |\chi_{\frac{1}{2}, -\frac{1}{2}}\rangle$.

- (i) Suppose the incoming spin-1/2 beam is in the pure spin state $|\chi^i\rangle$. Show that the differential cross section is in that case equal to the expectation value of the operator $\hat{F}^\dagger \hat{F}$ for this spin state $|\chi^i\rangle$.
- (ii) Next we switch to an incoming spin-1/2 beam with polarization vector \vec{P} . Demonstrate that the spin-ensemble average of the differential cross section is given by

$$\frac{d\bar{\sigma}}{d\Omega} = \frac{1}{2} \text{Tr}(F^\dagger F) + \frac{1}{2} \vec{P} \cdot \text{Tr}(\vec{\sigma} F^\dagger F) .$$

Scenario 2: an interesting scattering phenomenon is that the scattering process itself may give rise to polarization, irrespective of the initial polarization of the incoming beams.

- (iii) The scattering process turns a pure initial spin state $|\chi^i\rangle$ into a pure final spin state $|\chi^{\text{sc}}\rangle = c |\hat{F} \chi^i\rangle$ for the scattered spin-1/2 particles, where c is a spin-independent complex normalization factor that takes into account that only part of the incoming particles is actually scattered. Why is $|\chi^{\text{sc}}\rangle$ of this form?
- (iv) Consider an incoming spin-1/2 beam with density operator

$$\hat{\rho} = \sum_i W_i |\chi^i\rangle \langle \chi^i| .$$

Derive the following density operator for the scattered spin-1/2 particles:

$$\hat{\rho}' = |c|^2 \hat{F} \hat{\rho} \hat{F}^\dagger , \quad \text{with} \quad |c|^2 = 1 / \text{Tr}(\hat{F} \hat{\rho} \hat{F}^\dagger) .$$

- (v) Finally, assume the incoming spin-1/2 beam to be unpolarized. Determine the polarization vector \vec{P}' of the scattered spin-1/2 particles.